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Influence of Chain Length and Derivatization on the Lowest Singlet and Triplet States and Intersystem Crossing in Oligothiophenes

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Abstract: On the basis of configuration interaction calculations, we first describe the nature of the lowest singlet and triplet excited states in oligothiophenes ranging in size from two to six rings. We calculate the vertical excitation energies from the singlet ground state S0 to the first one-photon allowed singlet excited state S1 as well as the energy difference between the ground state and the lowest triplet state T1. The computed transition energies are in very good agreement with the measured values and indicate a strong confinement of the lowest triplet. We also uncover the nature of the higher-lying triplet excited state Tn that is coupled via a large oscillator strength to T1. The evolution with chain length of the T1−Tn excitation energies compares well with the experimental evolution based on photoinduced absorption data. Next, we investigate the geometry relaxation phenomena occurring in the S1 and T1 states; more pronounced and localized bond-length deformations are calculated in the triplet state than in the singlet, confirming the more localized character of T1. We also analyze the influence on the lowest excited states of grafting electroactive end-groups on the conjugated path of terthiophene. Finally, the various mechanisms involved in the nonradiative decay of the singlet excitations are discussed, and results are presented as a guide toward the optimization of light emission efficiency in conjugated systems.

I. Introduction

Due to their exciting electronic and optical properties, organic conjugated polymers have been the focus of an enormous amount of experimental and theoretical work in the last decade. Their high electrical conductivities upon doping and fast, strong nonlinear optical response make them suitable for a large number of potential applications, for instance in transistors, batteries, sensors, light amplificators, or all-optical switches.1 It has also been shown that conjugated polymers can be used as active layers in light-emitting diode (LED) devices.2,3 In this context, quantum-chemical calculations prove to be useful since they allow the characterization of both the singlet excited states involved in the emission mechanism and the triplet excited states, either formed directly by electron-hole recombination or resulting from an intersystem crossing between the singlet and triplet manifolds; note that such a crossing gives rise to a strong, long-lived, triplet−triplet absorption observed by photoinduced absorption (PA) experiments. Recently, we have presented a model including both electron−phonon and electron−electron interactions that provides a coherent description of the nature of the lowest singlet and triplet excitations in model oligomers of poly(paraphenylene vinylene) (PPV), a conjugated polymer which has been extensively used as active material in the LEDs.4 After PPV, similar LEDs have been successfully built by sandwiching oligothiophenes in between electron- and hole-injecting contacts.5−9 The choice for thiophene oligomers is mainly guided by the following reasons: (i) well-defined oligomers have been synthesized that possess very good processability and stability10 and (ii) the electronic photoexcitations of oligothiophenes have been well characterized experimentally;10−15 most importantly, fluorescence has been observed in thiophene oligomers,10−15 which has to be related to the fact that the lowest singlet excited state S1 in these molecules has a different parity from the ground state S0 and...
that, therefore, the \( S_1 \rightarrow S_0 \) electronic transition is dipole allowed.\(^{15-17}\) (Note, however, that even for a dipole-forbidden \( S_1 \) excited state, weak light emission can be induced by vibronic couplings.)

In this work, we apply techniques similar to those used for the description of the excitations in PPV oligomers\(^2\) to uncover the nature of the lowest singlet and triplet excited states of unsubstituted oligothiophenes ranging in size from two to six thiophene units (hereafter denoted as Th\(_n\), where \( n \) corresponds to the number of thiophene rings). We analyze the influence of the chain length on the singlet–singlet \( S_0 \rightarrow S_1 \) and triplet–triplet \( T_1 \rightarrow T_2 \) excitation energies as well as on the \( S_0 \rightarrow T_1 \) energy difference and compare our results to experimental data provided by one-photon optical absorption and photoinduced absorption measurements. We also investigate the lattice relaxations occurring in the \( S_1 \) and \( T_1 \) excited states of some model oligomers. Finally, we discuss some of the processes that lead to nonradiative decay of the singlet excitations and try to rationalize some experimental results indicating an increase of the fluorescence quantum yield either by increasing the chain length from two (Th\(_2\)) to six (Th\(_6\)) thiophene rings\(^3\) or by attaching electroactive groups on terthiophene (Th\(_3\)).\(^18\)

In this context, we study the effects on the lowest excited states of Th\(_3\) or grafting both electron-donating (OCH\(_3\) and NH\(_2\)) and electron-withdrawing (CHO and CN) moieties on the conjugated segment.

The structure and atom labeling of the various molecules investigated in this work can be deduced from those of the longest unsubstituted oligothiophene (Th\(_6\)) and those of the monosubstituted (Th\(_3\)-X, where \( X \) denotes the nature of the substituent) and bis-substituted (X-Th\(_3\)-X) trimers of thiophene, displayed in Figure 1.

