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Pulsed Power Corona Discharges for Air Pollution Control

Erwin H. W. M. Smulders, Bert E. J. M. van Heesch, and Sander S. V. B. van Paasen

Abstract—Successful introduction of pulsed corona for industrial purposes very much depends on the reliability of high-voltage and pulsed power technology and on the efficiency of energy transfer. In addition, it is of the utmost importance that adequate electromagnetic compatibility (EMC) is achieved between the high-voltage pulse source and the surrounding equipment.

Pulsed corona is generated in a pilot unit that produces narrow 50 MW pulses at 1000 pps (net average corona power 1.5 kW). The pilot unit can run continuously for use in industrial applications such as cleaning of gases (100 m³/h) containing NO or volatile organic compounds (VOC's) or fluids (e.g., waste water). Simultaneous removal of NO and ethylene to obtain clean CO₂ from the exhaust of a combustion engine was tested at an industrial site.

Various chemical processes, such as removal of toluene or styrene from an airflow are tested in the laboratory. We developed a model to analyze the conversion of these pollutants. To examine the discharges in the reactor we use current, voltage, and E-field sensors as well as a fast charge-coupled device (CCD) camera. Detailed energy input measurements are compared with CCD movies to investigate the efficiency of different streamer phases.

EMC techniques incorporated in the pilot unit are based on the successful concept of constructing a low transfer impedance between common mode currents induced by pulsed power and differential mode voltages in signal lines and external main lines.

Index Terms—Air pollution control, apparent activation energy, butane, charge-coupled device camera, E-field sensors, electromagnetic compatibility, ethylene, high-voltage, industrial applications, NO, pentane, pilot unit, pulsed corona, pulsed power, pulse source, reliability, streamers, styrene, toluene, transfer impedance, trichloroethane, volatile organic compounds, waste water.

I. INTRODUCTION

HIGH power pulsed corona is a promising type of high-voltage discharge that can be well controlled and has many interesting environmental applications. Energy is deposited in a gas or in a liquid in a highly concentrated form by pulsed corona discharges [1]–[3]. Operating over a wide pressure and temperature range [4], pulsed corona generates electrons, free radicals, excited molecules and, ultraviolet (UV) radiation. Direct bond cleavage or interactions through radicals can breakdown various hazardous organic pollutants. Nonhazardous fragments or compounds that can be further treated using conventional techniques are the result. (e.g., H₂O, CO₂, NO₂, and HCl).

Pulsed corona in water leads to the formation of high local electric fields, electrons, OH* and H* radicals, and the repetitive formation of shock waves. In gases, the energized electrons produce radicals through dissociative electron attachment or through electron impact dissociation [5].

Advantages of pulsed corona treatment are: simultaneous removal of several pollutants, high destruction efficiency, no demands on temperature and pressure, insensitive to contamination, no damage from high loads, widely applicable, simply installed, compact, little service, small scale, and no additives.

Specific parameters of the high-voltage pulse are dictated to some extend by the reactor processes. In addition, the parameters are governed by the demand of a match between the pulse source and the transient pulsed corona load. Pulsed corona processing needs pulsed power that is generated in an efficient and reliable way. We investigate whether spark gaps as a reliable long lifetime high-voltage power switch are a good candidate for use in continuous high power industrial applications.

With respect to electromagnetic compatibility (EMC), methods have been developed to effectively suppress interference, in a systematic and reproducible manner, by appropriate guidelines and rules for the layout [6], [7]. Sensitive apparatus for chemical analysis should operate reliably close to the pulse source. Computers and networks in close vicinity and in the surrounding building should remain undisturbed.

II. PULSE SOURCE AND CORONA REACTOR

The pulsed corona unit produces 100 kV pulses (10 ns risetime, 200 ns wide) at a maximum rate of 1000 pulses per second. An overall efficiency of 40–70% was obtained for the energy transfer from mains AC power to corona energy. The mean time between failures for the complete apparatus was approximately 100 h; it was limited by the problem of partial discharge growth toward breakdown in the cable of the transmission line transformer. After improvement of this part, failureless operation now passed the 350 h and is planned to reach 2500 h as testing continues.

Fig. 1 gives an overview of the layout and main components of the pulse source. The corona reactor is placed on top of the main EMC casing. For the production of the high-voltage pulse we use three stages: two resonant circuits coupled by a pulse transformer, followed by a transmission line transformer (TLT) as a third stage [8]. Each stage compresses the pulse duration and raises the pulse amplitude.

