Organized versus self-organized criticality in the abelian sandpile model
Fey - den Boer, A.C.; Redig, F.H.J.

Published: 01/01/2004

Citation for published version (APA):
SPOR-Report 2004-13

Organized versus self-organized criticality in the abelian sandpile model

A.C. Fey-den Boer
F.H.J. Redig
SPOR-Report
Reports in Statistics, Probability and Operations Research

Eindhoven University of Technology
Department of Mathematics and Computing Science
Probability theory, Statistics and Operations research
P.O. Box 513
5600 MB Eindhoven - The Netherlands

Secretariat: Main Building 9.10
Telephone: + 31 40 247 3130
E-mail: wscosor@win.tue.nl
Internet: http://www.win.tue.nl/math/bs/cosor.html

ISSN 1567-5211
Organized versus self-organized criticality in the abelian sandpile model

A. Fey-den Boer*
F. Redig†

August 24, 2004

Abstract: We define stabilizability of an infinite volume height configuration and of a probability measure on height configurations. We show that for high enough densities, a probability measure cannot be stabilized. We also show that in some sense the thermodynamic limit of the uniform measures on the recurrent configurations of the abelian sandpile model is a maximal element of the set of stabilizable measures. In that sense the self-organized critical behavior of the ASM can be understood in terms of an ordinary transition between stabilizable and non-stabilizable.

Key-words: Self-organized criticality, abelian sandpile model, activated random walkers, stabilizability.
AMS classification: 60K35 (primary), 60G60 (secondary)

1 Introduction

Self-organized criticality (SOC) is a concept introduced in [3] to model power-law behavior of avalanche sizes in various natural phenomena such as sand and rice piles, forest fires, etc. The conceptual point of view in [3] is that this kind of "criticality" is not tuned by parameters such as temperature, magnetic field, as is the case in critical systems of equilibrium statistical mechanics. This point of view has been questioned by several people, see e.g. [5, 1], where it is argued that the choice of the models exhibiting SOC involves an implicit tuning of parameters, and hence SOC is an (interesting) example of ordinary criticality. In the case of the abelian sandpile model, e.g. one can say that the choice of the toppling matrix (which governs the dynamics) having mass equal to zero is a fine tuning. Indeed in the massive or

*Faculteit exacte wetenschappen, Vrije Universiteit Amsterdam, De Boelelaan 1105, 1081 HV Amsterdam and EURANDOM, Technische Universiteit Eindhoven, Postbus 513, 5600 MB Eindhoven, The Netherlands, afey@eurandom.nl
†Faculteit Wiskunde en Informatica, Technische Universiteit Eindhoven, Postbus 513, 5600 MB Eindhoven, The Netherlands, f.h.j.redig@tue.nl
dissipative case (where in the bulk grains are lost upon toppling) the avalanche sizes exhibit exponential decay, so in that case there is no criticality.

Similarly, in [11], the authors investigate the relation between the critical density of some parametric model of random walkers with that of the abelian sandpile model, and prove in \( d = 1 \) that ASM density corresponds exactly to the transition point in the random walkers model. They further conjecture that this is also true in \( d \geq 2 \). In this paper we want to continue the relation between an ordinary critical phenomenon and the SOC-state of the abelian sandpile model. This is done through the notion of “stabilizability”. A height configuration is called stabilizable if upon stabilizing it in larger and larger volumes, the number of topplings at a fixed site does not diverge. This implies that we can “redistribute” the mass in infinite volume such that after the redistribution, all sites have a height between 1 and 2\( d \). Similarly a probability measure \( \nu \) on height configurations is called stabilizable if it concentrates on the set of stabilizable configurations. The conjecture in [11], inspired by [5] is that there exists \( \rho_c > 0 \) such that (modulo some restrictions on the measure \( \nu \)) if the \( \nu \) expected height \( \rho = \nu(\eta(0)) < \rho_c \), then \( \nu \) is stabilizable, if \( \rho > 2d \) it is not stabilizable, and for any \( \rho \in (\rho_c, 2d) \) there exist measures \( \nu \) with expected height \( \rho \) which are not stabilizable. Moreover, \( \rho_c \) is exactly the expected height in the stationary state of the abelian sandpile model, in the thermodynamic limit.