## II. Methodology

The geometries of all oligomers are optimized at the Hartree–Fock semiempirical Modified Neglect of Differential Overlap (MNDO)\(^19\) level. In the case of oligothiophenes, the MNDO technique has been shown to give bond length values in the ground state in good agreement with X-ray diffraction data.\(^20\) Note that all the calculations were performed on the basis of planar conformations. Crystal packing is likely to force the molecules to be planar in films,\(^21\) whereas torsions of the rings along the chain axis most probably occur in solution.\(^10\)

Based on the ground-state optimized geometries, we perform Intermediate Neglect of Differential Overlap (INDO)\(^22\)/Multi-

Reference Double-Configuration Interaction\(^23\) calculations to describe the lowest singlet and triplet excited states. The MRD-CI scheme has been described in detail elsewhere\(^24\) and successfully applied in the case of polyenes\(^25,26\) as well as oligothiophenes.\(^17\) oligopyrroles,\(^26\) and oligo(phenylene vinylene).\(^5\) Here, we consider in the CI expansion of the wave functions single and double excitations from the six highest \( \pi \)-levels to the four \( \pi^* \)-levels with respect to two reference configurations (dominating the description of the ground state \( S_0 \) and the lowest excited states \( S_1 \) and \( T_1 \)): (i) the Self-Consistent-Field (SCF) determinant itself and (ii) the configuration resulting from the excitation of one electron from the highest occupied molecular orbital (HOMO) to the lowest unoccupied molecular orbital (LUMO). In the case of the PPV oligomers, we found no major changes in the transition energies from the singlet ground state to the lowest excited states when considering from two to five reference determinants (including doubly excited determinants) in the MRD-CI formalism. As all truncated CI techniques, the MRD-CI scheme is not size consistent; however, an appropriate choice of the reference configurations provides a balanced description of the correlation effects in the ground state and in the excited states and therefore allows for a precise estimation of the transition energies to the lowest excited states, as has been shown elsewhere.\(^5,17,24-26\)

To investigate the lattice distortions taking place in the \( S_1 \) and \( T_1 \) excited states of oligothiophenes, we have considered the following approach. We first optimize the geometries of both states by means of the MNDO technique. The SCF calculations are then performed on the basis of a level occupancy characterized by the promotion of one electron from the HOMO to the LUMO level (flipping one spin for the triplet). The excited-state energy, obtained by a simple CI calculation involving a few electronic configurations, is used as a criterion during the geometry optimization. For oligo(phenylene vinylene)s, this technique leads to excited-state geometries that are very similar to those obtained on the basis of a bipolaron-
type relaxation (where geometry changes related to the formation of bipolaron defects are simulated) or BOBL-type relaxation (where bond lengths in the excited states are evaluated from the electronic bond orders).5

In the second step, on the basis of the MNDO-optimized excited-state geometries, INDO/MRD-CI vertical transition energies from the ground state to the lowest excited states are computed. Combining the ground-state deformation energy calculated at the MNDO level when going from the S0 geometry to the S1 or T1 geometry to the INDO/MRD-CI vertical S0→S1 and S0→T1 energy differences allows us to estimate the geometric relaxation energy differences in both the singlet S1 and triplet T1 excited states.

According to experiment and theory,27 the optimized 1B1u excited-state geometry of polyacetylene (PA) presents a lattice conformation corresponding to the formation of a soliton–antisoliton pair. Solitons are topological kinks, defining a domain wall between two geometric phases, characterized by an opposite alternation of the single and double carbon–carbon bonds; alternatively, they can be viewed as excitations of the twofold degenerate ground-state PA chains leading from one minimum of the potential to the other minimum of same energy.27 In oligothiophenes, soliton–antisoliton pairs are bound by the ground-state nondegeneracy (in opposite to the situation observed for polyacetylene, the two resonant geometric forms, obtained by exchanging the single and double bonds in the conjugated pathway have different energies), and we then parameterize the geometric deformations that characterize the excited states by fitting to the MNDO-optimized geometries the expression adopted by Sethna and Kivelson26 in the case of a loosely bound soliton–antisoliton pair

\[ r_n^e = r_n^g + (-1)^n \times \frac{\phi}{2\xi} + \frac{n - N_0 - R}{\xi \tanh \frac{n + 1 - N_0 - R}{\xi} - \frac{n - N_0 + R}{\xi \tanh \frac{n + 1 - N_0 + R}{\xi}}} \]

where \( r_n^g \) is the bond length between carbon atoms \( n \) and \( n+1 \) in the MNDO-optimized ground-state geometry, \( N_0 \) is the center of the chain, \( \xi \) is the coherence length of the defect, \( \phi \) is the amplitude of the deformations, and \( 2\xi \) is the distance between the “soliton” centers; the smaller the \( \phi \) value, the stronger the confinement of the charged species. In Figure 2, we present the bond-length distortions \( (r_n^e - r_n^g) \) calculated for various values of the parameters \( R, \phi, \) and \( \xi \). Note that the carbon–sulfur bonds, that are only slightly modified in the excited states with respect to their values in the ground state,17,27 are not included in this fitting procedure.

III. Vertical Excitation Energies

The ground-state geometry of all oligomers has been optimized at the MNDO level. Within the thiophene rings, the difference between the single and double carbon–carbon bond lengths, \( \Delta r \), amounts to 0.05–0.06 Å, and the carbon–sulfur C–S bonds are close to 1.7 Å; the interring carbon–carbon C–C bonds are calculated to be equal to \( \sim 1.45 \) Å. These results are in very good agreement with the available X-ray diffraction data, which indicate values of \( \Delta r \) and C–S bond lengths close to 0.07 and 1.7 Å, respectively, and interring C–C bond lengths of approximately 1.46 Å.30 At the 6-31G*/MP2 level, the corresponding values for the planar conformation of biphenylene are \( \Delta r = 0.03–0.04 \) Å; \( r(C-C) \sim 1.73 \) Å; and interring \( r(C-C) \sim 1.45 \) Å.31 We are thus confident that the MNDO technique is appropriate to reproduce the equilibrium geometries of thiophene oligomers.