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Fig. 1. Schematic overview of the Eindhoven University of Technology (EUT) pilot with high-voltage pulse source. The layout is shown of main components, voltage, and current measuring systems and EMC cabinets. Component values are: \( C_0 = 36 \) mF, \( L_1 = 27 \) \( \mu \)H, \( C_1 = 12 \mu F \), \( C_2 = 20 \mu F \), \( L_2 = 26 \mu H \), \( C_{HV} = 6 \) nF, HV transformer ratio is 10:600, and it has two symmetrical coils on a C-core with a 0.05 mm core lamination and a core gap of 2 \( \times \) 0.2 mm. The operation, pulse formation, control, and EMC details as well as the dimensions of the reactor are described in the text.

The first resonant circuit \((C_0, L_1, C_1)\) at, respectively, 26 mF, 27 \( \mu \)H, 12 \( \mu \)F) and part of the second resonant circuit \((C_3, L_2 = 28 \mu H)\) are at the low-voltage side (up to 750 V) of the pulse transformer. The pulse transformer (10:600) and the high-voltage capacitor \( C_{HV} \) (6 nF) to which it is connected complete the second circuit. This capacitor is pulse charged by the resonant circuits in 40 \( \mu \)s to a level of 25–35 kV.

The low inductance spark gap (50 nH, coaxial with \( C_{HV} \)) discharges \( C_{HV} \) into the TLT. Diode \( D_1 \) (with snubber \( Z_1, Z_2 \)) holds the voltage at \( C_{HV} \). Although the spark gap is free running, an automatic triggering, derived via 2 M\( \Omega \) from the voltage reversal at the transformer secondary is added to avoid misfiring. The transformer voltage reversal is damped via diode \( D_2 \) and the 28 k\( \Omega \) resistor.

After each high-voltage pulse, part of the energy returns to the low voltage resonant circuits. This surplus of energy depends on the actual breakdown voltage of the spark gap. Depending on the value of this energy, the \( V_{C1} \) control system first dumps part of the previously stored energy of buffer capacitor \( C_2 \) (20 \( \mu \)F). Next the surplus energy is shared by \( C_2 \) and \( C_3 \) for partial reuse in the next pulse.

The reliability of the spark gap is excellent, after \( 10^9 \) pulses (total transferred charge 200 kC), only minor electrode wear is visible. The gap has to be flushed continuously with air (30 Nm\(^3\)/h typically). The spark gap transfers the energy of capacitor \( C_{HV} \) to the TLT: four coaxial cables, each 50 \( \Omega \) and 20 m long (RG214 later replaced by RG218 of 12 m length).

We found that four cables is an excellent number with respect to amount of cable, impedance match, and voltage gain. At no load the TLT multiplication ratio is approximately 5.5, resulting in a peak output voltage of 160 kV (10 ns rise time). Loaded by corona the voltage is 100 kV.

Measured voltage and current and calculated power and energy are shown in Figs. 2 and 3. The corona power is the

Fig. 2. Typical corona voltage and current pulse. The output high-voltage pulse is measured with a differentiating sensor in the feedthrough; the current pulse is measured with one-turn Rogowski coil at the base of the reactor. Passive integrators for both signals are housed in the EMC cabinet for control electronics.

Fig. 3. Power \((VI)\) and energy \((\int VI \, dt)\) of a corona discharge calculated from the measured voltage and current pulses given in Fig. 2.
calculated product \((VI)\) of the measured voltage and current. The energy is the time integral of this product.

The gas to be cleaned flows through the 3.5 m long stainless steel tube reactor, inner diameter 250 mm. The inside wall is smooth or can be fitted with needles. The center electrode is 3 m long, either a wire, diameter 0.25–1.0 mm, or a steel M8 stud. A drawing of the base of the reactor tube and the lower part of the tube (extending unaltered to the end flange) is given in Fig. 1.