The aim of this paper is to prove the last two items of this conjecture, and to give some more insight in the regime \( \rho \in (\rho_c, 2d) \). Our paper is organized as follows: in section 2 we define the notion of stabilizability. In section 3 we precisely state the main conjecture of [11] and prove item 3 of it. In section 4 we prove item 2 of the conjecture, and show that in some sense the infinite volume limit of the stationary measure of the abelian sandpile model is “maximal stabilizable”. Finally in section 5 we introduce the concept of meta-stabilizability, and give a class of examples of probability measures on height configurations having this property.

2 Basic definitions

A height configuration is defined as a map \( \eta : \mathbb{Z}^d \to \{1, 2, \ldots\} \), where by convention the minimal height is chosen to be 1. The set of all height configurations is denoted by \( \mathcal{H} \). The set of probability measures on the Borel sigma-field of \( \mathcal{H} \) is denoted \( \mathcal{P}(\mathcal{H}) \). A configuration is called stable if for all \( x \in \mathbb{Z}^d \), \( \eta(x) \leq 2d \). The set of stable configurations is denoted by \( \Omega \). Similarly, for \( V \subset \mathbb{Z}^d \), \( \Omega_V \) denotes the set of stable configurations \( \eta : V \to \{1, \ldots, 2d\} \).

For \( V \subset \mathbb{Z}^d \) the abelian sandpile toppling matrix is defined as

\[
(\Delta_V)_{x,y} = 2d\delta_{z,y} - 1_{z,y \in V, |z - y| = 1}\tag{2.1}
\]

Its inverse is denoted by

\[
G_V(x, y) = (\Delta_V^{-1})_{xy}\tag{2.2}
\]

The probabilistic interpretation of \( G_V \) is:

\[
G_V(x, y) = \frac{1}{2d} \mathbb{E}_x^V \text{ (number of visits at site } y)\tag{2.3}
\]

2
where $E^V_x$ denotes expectation in the random walk started at $x$, and killed upon exiting $V$.

**Definition 2.4.** A height configuration $\eta \in \mathcal{H}$ is called called stabilizable if for any sequence of volumes $V_n \uparrow \mathbb{Z}^d$, there exists $m_{V_n} \in \mathbb{N}^V$ such that

$$\eta_{V_n} - \Delta_{V_n} m_{V_n} = \xi_{V_n} \in \Omega_{V_n}$$

(2.5)

and for any $x \in \mathbb{Z}^d$ $m_{V_n}(x) \to m(x)$ as $n \to \infty$.

The set of all stabilizable configurations is denoted by $\mathcal{S}$. It follows immediately from the definition that for $\eta \in \mathcal{S}$, and $m = \lim_{n \to \infty} m_{V_n}$,

$$\eta - \Delta m = \xi \in \Omega$$

(2.6)

where $\Delta_{x,y} = 2d \delta_{xy} - \mathbb{1}_{|x-y|=1}$ is the infinite volume toppling matrix.

**Definition 2.7.** A probability measure $\nu$ on height configurations is called stabilizable if $\nu(\mathcal{S}) = 1$.

It is clear that the set of stabilizable configurations is a translation invariant subset of $\mathcal{H}$. Therefore any stationary and ergodic probability measure $\mu$ on $\mathcal{H}$ satisfies $\mu(\mathcal{S}) \in \{0, 1\}$.

Remark that in finite volume $V \subset \mathbb{Z}^d$, the abelian sandpile model is “well-defined”. This means that for any height configuration $\eta \in \mathbb{N}^V$, the equation

$$\eta_V - \Delta_V m_V = \xi_V$$

(2.8)

with unknowns the couple $(m_V, \xi_V)$ has at least one solution, namely for $x \in V$, $m_V(x)$ equals the number of topplings at $x$ needed to stabilize $\eta$ in $V$ (see e.g. [8]). Notice that the couple $(m_V, \xi_V)$ is not unique, but if $m_V$ is such that $\xi_V$ is stable, then $m_V(x)$ is greater or equal than the number of topplings at $x$ needed to stabilize $\eta_V$. So the vector collecting the number of topplings (needed to stabilize $\eta_V$) is the minimal solution $m_V$ of the equation (2.8). We will always choose this solution in the sequel.

By abelianness, $m_V(x)$ is non-decreasing in $V$. Therefore a configuration $\eta$ in not stabilizable if and only if there exists $x \in \mathbb{Z}^d$ such that $m_V(x) \uparrow \infty$. By abelianness, the $m_V$ are also non-decreasing functions of the configuration $\eta$, w.r.t. the pointwise ordering $\eta \leq \xi$ if $\forall x \in \mathbb{Z}^d$, $\eta(x) \leq \xi(x)$. Therefore we have the following immediate properties of the set of stabilizable configurations.