On the basis of the MNDO ground-state geometries, we have first calculated the electronic excitation energies from the singlet ground state \( S_0 \) to the lowest singlet excited state \( S_1 \), by means of the INDO/MRD-CI approach. Note that the CI expansion of the \( S_1 \) wave function is dominated by the electronic configuration built by promoting one electron from the HOMO level to the LUMO level. The theoretical \( S_0 \rightarrow S_1 \) transition energies are compared in Table 1 to the experimental values, extracted from optical absorption measurements in solution;11,14 the measured and calculated evolutions with the inverse number of thiophene rings of the \( S_0 \rightarrow S_1 \) excitation energies are illustrated in Figure 3. The agreement between the experimental and theoretical results is impressive, which shows the suitability of the INDO/MRD-CI technique for the description of the electronic excitations in these oligomers. Both the spectroscopic data and the calculated values indicate a red-shift of the first one-photon allowed electronic transition with the size of the oligomers, due to the extension of the \( \pi \)-delocalized system. Extrapolation to an infinite chain length of the calculated transition energies, assuming a linear evolution of \( S_0 \rightarrow S_1 \) with respect to \( 1/n \), leads to an optical gap close to 2 eV, which again is in good agreement with the spectroscopic data on polythiophene; recent optical measurements performed on regioregular polythiophenes show indeed an absorption with an onset at \( \sim 1.7–1.8 \) eV and a peak at \( \sim 2.5 \) eV.32,33

In Table 1, we also report the energy difference between the \( S_0 \) ground state and the \( T_1 \) lowest triplet excited state, as calculated at the INDO/MRD-CI level on the basis of the \( S_0 \) geometry. As illustrated in Figure 3, the evolution with chain length of the \( S_0 \rightarrow T_1 \) energy difference is much slower than that of the \( S_0 \rightarrow S_1 \) excitation: the singlet–triplet energy difference is only lowered by \( \sim 0.2 \) eV when going from the dimer to the hexamer of thiophene while the corresponding singlet–singlet absorption is characterized by a bathochromic shift that amounts to \( \sim 1.4 \) eV. The slower evolution with the number of thiophene rings of the triplet excitation energy (with respect to that of the singlet) reflects its more localized character. Such a difference has clearly to be related to the exchange term.

Table 1. INDO/MRD-CI Singlet–Singlet \( S_0 \rightarrow S_1 \), Singlet–Triplet \( S_0 \rightarrow T_1 \), and Triplet–Triplet \( T_1 \rightarrow T_1 \) Energies, Calculated on the Basis of the Ground-State Geometries of Oligothiophenes

<table>
<thead>
<tr>
<th></th>
<th>Th2</th>
<th>Th3</th>
<th>Th4</th>
<th>Th5</th>
<th>Th6</th>
<th>Th7</th>
<th>Th8</th>
<th>Th9</th>
<th>Th11</th>
</tr>
</thead>
<tbody>
<tr>
<td>( S_0 \rightarrow S_1 ) calcd</td>
<td>4.23</td>
<td>3.58</td>
<td>3.16</td>
<td>2.92</td>
<td>2.72</td>
<td>2.58</td>
<td>2.45</td>
<td>2.40</td>
<td></td>
</tr>
<tr>
<td>( S_0 \rightarrow T_1 ) calcd</td>
<td>4.37</td>
<td>3.42</td>
<td>3.01</td>
<td>2.62</td>
<td>2.40</td>
<td>2.20</td>
<td>2.21</td>
<td>1.75</td>
<td>1.60</td>
</tr>
<tr>
<td>( S_0 \rightarrow S_1 ) exp</td>
<td>1.84</td>
<td>1.68</td>
<td>1.61</td>
<td>1.61</td>
<td>1.65</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Footnotes:
1. Reference 11.
2. Reference 14.
3. Reference 41.
4. Reference 36.
5. Experimental data taken from optical measurements in solution are also included. All the energies are given in eV.
that allows a more pronounced confinement of the triplet exciton with respect to the singlet. Very similar conclusions have been drawn from similar calculations on PPV oligomers.\textsuperscript{5} Note that Optical-Detected Magnetic Resonance (ODMR) experiments on polythiophene indicate that the T\textsubscript{1} triplet state extends over not much more than one thiophene ring,\textsuperscript{34} which agrees well with our results. Moreover, similar ODMR measurements performed on oligomers ranging in size from two to five rings\textsuperscript{35} indicate a very small chain-length dependence of the average spin–spin distance, \( r \), in the triplet, as deduced within the dipolar approximation from the zero-field splitting, \( D \) (using the relation \( D/hc = 2.6017r^{-3} \), where \( D/hc \) is in cm\textsuperscript{-1} and \( r \) in Å, we find \( r = 3.27, 3.36, \) and 3.49 Å, in Th3, Th4, and Th5, respectively).