III. MEASURING SYSTEMS

The base of the reactor vessel is a one-turn Rogowski coil (a toroid of rectangular minor cross section) to measure the external (corona) current \(I_E\). The central wire is fed into the reactor via a HV feedthrough, which also contains a capacitive sensor to form a differentiating-integrating (DI) measuring system for the external (corona) voltage \(V_E\) \([9], [10]\). In a DI system it is the strong differentiated signal that is transported via the coaxial cable to the passive RC input section of an integrator at the wall of an EMC cabinet. The integrator restores the original waveform but it also acts as an effective EMC filter. Large common mode currents are allowed to flow on the signal cables from sensor to integrator and back via grounding systems. A sufficiently low transfer impedance of these cables leaves the measured signal undisturbed. To avoid interference it is crucial that common mode currents do not enter the EMC cabinet behind the integrator \([7]\). The bandwidth of the DI systems used is 30 kHz–50 MHz. Passive integrators have been applied. Signal evaluation and corrections for the 50 \(\Omega\) load of the integrators were made as follows:

\[
V_X = k_1 V_{o1} + \alpha_1 \int k_2 V_{o2} \, dt \tag{1}
\]

and

\[
I_X = k_2 V_{o1} + \alpha_2 \int k_2 V_{o2} \, dt \tag{2}
\]

where \(V_{o1}\) and \(V_{o2}\) are the integrator outputs, \(k_1\) and \(k_2\) are the calibration factors, and \(\alpha_1\) and \(\alpha_2\) are the known correction factors for the integrator droop.

The E field at the cylinder electrode can be measured by means of a grid sensor, mounted flush with the surface of this electrode \([11]\). Fig. 4 shows the construction of this sensor.

It consists of a measuring electrode (a brass plate with a diameter of 30 mm) behind a grounded stainless steel grid (mesh size 1.4 mm). The distance between the plate and the grid is 2 mm. Charge carriers, which would normally produce a conductive current at the cylinder, are prohibited to enter the grid sensor by a repelling E field between the grounded grid and the plate, which is at a DC bias voltage of 15 V. This bias voltage and the voltage pulse at the central wire have the same polarity. Since the grid is not perfectly shielding, still some 2% of the corona E field is seen by the plate. This transient field produces a displacement current to the plate. The resulting signal is easily distinguished from the DC bias of the sensor. Related techniques to detect the E field are given by other authors \([12], [13]\). To obtain the total displacement current \(I_D\) to the cylinder electrode, the current to the plate of the E-field sensor needs to be multiplied by a factor that accounts for the ratios of surface area and grid penetration. Finally the current must be integrated in the time domain to obtain the displacement charge \(Q_D\), which is proportional to the E field at the cylinder. The bandwidth of this system is 0.5 Hz–50 MHz.

Operating as ion sensor, a forward DC-bias voltage (900 V) on the plate is applied to attract the ions into the grid sensor toward the plate. The resulting conduction current to the plate is a measure for the number of ions arriving at the sensor. The mesh size of the grid is small enough to strongly reduce the corona E field, as is necessary for the ion measurements.

For registrations of the voltage, current, and E-field signals a Nicolet 450 digital oscilloscope, 200 MSA/s, and an HP 54542 A digital oscilloscope, 2 GSA/s were used.

Chemical analysis was done with an HP 5880 gas chromatograph and a flame ionization detector (FID). For propane, butane, pentane, and ethylene a packed column porapak Q was used, for styrene, toluene, and 1,1,1-trichloro-ethane we used a packed SE 54 column. For styrene and toluene temperature programmed runs were made. The NO content of the flue gas of a gas-fueled engine at an industrial site was detected with a chemoluminescence method. All NO measurements in the laboratory were made with a testo 342-2 flue gas analyzer.

In the laboratory, the gas flow to be cleaned was composed by mixing contaminant with an airflow. This mix was fed to the main airflow, produced by pumping environmental air (1.5% \(H_2O\)) through the reactor \([14]\). Main airflow was measured with a testovent 4300 anemometer. Contaminants were fed in through Brooks mass flow controllers. In the case of styrene, toluene, pentane, and 1,1,1 trichloro-ethane, a heated bubbler was used for evaporation. The bubbler was followed by a cool down vessel (20°C) and the dry airflow into the bubbler was set with a mass flow controller; see Fig. 5. The flow from this system was added to the bulk flow of environmental air through the reactor. In and output concentrations to the reactor were determined from the known input flow and from subsequent measurements with the same instrument while the reactor was in the on or off state.
The total corona current, measured externally as $I_X$, can be represented [1] as a conductive component $I_C$ (passing charge carriers) plus a displacement component (changing E field) $I_D$. This component, measured separately with the E-field sensor, equals the current $I_0$ due to the changing external voltage $V_X$ across the vacuum capacitance $C_0$ between the electrodes, plus the current $I_S$ due to the movement of space charge in the gas

$$I_X = I_C + I_D = I_C + I_0 + I_S \quad (3a)$$

and

$$I_0 = C_0 \frac{dV_X}{dt}. \quad (3b)$$

A number of CCD images (5 ns gating) are recorded in successive discharges to display the development of the corona discharge. The direction of view of the CCD camera was axial, parallel to the corona wire. CCD images cover the full cross section of the cylindrical reactor. From each recorded image we copy a strip running from wire to wall and combine these strips into one picture, a CCD movie, as shown in Fig. 6.