**Proposition 2.9.** a) $\mathcal{S}$ is a translation invariant measurable set.

b) If $\eta \in \mathcal{S}$ and $\xi \leq \eta$, then $\xi \in \mathcal{S}$.

c) If $\mu$ is a stabilizable probability measure, and $\nu \leq \mu$, then $\nu$ is a stabilizable probability measure.
We then define the following "critical densities":

**Lemma 2.10.** Define

\[
\rho_c^+ = \inf\{\rho \geq 1 : \text{such that } \exists \nu \in \mathcal{P}(\mathcal{H}) \text{ with } \nu(\eta(0)) = \rho, \text{ and } \nu \text{ is not stabilizable}\}
\]

\[
\rho_c^- = \sup\{\rho \geq 1 : \text{such that } \forall \nu \in \mathcal{P}(\mathcal{H}) \text{ with } \nu(\eta(0)) = \rho, \nu \text{ is stabilizable}\}
\]

Then \(\rho_c^+ = \rho_c^-\).

**Proof.** It suffices to see that the set

\[
S = \{\rho \geq 1 : \text{such that } \forall \nu \in \mathcal{P}(\mathcal{H}) \text{ with } \nu(\eta(0)) = \rho, \nu \text{ is stabilizable}\}
\]

is an interval. Suppose that \(\rho \in S\) and \(\rho' < \rho\). Consider a measure \(\nu' \in \mathcal{P}(\mathcal{H})\) with \(\nu'(\eta(0)) = \rho'\). Then there exists a measure \(\nu \in \mathcal{P}(\mathcal{H})\) such that \(\nu(\eta(0)) = \rho\) and \(\nu \geq \nu'\). Since \(\nu\) is stabilizable, by the monotonicity property 2.9 item 3, \(\nu'\) is stabilizable. \(\square\)

We now introduce the "critical state" of the sandpile model, and its thermodynamic limit. Define a configuration allowed in a volume \(V\) if for any subset \(W \subset V\), the inequality

\[
\eta(x) \leq |\{y \in W, |y - x| = 1\}|
\]

is violated for at least one \(x \in W\). The set of allowed configurations in volume \(V\) is denoted by \(\mathcal{R}_V\). It is well-known that the stationary measure of the abelian sandpile model in finite volume \(V\) is uniform on \(\mathcal{R}_V\) (see e.g. [2] or the basic reference [4]). We denote this measure by \(\nu_v\). Recently, it has been proved in [9], [2] that the weak limit \(\mu = \lim_{V \to 2^d} \mu_v\) exists and defines a measure on infinite volume height configurations. Moreover, its support \(\mathcal{R}\) is the set of those configurations such that all restrictions \(V\) have the property \(\eta_V \in \mathcal{R}_V\). We will call this measure \(\mu\) the uniform measure on recurrent configurations (UMRC). We will always use the symbol \(\mu\) for the UMRC. Put \(\rho_c = \mu(\eta(0))\).

## 3 Main conjecture and results

In the rest of the paper we will prove point 2 and 3 of the following conjecture appearing in [11], cf. also [5], and we will also give some additional new results and examples.

**Conjecture:**
Let \(\nu\) be a stationary and ergodic probability measure on \(\mathbb{N}^{2d}\). Put \(\rho = \nu(\eta(0))\).

1. For \(\rho < \rho_c\), \(\nu\) is stabilizable
2. For \(\rho_c < \rho \leq 2d\) there exist \(\nu\) which are not stabilizable
3. For \(\rho > 2d\), \(\nu\) is not stabilizable
The following theorem settles point 3 of the conjecture.

**Theorem 3.1.** Suppose that \( \eta \) has a distribution \( \nu \) such that \( \nu(\eta(0)) = \rho > 2d \). Then \( \eta \) is almost surely not stabilizable.

**Proof.** Let \( G_V \) be the Green function introduced in (2.1), (2.2). Simple random walk killed upon exiting \( V \) will be denoted by \( \{X_n, n \in \mathbb{N}\} \), and corresponding expectation by \( \mathbb{E}_V^\nu \). Finally, let \( \tau_V \) denote the lifetime of this walk. The infinite volume random walk expectation is denoted by \( \mathbb{E}_0 \). Of course \( \mathbb{E}_0 \) and \( \mathbb{E}_V^\nu \) expectation of events before \( \tau_V \) coincide.