Optical absorption measurements in a solvent containing heavy atoms (i.e., 1,2-dibromomethane) have allowed the T\textsubscript{1} excited state in thertiophene (Th3) to be located.\textsuperscript{36} In this case, the intensity of the S\textsubscript{0} \( \rightarrow \) T\textsubscript{1} excitation is raised due to the increase of the spin-orbit coupling related to the “heavy-atom effect” induced by the solvent molecules.\textsuperscript{37} The experimental absorption spectrum shows a broad and weak band that peaks at \( \sim 1.7 \text{ eV} \), in very good agreement with the 1.68 eV calculated value. Recently, Janssen \textit{et al.} have observed that addition of C\textsubscript{60} to oligothiophenes (from Th6 to Th11) in solution results in a quenching of the thiophene oligomers triplet state via energy transfer to C\textsubscript{60} and produces triplet-state C\textsubscript{60}.\textsuperscript{14} Therefore, the triplet T\textsubscript{1} state of oligothiophenes ranging in size from six to eleven rings can be estimated to lie between the triplet energy of C\textsubscript{60} (1.57 eV)\textsuperscript{38,39} and the triplet energy of Th3 (1.71 eV).\textsuperscript{36} Moreover, the efficiency of the energy transfer (from the triplet state of Thn to the triplet state of C\textsubscript{60}) was found to slightly depend on the length of the conjugated oligomers and to be less efficient in the longest chains; this feature can be explained by a small red-shift of the S\textsubscript{0} \( \rightarrow \) T\textsubscript{1} transition with extension of the \( \pi \)-system, consistent with the theoretical evolution. We finally note that Xu and Holdcroft have measured a phosphorescence peak at \( \sim 1.5 \text{ eV} \) in polythiophene,\textsuperscript{40} a value very close to the one we can extrapolate for an infinite chain length (1.49 eV) on the basis of the calculated oligomer excitation energies.

One of the possible decay routes of the singlet excitons in conjugated systems consists in an intersystem crossing to form a metastable triplet exciton. Such a triplet gives rise to a single, intense triplet–triplet T\textsubscript{1} \( \rightarrow \) T\textsubscript{1} electronic transition, which has been observed by photoinduced absorption (PA) measurements on oligothiophenes in solution.\textsuperscript{14,41} The INDO/MRD-CI calculations indicate that the T\textsubscript{1} state is mainly coupled to one triplet excited state (hereafter denoted T\textsubscript{3}), via a large oscillator strength whose value is comparable to that of the S\textsubscript{0} \( \rightarrow \) S\textsubscript{1} excitation. The T\textsubscript{1} wave function, where \( n \) varies from 8 to 10 depending on chain length, is quite complex and includes both singly and doubly excited configurations with respect to the Hartree–Fock determinant. The T\textsubscript{1} \( \rightarrow \) T\textsubscript{3} transition energies are collected in Table 1 together with the experimental data obtained from the PA measurements.\textsuperscript{14,41} Although the theoretical calculations lead to a systematic overestimation of the triplet–triplet excitation energies, the calculated chain-length dependence of the T\textsubscript{1} \( \rightarrow \) T\textsubscript{3} transition energy is in good agreement with the experimental one and shows a strong red-shift when increasing the chain length (Figure 4). It thus appears that, as is the case for PPV oligomers,\textsuperscript{5} the T\textsubscript{3} state in oligothiophenes is not as localized as T\textsubscript{1}, a feature which can be accounted for by the fact that, as the T\textsubscript{3} excited state lies at higher an energy than T\textsubscript{1}, it has a smaller binding energy and hence its wave function is less confined.

IV. Geometry Relaxation Phenomena

The equilibrium bond lengths in the S\textsubscript{1} and T\textsubscript{1} excited states of the dimer (Th2), trimer (Th3), and tetramer (Th4) of thiophene have been calculated at the MNDO level, by considering the approach described in the methodology section.\textsuperscript{32} The typical bond-length modifications obtained from these calculations lead to the appearance of a quinonoid character, fully

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure2.png}
\caption{Bond-length deformations (in Å) calculated on the basis of eq 1 with (a) \( \phi = 0.1; \xi = 2; \) and \( R = 2 \) (circles); (b) \( \phi = 0.1; \xi = 2; \) and \( R = 1 \) (squares); (c) \( \phi = 0.1; \xi = 1; \) and \( R = 1 \) (crosses); and (d) \( \phi = 0.1; \xi = 1; \) and \( R = 2 \). \( \phi \) is in Å, \( \xi \) and \( R \) are given in site units.}
\end{figure}

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure3.png}
\caption{Comparison between the evolutions with the inverse number of thiophene rings, \( 1/n \), of (i) the INDO/MRD-CI singlet–triplet S\textsubscript{0} \( \rightarrow \) S\textsubscript{1} (solid line, closed circles) and singlet–triplet S\textsubscript{0} \( \rightarrow \) T\textsubscript{1} (dashed line, open circles) energies and (ii) the experimental S\textsubscript{0} \( \rightarrow \) S\textsubscript{1} (closed squares,\textsuperscript{11} closed triangles\textsuperscript{14} and S\textsubscript{0} \( \rightarrow \) T\textsubscript{1} (open squares\textsuperscript{36}) energies obtained from measurements in solution.}
\end{figure}


Singlet and Triplet States in Oligothiophenes

Figure 4. Comparison between the evolutions with the inverse number of thiophene rings, 1/n, of (i) the INDO/MRD-CI triplet–triplet T₁ → T₄ excitation energies (solid line, closed circles) and (ii) the experimental T₁ → T₄ (open squares, open triangles) transition energies obtained from photoinduced absorption measurements in solution.