IV. DISCHARGE DEVELOPMENT

Both positive and negative corona show similar behavior. First, many streamers develop from the wire toward the cylinder during the initial streamer phase [1]. The strongly inhomogeneous space charge at the streamer head creates a large E field, which causes an enhanced ionization and further growth of the initial streamers toward the cylinder. Each streamer carries a current of 0.1–1 A. The capacitance between the streamer head and the cylinder acts as a limiting impedance for the current through the initial streamer. The conductive streamer phase starts after the arrival of the initial streamers at the cylinder. The development of the initial streamer into a complete channel between the wire and the cylinder results in a disappearance of the capacitance between the streamer head and the cylinder. This allows a much larger current (1–10 A) to flow through each streamer; many streamers in parallel carry the large current of the pulsed corona discharge.

The curves for the measured charges $Q_X$ (total) and $Q_D$ (displacement only) are equal until the arrival of the initial streamers at the cylinder. After arrival a conductive current at the cylinder causes a separation of the curves $Q_X$ and $Q_D$. With a negative polarity at the wire, corona also develops at the needles covering the cylinder electrode. The resulting conductive current causes an early separation between the curves $Q_X$ and $Q_D$.

Based on current and voltage measurements a calculation was made of the energy input during the initial and the conductive streamer phase of a positive polarity discharge. The results are given in Fig. 7. It turns out that after the completion of a pulse, about one third of the corona energy is a result of the initial streamer phase and 2/3 results from the conductive streamer phase. If the indications are correct that the amount of corona energy is the important parameter for chemistry then the conductive streamer phase can add considerably to the conversion process.

V. REACTOR MODEL

An exponential behavior for the VOC's removal by the pulsed corona reactor was found experimentally and was given as [15]

$$X = 1 - \exp(-E/\Delta E), \quad (4)$$

The dimensionless $X$ is the amount removed divided by the amount on input, where both amounts can be consistently expressed either in terms of $g/h$ or as concentration, e.g., mole fraction or parts per million. The parameter $E$ is the applied corona energy per unit volume and it will be shown now that $\Delta E$ is an apparent activation energy per unit volume for a mixture of a specific VOC plus air. Within the concentration ranges that we tested this apparent activation energy for the mixture is independent of the VOC's concentration. An analysis based on gas dynamics and reaction kinetics is used to explain the exponential behavior and apparent activation energy. The reactions are activated by use of corona energy in an oxidating environment. Since the concentration of reactants formed in the environmental air by the corona discharge is high compared to the VOC's concentration only first order reactions are considered. As the mechanism for the reaction is not known, an alternative expression for the rate constant $k$ [number of transitions per second] of the reaction will be needed.

With the experimental data we can find Peclet numbers of $10^4$–$10^5$ which means we are dealing with a plug flow reactor. In our tubular plug flow reactor we have Reynolds numbers of $2000$–$10^4$ which means that the flow is turbulent. With this knowledge we can make the following assumptions:

• only energy from the corona discharge is used for decomposition;
• the apparent activation energy is independent of the VOC's concentration (ad hoc, verified by experiment);
• the reactor is an ideal plug flow reactor;
• only first order reactions are considered;
• an alternative expression for the rate constant $k$ is postulated.

A negative free energy change of a reaction means that after the reactions the products are at a lower free energy level. However, if bonds are broken, the reactants must go up an energy hill first, before they will go downhill. The activation of VOC’s plus air can be reached by using corona energy [14]. We are concerned about the apparent activation
energy $\Delta E [J/l]$ needed for the decomposition of VOC’s in the mixture. For the rate constant $k$ [number of transitions per second] of the reaction we postulate a simple alternative expression [14]

$$k = \gamma P/\Delta E$$

(5)

where $P$ is the corona power per $m^3$ and $\gamma$ the used fraction of this power. The apparent activation energy per $m^3$, $\Delta E$, varies with the concentration of the mixture air plus VOC’s

$$\Delta E = \Delta E_0 C_{\text{mix}}/C_{\text{mix0}}$$

(6)

Here $C_{\text{mix}}$ is number of moles of mixture per unit volume [mol/m³] and $C_{\text{mix0}}$ and $\Delta E_0$ are the values at standard temperature and pressure (STP) conditions.