Suppose that \( \eta \) drawn from \( \nu \) is stabilizable. Then we have

\[
m_\nu(0) = \sum_{x \in V} G_V(0, x)(\eta(x) - \xi(x))
\]

\[
= (2d)^{-1} \mathbb{E}_0 \left( \sum_{n=0}^{\tau_V} (\eta(X_n) - \xi(X_n)) \right)
\]

and \( m_\nu(0) \uparrow m(0) < \infty \) as \( V \uparrow \mathbb{Z}^d \). Since the random field \( \eta \) is stationary and ergodic, we have

\[
\lim_{V \uparrow \mathbb{Z}^d} \frac{1}{\tau_V} \sum_{n=1}^{\tau_V} \eta(X_n) = \rho
\]

\( \mathbb{P}_0 \times \nu \) almost surely, where \( \mathbb{P}_0 \) denotes the pathspace measure of the simple random walk starting at 0. For \( \xi \) we cannot conclude such a strong statement but we have, by stability

\[
\limsup_{V \uparrow \mathbb{Z}^d} \frac{1}{\tau_V} \sum_{n=1}^{\tau_V} \xi(X_n) \leq 2d
\]

Therefore since \( \rho > 2d + \delta \) for some \( \delta > 0 \),

\[
\liminf_{V \uparrow \mathbb{Z}^d} \frac{1}{\tau_V} \sum_{n=1}^{\tau_V} (\eta(X_n) - \xi(X_n)) \geq \delta
\]

This implies, using Fatou's lemma, and the fact \( \tau_V \to \infty \) that for any \( A > 0 \)

\[
2d \liminf_{V \uparrow \mathbb{Z}^d} m_\nu(0) = \liminf_{V \uparrow \mathbb{Z}^d} \mathbb{E}_0 \left( \sum_{n=0}^{\tau_V} (\eta(X_n) - \xi(X_n)) \right)
\]

\[
\geq \liminf_{V \uparrow \mathbb{Z}^d} \mathbb{E}_0 \left( AI(\tau_V > A) \frac{1}{\tau_V} \left( \sum_{n=0}^{\tau_V} (\eta(X_n) - \xi(X_n)) \right) \right)
\]

\[
\geq \mathbb{E}_0 \left( \liminf_{V \uparrow \mathbb{Z}^d} AI(\tau_V > A) \liminf_{V \uparrow \mathbb{Z}^d} \frac{1}{\tau_V} \left( \sum_{n=0}^{\tau_V} (\eta(X_n) - \xi(X_n)) \right) \right)
\]

\[
\geq A \delta
\]

Since \( A > 0 \) is arbitrary, we arrive at a contradiction.

\( \square \)
4 Adding to the stationary measure

In this section we settle point 2 of the conjecture.

The UMRC $\mu$ is obtained as a limit of finite volume stationary measure $\mu_\nu$. These $\mu_\nu$ are in turn obtained by running the finite volume addition and relaxation process for a long time. Therefore, one can believe that $\mu$ is “on the edge” of stabilizability. More precisely if one could still “add mass” to $\mu$, then $\mu$ would not be stationary.

However, it is not true that $\mu$ is a maximal stabilizable measure in the sense of the FKG ordering of measures. Indeed, one can create the following translation invariant $\nu$: pick a configuration according to $\mu$ and flip all the height ones to height four. This measure is strictly dominating $\mu$ in FKG sense, but it concentrates on stable configurations. In the last section of this paper we will show that such “artificially stable” measures are in some sense “metastable”.

The idea of formalizing the maximality of $\mu$ is that “one cannot add mass to $\mu$”. For $\mu, \nu$ probability measure on $\mathcal{H}$ we denote $A_{\nu}$ the distribution of $TJ + A$ where $TJ$ is distributed according to $\mu$ and $A$ is independent of $TJ$ and distributed according to $\nu$.

Definition 4.1. A probability measure $\mu$ on $\mathcal{H}$ is called maximal stabilizable if for any ergodic translation invariant $\nu$ with $\nu(\eta(0)) > 0$, $\mu \oplus \nu$ is not stabilizable.