Figure 5. Evolution with site position, i, on the thiophene tetramer, of the bond-length alternation, ∆r, calculated at the MNDO level as the difference between the lengths of the (i,i+1) and (i,i−1) carbon–carbon bonds, in the S₀ (solid line), S₁ (dashed line), and T₁ (dotted line) states.

consistent with the formation of a polaron-like defect.43 We note that, for the monomer and planar dimer of thiophene, similar geometric deformations have been predicted at the 3-21G*/CIS ab initio level and from semiempirical calculations including single and double excitations (QCFF/PI + CISD technique).29 Finally, it should be emphasized that, using a similar geometric deformations have been predicted at the MNDO level as the difference between the lengths of the (i,i+1) and (i,i−1) carbon–carbon bonds, in the S₀ (solid line), S₁ (dashed line), and T₁ (dotted line) states.

From Table 2 and Figure 5, we can draw the following conclusions.

(i) Rather small lattice distortions occur in the S₁ excited state, which translates into small relaxation energies (on the order of 0.1—0.2 eV). The geometry deformations lead to a reduction in bond-length alternation along the conjugated path formed by the carbon atoms; alternatively, one can view the same bond-length modifications as giving rise to the appearance of some quinoid character on the thiophene rings. The absolute value of the bond-length alternation in Th₄ drops from ~0.04—0.05 Å in the external rings to ~0.0—0.02 Å in the central part of the molecule. These lattice distortions are similar to those associated with the formation of charged soliton–antisoliton pairs in polyacetylene;27,45 however, in oligothiophenes, the electron-hole pairs are bonded in the S₁ excited state by (i) the Coulombic attraction between opposite charges; (ii) the non-degenerate ground state of the thiophene oligomers; and (iii) the finite size of the systems, which prevents the separation between the opposite charge carriers occurring upon photoexcitation in polyacetylene.27,45 From the R value obtained by fitting eq 1 to the singlet excited-state geometry of Th₄ and the comparison between the bond-length deformations in the different oligomers, we estimate the geometric defect to extend over 3—4 thiophene units.

(ii) Much more pronounced lattice distortions take place in the lowest triplet excited state, as indicated by the larger amplitude, φ, value found for T₁. The relaxation energies with respect to vertical excitations are significantly larger than those found for the singlet and on the order of 0.3—0.4 eV. In Th₄, the bond-length alternation absolute value, ∆r, in the external rings is close to 0.05—0.06 Å, i.e., the value calculated in the ground state. Going from the ends of the chain toward the center, ∆r decreases (the single bonds are shortened and the double bonds are elongated), reaches 0 at the connection between the first and second rings (the bonds are there equal), and then evolves strongly and peaks at the center of the tetramer (where the single and double bond characters have exchanged with respect to the ground state and the maximum absolute value of the bond-length alternation is recovered). In the triplet excited state as in the singlet, the formation of a soliton–antisoliton pair clearly emerges from the evolution of the geometry deformations along the chain axis. For essentially the same reasons as for the singlet defect, the electron-hole pair is strongly bonded in the triplet excited state; due to the exchange potential, the confinement of the triplet is even more pronounced than that of the singlet and, from the R value fitted for T₁, it seems that 1—2 thiophene ring(s) is(are) required for a proper accommodation of the triplet defect, in accord with the ODMR data.34,35 In fact, the MNDO results indicate that, in the triplet T₁ state, the bond lengths in the two central rings of Th₄ are close to those calculated for the same state in Th₂; the external rings of the tetramer are almost unaffected when going from S₀ to T₁.

Table 2. Parameters in Eq 1 (φ in Å, ξ and R in site units) Obtained by Fitting to the MNDO-Optimized Geometries of the Singlet S₁ and Triplet T₁ Excited States of the Thiophene Tetramer (Th₄).

<table>
<thead>
<tr>
<th></th>
<th>φ</th>
<th>ξ</th>
<th>R</th>
</tr>
</thead>
<tbody>
<tr>
<td>S₁</td>
<td>0.019</td>
<td>2.25</td>
<td>4.57</td>
</tr>
<tr>
<td>T₁</td>
<td>0.029</td>
<td>1.58</td>
<td>3.45</td>
</tr>
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</table>

The substitution effects on the \( S_0 \), \( S_1 \), and \( T_1 \) geometries of terthiophene have been investigated at the MNDO level. The INDO/MRD-CI \( \pi \)-charge distributions, calculated for the unsubstituted, monosubstituted, and bis-substituted trimers, are given in [e] in Table 3. Similar amounts of charge transfer, \( ∼0.1 \text{ eV} \), occur for the unsubstituted and substituted oligomers, that amounts to \( ∼0.4\text{ eV} \) when going from \( \text{Th3-OCH}_3 \) oligomer to the \( \text{CHO-Th3-OCH}_3 \) oligomer. The experimental data indicate a red-shift of the maximum absorption between the unsubstituted and substituted oligomers, that is, similar amounts of charge transfer.