The reaction rate $R$ (number of moles converted per second per $m^3$) of a first order reaction is

$$R = kC.$$  

(7)

The VOC’s concentration is $C$. The reaction rate is substituted in the continuity equation

$$\partial C/\partial t + u \partial C/\partial z = -R$$

(8)

for the flow in the reactor. The flow velocity is $u$ [m/s] in the direction along the axial $z$ coordinate of the cylinder.

To find the stationary solution of (8) we integrate along the length of the reactor and arrive at the desired exponential relation given above in (4).

The exponential relationship is very useful for fitting measured data in many cases. Some removal processes however do not fit the model.

VI. CONVERSION OF VOC’S AND NO

Simultaneous removal of NO and ethylene to obtain clean $CO_2$ from the exhaust of a power generating combustion
engine is an attractive option for CO₂ recycling applied to crops growing. Initial tests were performed at an industrial combined heat and power (CHP) site. An overview of the setup is given in Fig. 8.

During generation of 110 kW of electrical power, the flue gas was treated by the pulsed corona pilot unit. The unit was running at a corona power level of up to 600 W. NO removal (oxidation into NO₂) was determined in terms of g/h (NO and total NOₓ measurements were taken before the scrubber and before the reactor, nitric acid was not measured). Removal of other flue gas fractions such as ethylene was found to be minimal, most probably due to the large NO load of 1.6 kg/h. NO has the lowest apparent activation energy and is therefore processed more easily. Fig. 9 shows the resulting NO-removal rate as a function of corona power. Since a relatively small fraction of the high NO load was removed, we see a high removal efficiency (95 g/kWh). It is clear that we are in the initial part of the exponential function X(E) given in (4).

In the laboratory it was clearly shown that styrene, toluene, ethylene, and NO-ethylene mixtures can be decomposed to a high degree with our pulsed corona reactor. The VOC’s propane, butane, pentane, and 1,1,1 trichloro-ethane are not completely decomposed. Table I gives an overview of chemical measurements. We now compare the calculated conversion using (4) and the measured conversion. It is found that the model fits in the cases of: styrene, toluene, 1,1,1 trichloro-ethane, ethylene, NO, and pentane. The conversion for propane and butane calculated with (4) does not fit the experimental data. Part of this conversion may require additional intermediate products which is not in line with the assumption of a first order reaction. Fitted values for ΔE₀ are summarized in Table I. The fit for toluene is illustrated in Fig. 10 by a graph of measured values X_EXP (amount removed divided by amount on input both taken from toluene measurements) versus calculated values for the conversion, X_MODEL (amount removed divided by amount on input as calculated from (4) by using the measured corona energy density and the fitted value for the apparent activation energy). Here the full range of residence times, concentrations, and corona energies is covered by (4) and a single value for ΔE₀.

Apparent activation energies found in this work are compared with ΔE₀ values derived from results of other researchers [16]–[22]. A summary can be found in Table II.
TABLE I
OVERVIEW OF CHEMICAL MEASUREMENTS AND DECOMPOSITION OF VOC’S IN AIR. THE VALUES FOR THE APPARENT ACTIVATION ENERGY $\Delta E_0$ ARE FITTED ACCORDING TO (4) AND (6)