In order to state our main result of this section, we need some more conditions on the UMRC $\mu$. For a configuration drawn from $\mu$, define the addition operator $a_x, \nu$ by $a_x, \nu(\eta) = (a_x, \nu \eta_\nu) \eta_\nu$. In words in $V$ the configuration changes as if we added and stabilized in $V$, outside $V$ the configuration does not change. We say that the infinite volume addition operator $a_x = \lim_{V \to \mathbb{Z}^d} a_x, \nu(\eta)$ exists (in the product topology). We now can state our conditions

Definition 4.2. The UMRC is called canonical if

1. The infinite volume addition operators $a_x$ are well-defined w.r.t. the UMRC for any $x \in \mathbb{Z}^d$.

2. The UMRC is stationary w.r.t. the action of $a_x$, i.e., if $\eta$ is distributed according to the UMRC, then so is $a_x \eta$.

In [6] we prove that these conditions are satisfied on $\mathbb{Z}^d$, $d \geq 5$. The restriction $d \geq 5$ is however of a technical nature, and we strongly believe that these conditions are satisfied as soon as $\mu$ exists. If the UMRC $\mu$ is canonical, then one can easily see that finite products of addition operators are well-defined $\mu$ a.s. and leave $\mu$ invariant. See [6] for a complete proof.

Our main result in this section is the following.

Theorem 4.3. If the UMRC is canonical, then it is maximal stabilizable.
Proof. We have to prove that $\mu \oplus \nu$ is not stabilizable for any $\nu$ stationary such that $\nu(\eta(0)) > 0$. A configuration drawn from $\mu \oplus \nu$ is of the form $\eta + \alpha$, where $\eta$ is distributed according to $\mu$ and $\alpha$ independently according to $\nu$.

Suppose $\eta + \alpha$ can be stabilized, then we can write

$$\eta_V + \alpha_V - \Delta_V m^1_V = \xi^1_V$$

for $m^1_V \uparrow m^1_{\infty}$. We define $m^{2,V} \in \mathbb{N}^{\mathbb{Z}^d}$ by

$$\eta + \alpha^0_V - \Delta m^{2,V} = \xi^{2,V}$$

where $\alpha^0_V : \mathbb{Z}^d \to \mathbb{N}$ is defined as $\alpha^0_V(x) = \alpha(x)1_{x \in V}$. In words this means that we add according to $\alpha$ only the finite volume $V$ but we stabilize in infinite volume. The fact that $m^{2,V}$ is finite follows from the fact that the addition operators $\alpha_x$ and finite products of them are well-defined infinite volume on $\mu$ almost every configuration. Since for $W \supset V$

$$\alpha^0_V \leq \alpha^0_V$$

and $m^1_V$ does not diverge, it is clear that $m^{2,V}$ is well-defined, by approximating the equation (4.5) in growing volumes. Moreover, for $\Lambda \subset \mathbb{Z}^d$ fixed, it is also clear that $(m^{2,V})_{\Lambda}$ and $(m^1_V)_{\Lambda}$ will coincide for $V \supset V_0$ big enough. Otherwise, the stabilization of $\eta_V + \alpha_V$ would require additional topplings in $\Lambda$ for infinitely many $V$'s, which clearly contradicts that $m^1_V$ converges (and hence remains bounded). But then we have that for $V$ big enough, $(\xi^1_V)_\Lambda$ and $\xi^{2,V}_\Lambda$ coincide. For any $V$, the distribution of $\xi^{2,V}$ is $\mu$, because $\mu$ is stationary under the infinite volume addition operators. Therefore, we conclude that the limit $\lim_V \xi^1_V = \xi^0_V$ is distributed according to $\mu$. Hence, passing to the limit $V \uparrow \mathbb{Z}^d$ in (4.5) we obtain

$$\eta + \alpha - \Delta m = \xi$$

where $\eta$ and $\xi$ have the same distribution $\mu$, and where $m \in \mathbb{N}^{\mathbb{Z}^d}$. Moreover, $m$ is the non-decreasing limit of $m^{2,V}$ defined in (4.5), i.e.,

$$m(0) \geq m^{2,\Lambda}(0) = \lim_{V \uparrow \mathbb{Z}^d} \sum_x G_V(0,x)(\eta(x) + \alpha^0_\Lambda(x) - \xi^{2,\Lambda}(x))$$

$$\geq \lim_{V \uparrow \mathbb{Z}^d} \mathbb{E}^V_{\eta}(\sum_{n=0}^{\tau_{V,\Lambda}} E\eta - \xi_\Lambda(X_n) + \sum_{n=0}^{\tau_{V,\Lambda}} \alpha(X_n))$$