The substitution effects on the \( S_0 \), \( S_1 \), and \( T_1 \) geometries of terthiophene have been investigated at the MNDO level. The INDO/MRD-CI \( \pi \)-charge distributions, calculated for the unsubstituted, monosubstituted, and bis-substituted trimers, are collected in Table 3, and the singlet—singlet, singlet—triplet, and triplet—triplet vertical excitation energies are reported in Table 4. Analysis of these results leads to the following observations:

(i) The effects on the ground-state geometry of Th3, of the donor/acceptor groups grafted to the conjugated segment are very local; with respect to Th3, only the geometry of the thiophene ring(s) to which the substituent(s) is(are) attached is significantly modified. These geometric deformations lead to a small decrease in bond-length alternation in the external rings and are rather insensitive to the nature of the end-group(s) we have considered.

(ii) In the singlet excited state, similar geometry deformations are induced by substituting the terthiophene molecule with electroactive moieties. However, we note that, if the main lattice distortions occur in the thiophene ring(s) the end-group(s) is-(are) connected to, the geometry of the other ring(s) is also modified. We explain this difference by the fact that, in \( S_1 \), all the thiophene rings are involved in the charge transfer process, while, in \( S_0 \), the charge provided by the electroactive group(s) is mainly injected to the external ring(s). As in substituted polyenes \(^{46}\) and oligo(phenylene vinylene) derivatives, \(^{47}\) the electronic transition to the first allowed singlet excited state, \( S_1 \), is red-shifted when end-groups are attached to the terthiophene molecule; the amplitude of this shift is proportional to the amount of charge transfer. For similar strengths of the acceptor and the donor (i.e., similar amounts of charge transfer), electron-withdrawing groups have a stronger effect on the \( S_0 \rightarrow S_1 \) excitation energies than electron-donating ones, a feature that is rationalized by the effective chain lengthening induced by grafting acceptor groups.

(iii) In the triplet \( T_1 \) excited state, the bond-length modifications induced by substitution can be related to the evolution of the spin density distributions. As expected from the confined character of the triplet excitation in oligothiophenes, the spin density in the triplet state of Th3 is mainly localized on the central aromatic ring. In the monosubstituted oligomers, we observe an overall increase in spin density both on the substituted ring and on the central ring at the expense of the third ring; hence, more pronounced geometry distortions occur in the former two thiophene units. In the bis-substituted compounds, the two external rings are, of course, equivalent, and the triplet-state geometry is symmetric. The influence of the end-groups we have considered in this work, on the energy of the triplet state with respect to the ground state, can be understood by considering the following effects: (i) attaching end-groups on the conjugated molecule tends to stabilize the triplet state due to an enhancement of its confinement (as indicated by the increase in spin density on the central thiophene unit of the substituted molecules except for \( \text{Th3-NH}_2 \)); (ii) the \( T_1 \) state energy is lowered by raising the electronic density in the central part of the molecule; and (iii) \( T_1 \) gets destabilized when removing electrons from the conjugated path. In the trimers substituted with donors, effects (i) and (ii) (localization and increase in electron density) contribute to stabilize the triplet, while, in acceptor derivatives, effects (i) and (iii) (localization and decrease in electronic density) are opposite, leading to a weaker effective energy stabilization of \( T_1 \). In the case of the \( \text{CHO-Th3-CHO} \) oligomer, the \( S_0 → T_1 \) energy difference is close to the one calculated for \( \text{Th3} \).

(iv) Whatever the mode of substitution, the \( T_n \) state energy is lowered when going from unsubstituted to substituted compounds. For this high-energy lying triplet excited state, as mentioned previously, the localization of the wave function is much less pronounced than in \( T_1 \) and the evolution of the \( T_n \) excited state energy with substitution follows that of the \( S_1 \) singlet excited state. As a consequence, the theoretical triplet—triplet transition energies show a bathochromic displacement when electroactive groups are grafted to the conjugated bridge: this shift depends on the nature of the end-groups and is calculated to be larger for acceptor derivatives.


VI. Decay of the Singlet Excitations

In LED devices based on conjugated polymers, the first step in the light emission process is the recombination of electrons and holes injected at electrodes to form singlet excitons; the second step consists in the radiative decay of the singlet excitations, i.e., luminescence. Competing nonradiative processes provide additional means of decay and therefore reduce the efficiency of the device. These processes are numerous and include interchain processes (e.g., excimer formation), quenching of excitons by extrinsic or conformational defects or by low-lying two-photon states, internal conversion (IC), intersystem crossing (ISC) from the singlet to the triplet manifold, singlet fission (SF), or charge separation (CS). IC corresponds to a nonradiative decay to the ground state through phonons (vibrational quanta) emission; ISC denotes the crossing process that can occur between singlet and triplet potential energy curves, as a consequence of spin-orbit interaction; SF is a bimolecular process that produces triplet excitons from the collision between singlets and requires the triplet state to be located at roughly half the energy of the singlet state (such a nonradiative decay has been shown to occur in polydiacetylenes as well as in tetracene and anthracene crystals). Note that, while ISC and IC are intrinsic molecular phenomena, SF and CS are solid state effects.