<table>
<thead>
<tr>
<th>Substance</th>
<th>M [kg/kmol]</th>
<th>Concentration [ppm]</th>
<th>Mass flow [g/h]</th>
<th>Corona Power [kW]</th>
<th>Corona energy kWh per kg load to reach 63% removal</th>
<th>$\Delta E_0$ [J/liter]</th>
</tr>
</thead>
<tbody>
<tr>
<td>NO</td>
<td>30</td>
<td>213</td>
<td>6.5-23</td>
<td>0.15 - 1.10</td>
<td>15</td>
<td>16.5</td>
</tr>
<tr>
<td>Toluene</td>
<td>92</td>
<td>125 - 450</td>
<td>18-40</td>
<td>0.13 - 1.16</td>
<td>24</td>
<td>98.8</td>
</tr>
<tr>
<td>Styrene</td>
<td>104</td>
<td>30 - 190</td>
<td>15-87</td>
<td>0.25 - 2.54</td>
<td>7</td>
<td>11.3</td>
</tr>
<tr>
<td>1,1,1 TCA</td>
<td>133</td>
<td>80 - 1000</td>
<td>4-25</td>
<td>0.13 - 1.13</td>
<td>180</td>
<td>135</td>
</tr>
<tr>
<td>Pentane</td>
<td>72</td>
<td>80 - 1000</td>
<td>6-73</td>
<td>0.14 - 1.04</td>
<td>88</td>
<td>185</td>
</tr>
<tr>
<td>Ethylene</td>
<td>28</td>
<td>150 - 2500</td>
<td>17-114</td>
<td>0.16 - 1.23</td>
<td>12</td>
<td>41.2</td>
</tr>
<tr>
<td>Propane</td>
<td>44</td>
<td>10 - 500</td>
<td>1.7-36</td>
<td>0.16 - 1.15</td>
<td>180</td>
<td></td>
</tr>
<tr>
<td>Butane</td>
<td>58</td>
<td>50 - 2000</td>
<td>9-111</td>
<td>0.14 - 1.19</td>
<td>75</td>
<td></td>
</tr>
</tbody>
</table>

Fig. 10. Measured conversion of toluene versus calculated conversion using the model represented by (4). The full range of flow rates, input concentrations, and applied corona power is covered.

TABLE II

<table>
<thead>
<tr>
<th></th>
<th></th>
<th></th>
<th></th>
<th></th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>NO</td>
<td>16.5</td>
<td>50</td>
<td>10</td>
<td>10-21</td>
<td></td>
</tr>
<tr>
<td>Toluene</td>
<td>98.8</td>
<td>120</td>
<td>91-253</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Styrene</td>
<td>11.3</td>
<td></td>
<td>185</td>
<td></td>
<td></td>
</tr>
<tr>
<td>CCl₄</td>
<td>219</td>
<td>555</td>
<td>9</td>
<td></td>
<td></td>
</tr>
<tr>
<td>TCA (1,1,1)</td>
<td>135</td>
<td></td>
<td>80</td>
<td></td>
<td></td>
</tr>
<tr>
<td>TCE</td>
<td>16</td>
<td>9</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Benzene</td>
<td>500</td>
<td>39</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Pentane</td>
<td>185</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Ethylene</td>
<td>41.2</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Propane</td>
<td>do not fit</td>
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<tr>
<td>Butane</td>
<td>model</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Remarks:
- LANL in N₂, 20 % O₂, 1% H₂O
- LLNL normal: in N₂
- Italic: in dry air
- Roush-Hutherson values depending on pulse parameters

VII. EFFICIENCY

The transport of energy from mains AC power to chemical processes can be divided in three separate conversion steps [23].

1) Energy from mains power into high-voltage pulses.
2) High-voltage energy into corona discharge energy in the gas.
3) Corona energy finally used for chemical processes.

The NO removal process was chosen for investigation of the efficiencies of the three separate steps of energy transport. At the same time we investigated the effect of polarity on these efficiencies. An overview of the results is given in Fig. 11.

The conversion in steps 1) and 2) was determined from the electrical measurements on mains power and on corona power.

The application of corona energy for chemical processes, step 3), was expressed as amount of substance processes per kWh of the corona energy from step 2). Step 3) is found to be independent of polarity, i.e., chemistry is polarity independent here.

Step 2) energy transfer, however, is very much affected by polarity choice. A positive wire gives a better energy transfer when polarity effects are compared at the same external voltage. As we know from previous work [1], a negative wire
needs a higher voltage to produce intense coronas. This effect is not accounted for in this study, we applied the same voltage level for both polarities.

VIII. CONCLUSIONS

Simultaneous recordings with CCD camera and electrical diagnostics allow a clear analysis of the high intensity corona development. Initial streamer phase and transition to a conductive streamer phase are detected and an excellent correlation is found between all data from electrical records and CCD movies. The energy input during the various streamer developments was followed.

Chemical conversion processes initiated by the pulsed corona were tested in the lab but also at industrial locations. A model that only needs two parameters, apparent activation energy and corona energy, can in many cases explain the conversion processes. Initial streamer phase and transition to a conductive streamer phase are detected and an excellent correlation is found between all data from electrical records and CCD movies. The energy input during the various streamer developments was followed.

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REFERENCES


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