In oligothiophenes, the possibility of quenching the singlet excitons by lower-lying two-photon states can be ruled out, since the 2A state has been shown both experimentally and theoretically to lie above the 1B state. Recent time-resolved fluorescence studies on unsubstituted thiophene oligomers (for which migration of the excitons towards trapping centers is minimized by the finite size of the chains and the smaller disorder) in solution (where excimer formation is expected to be less efficient than in the solid state) indicate a sharp increase of the fluorescence quantum yield \( \Phi_f \) with the number of thiophene rings from two to seven and by attaching electroactive end-groups on the terthiophene molecule. In both cases, no significant alteration in the radiative decay rate constant, \( k_R \), was observed, and the evolution of \( \Phi_f \) with chain length and substitution was related to the decrease of the nonradiative decay rate, \( k_{NR} \). It was also found that internal conversion plays a minor role in the nonradiative decay of the singlet excitations and that, for oligothiophenes, intersystem crossing gives the main contribution to \( k_{NR} \).

Radiationless transition theory expresses the intersystem rate constant in terms of (i) a state density factor; (ii) the matrix element of the spin-orbit coupling hamiltonian between the singlet and triplet wave functions; and (iii) an overlap factor, which accounts for the decrease in rate with an increasing energy gap between the two states involved. A precise description of the intersystem crossing process definitively requires the calculation of the spin-orbit coupling interaction (now in the spin-orbit coupling term). Therefore, we believe that, as was first suggested by Ponterini and co-workers, these changes should be related to the variation in the overlap factor. As pointed out by Robinson and Frosch in the case of a series of aromatic hydrocarbons, this factor is primarily controlled by the singlet-triplet energy separation.

The singlet energy difference that we calculate in unsubstituted oligothiophenes is much too large to give efficient singlet-triplet overlap, and hence the probability for \( S_1 \rightarrow T_1 \) crossing is very weak. The INDO/MRD-CI calculations indicate that there is one triplet excited state (\( T_1 \), for which the wave function involves excitations from deep \( \pi \)-levels with large weights on the sulfur atoms, that lies at an energy close to that of \( S_1 \). Although the restricted configuration space we consider in the MRD-CI formalism does not allow us to obtain a precise estimation of the high-lying excited states (the theoretical \( S_1 \rightarrow T_1 \) transition energies are indeed overestimated with respect to the measured values), the INDO/MRD-CI scheme adequately reproduces the low-lying excited states (as indicated by the good agreement found between the measured and calculated \( S_0 \rightarrow S_1 \) and \( S_0 \rightarrow T_1 \) transition energies). We feel therefore confident in using the energy of \( T_4 \) as an important criterion in determining the rate of intersystem crossing from the singlet to the triplet manifold. In Figure 6, we plot the evolution of the \( S_0 \rightarrow S_1 \) and \( S_0 \rightarrow T_1 \) transition energies with respect to the inverse number of thiophene rings. In bithiophene, the triplet \( T_4 \) lies below the singlet \( S_1 \), while for longer chains, due to the larger stabilization of the \( S_1 \) excited state with chain length, the state ordering is reversed, \( S_1 \) being below \( T_4 \). Both the singlet–singlet \( S_0 \rightarrow S_1 \) and singlet–triplet \( S_0 \rightarrow T_4 \) transition energies follow a quasi linear evolution with respect to \( 1/n \); we note, however, that the calculated energy of the triplet \( T_4 \) excited state of Th3 is overestimated by comparison to the value interpolated by considering a linear relationship between the excitation energies and the inverse chain length (Th3 has...
energies, 

\[ \Delta E(S_1 - T_1) \]

and the Inverse Number of Thiophene Rings in the Unsubstituted Oligomers and the Fluorescence Yields, \( \phi_F \), as well as the Nonradiative Decay Rates, \( k_{NR} \). Measured in Solution (\( k_{NR}^{exp} \)) and Calculated on the Basis of Simple Assumptions (\( k_{NR}^{th} \), see Text).

<table>
<thead>
<tr>
<th>Substance</th>
<th>( \Delta E(S_1 - T_1) ) (eV)</th>
<th>( k_{NR}^{th} ) ((10^9 \text{ s}^{-1}))</th>
<th>( k_{NR}^{exp} ) ((10^9 \text{ s}^{-1}))</th>
<th>( \phi_F^{exp} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Th2</td>
<td>0.10 (0.10)</td>
<td>20.1</td>
<td>19.7</td>
<td>0.018</td>
</tr>
<tr>
<td>Th3</td>
<td>0.24 (0.05)</td>
<td>3.70</td>
<td>3.88,8.7,7.4,5.3</td>
<td>0.07,0.05,0.07</td>
</tr>
<tr>
<td>Th4</td>
<td>0.16 (0.13)</td>
<td>1.23</td>
<td>1.54,1.5</td>
<td>0.20,0.18</td>
</tr>
<tr>
<td>Th5</td>
<td>0.18 (0.18)</td>
<td>1.14</td>
<td>0.82,0.74</td>
<td>0.28,0.36</td>
</tr>
<tr>
<td>Th6</td>
<td>0.22 (0.21)</td>
<td>1.13</td>
<td>0.70,0.73</td>
<td>0.42</td>
</tr>
<tr>
<td>Th3-CHO</td>
<td>0.36</td>
<td></td>
<td>0.24</td>
<td>0.50</td>
</tr>
<tr>
<td>TH3-CN</td>
<td>0.34</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Th3-NH2</td>
<td>0.29</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Th3-OCH3</td>
<td>0.25</td>
<td></td>
<td>4.6</td>
<td>0.061</td>
</tr>
<tr>
<td>OHC-Th3-CHO</td>
<td>0.47</td>
<td></td>
<td></td>
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<tr>
<td>NC-Th3-CN</td>
<td>0.40</td>
<td></td>
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<tr>
<td>H2N-Th3-NH2</td>
<td>0.26</td>
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</tr>
<tr>
<td>MeO-Th3-OMe</td>
<td>0.24</td>
<td></td>
<td></td>
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</tr>
</tbody>
</table>

\( ^{a} \) Reference 11. \( ^{b} \) Reference 18. \( ^{c} \) Reference 50.

not been included in the fit). As the crossing between the evolutions of the singlet \( S_1 \) and triplet \( T_4 \) excited states occurs for a chain length close to that of the Th3 oligomer, we associate this discrepancy to the fact that we did not include in the calculations the spin-orbit coupling which, in the crossing region, is expected to mix efficiently the singlet and triplet wave functions.

In Table 5, we collect (i) the \( S_1 - T_4 \) energy differences, calculated at the INDO/MRD-CI level for the different oligomers investigated, as well as the values extrapolated from a linear relationship between the transition energies and the inverse number of thiophene rings and (ii) the fluorescence quantum yields and the nonradiative decay rates measured in solution. In the unsubstituted thiophene (Th3), Rossi et al. have pointed out the existence of a thermally activated decay route of \( S_1 \), which they express the total nonradiative decay rate constant as

\[
k_{NR} = k_1 + k_2(T)
\]

where \( k_1 \) is a nonactivated decay rate (including nonactivated ISC processes, internal conversion, singlet fission, ...) and \( k_2 \) is the intersystem crossing temperature-dependent decay rate. On the basis of the temperature dependence of \( k_{NR} \), they estimate the pre-exponential factor, \( A_2 \), as equal to \( 1.9 \times 10^{10} \text{ s}^{-1} \), and an activation energy of \( \sim 0.05 \text{ eV} \), in excellent agreement with the interpolated \( S_1 - T_4 \) INDO/MRD-CI energy difference for Th3. Assuming that \( k_1 \) remains constant when elongating the chain, we have calculated the \( k_{NR} \) rate constants on the basis of the experimental \( k_1 \) and \( A_2 \) values and of the calculated activation energies, \( E_{ISC} \), \( k_{NR} = k_1 + A_2 \exp \left( \frac{-\Delta E_{ISC}}{kT} \right) \) (2)

Table 5. INDO/MRD-CI \( S_1 - T_4 \) Energy Difference (in eV) Calculated for the Unsubstituted, Monosubstituted (Th3-X), and Bis-Substituted (X-Th3-X) Oligothiophenes and (between Parentheses) Extrapolated from a Linear Relationship between the Excitation Energies and the Inverse Number of Thiophene Rings in the Unsubstituted Oligomers and a Fluorescence Yields, \( \phi_F \), as well as the Nonradiative Decay Rates, \( k_{NR} \). Measured in Solution (\( k_{NR}^{exp} \)) and Calculated on the Basis of Simple Assumptions (\( k_{NR}^{th} \), see Text).
Geometry relaxation phenomena occur in the lowest singlet and triplet states of oligothiophenes that lead to the appearance of a quinoid character, especially pronounced for the T<sub>1</sub> state. Referring to the soliton—antisoliton picture used to describe excitations in polyacetylene, the electron-hole pair appears to be strongly bonded in both states, with a stronger localization of the triplet state. These results are in line with those obtained recently in PPV oligomers, which indicates that this picture holds true for the nondegenerate ground-state polymers.

We have analyzed in depth the influence on the lowest singlet and triplet excitations of grafting electroactive groups on the conjugated segment of a model thiophene oligomer: terthiophene. Different effects have been calculated on the singlet and triplet states, which can be rationalized on the basis of the different natures of these states, as highlighted by the study of the unsubstituted molecules.

Finally, the various mechanisms giving rise to nonradiative decay of the singlet excitons have been discussed. In oligothiophene solutions, intersystem crossing between the S<sub>1</sub> singlet state and a higher-lying T<sub>4</sub> triplet state (followed by fast nonradiative decay to T<sub>1</sub>) is expected to be very efficient in view of the low activation energy required. (Note that the importance of the ISC effect as a nonradiative decay route of the singlet excitations has been demonstrated by recent experimental investigations on polythiophene). On the basis of simple assumptions, we have presented a model that allows one to rationalize the huge increase in fluorescence yield observed experimentally when increasing the chain length and by substitution with acceptor groups. It is worth stressing that, in films, interchain effects are expected to be very important and to lead to a systematic decrease of the light emission yields, as has been observed recently in a cyano-substituted PPV.

We believe that the information provided by this work should be of prime importance in understanding the photophysical properties of oligothiophenes and derivatives and in defining new routes toward the optimization of the light emission efficiency in LED devices where conjugated compounds are used as active layer.

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**Supporting Information Available:** Tables providing the MNDO-optimized carbon–carbon and carbon–sulfur bond lengths in the S<sub>0</sub>, S<sub>1</sub>, and T<sub>1</sub> states of the dimer, trimer, and tetramer of thiophene and substituted terthiophenes (4 pages). Ordering information is given on any current masthead page.

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