(4.8)

where $\tau_{V,\Lambda}$ denotes the exit time of $\Lambda$ of the simple random walk killed upon exiting $V$. Since $\tau_{V,\Lambda}$ diverges in the limit $V \uparrow \mathbb{Z}^d$ and then $\Lambda \uparrow \mathbb{Z}^d$, we conclude by ergodicity of $\mu$, and the fact that both $\xi$ and $\eta$ have this distribution $\mu$, that the term containing $\alpha$ in the r.h.s. of (4.8) diverges faster than the the other term. Hence $m(0)$ is infinite, which gives the desired contradiction.

\[ \blacksquare \]

In dimension $d = 1$ the situation is simpler, see [11] for the proof.
Theorem 4.9. A stationary and ergodic measure \( \nu \) on \( \mathbb{N}^Z \) such that \( \nu(\eta(0)) < 2 \) is stabilizable. If on the contrary \( \nu(\eta(0)) > 2 \), then \( \nu \) is not stabilizable.

Remark 4.10. For \( \nu(\eta(0)) = 2 \) one can have both stabilizability and non-stabilizability: e.g. the configuration 31313131... and its shift 13131313... are not stabilizable.

5 Constructive example

The following is a constructive example of an addition which leads to infinitely many topplings at the origin in the limit \( V \uparrow \mathbb{Z}^d \). It settles point 2 of the conjecture, even in the case when the UMRC is not canonical (in particular for \( d \leq 4 \)), and shows that point 1 is not true in the generality of stationary and ergodic probability measures on height configurations.

We give the example for \( d = 2 \); the generalization to \( d > 2 \) is obvious. Let \( \omega, \omega' \) be independent and distributed according to a Bernoulli measure \( \mathbb{P}_p \) on \( \{0, 1\}^Z \), with \( \mathbb{P}_p(\omega(x) = 1) = p \). Consider the following two dimensional random field \( \zeta(x, y) = \omega(x) + \omega(y) \). In \( \zeta \) we have almost surely infinitely many rectangles \( R_1, \ldots R_n, \ldots \) surrounding the origin with corner sites having height two and the other sites having height one. If we add such a configuration \( \zeta \) to any recurrent configuration \( \eta \) drawn from \( \mu \), then we have that the number of topplings at the origin in the finite volume \( V \) is at least the number of rectangles \( R_i \) that are inside \( V \). Indeed, upon addition according to \( \zeta \) on such a rectangle, every site inside the rectangle will topple at least once, see e.g. [10]. Therefore the distribution \( \mu_p \) of \( \eta + \zeta \) where \( \eta \) is drawn from the UMRC \( \mu \), is not stabilizable. Since we can choose \( p \) arbitrary close to zero, any density \( \rho \in (\rho_c, \rho_c + 2) \) can be attained by \( \mu_p \).

To show that we can get below \( \rho_c \) consider a translation invariant probability measure \( \mu' \) concentrating on some subset \( \mathcal{R}' \) of \( \mathcal{R} \). Then with the same reasoning, the distribution \( \mu'_p \) of \( \eta + \zeta \) where \( \eta \) is drawn from the \( \mu' \), is not stabilizable. Consider therefore \( \mu' \) to be a weak limit point of the uniform measures on minimal recurrent configurations, then the distribution \( \mu'_p \) has the expectation \( \mu_p(\eta(0)) \) arbitrary close to \( \mu'(\eta(0)) < \rho_c \). This shows that point 1 of the conjecture cannot hold in that generality.

Combining our results so far with proposition 2.4 from [11], we conclude

Theorem 5.1. Let \( \rho_c^+ = \rho_c^- \) be as in lemma 2.10. Then

\[
\rho_c^+ = \inf \{ \nu(\eta(0)) : \nu \text{ is translation invariant and } \nu(\mathcal{R}) = 1 \} \quad (5.2)
\]

For \( d \geq 2 \), \( \rho_c^+ < \rho_c \).
6 Other notions of stabilizability

6.1 Stabilization in infinite volume

Definition 6.1. A configuration $\eta \in \mathbb{N}^{\mathbb{Z}^d}$ is called weakly stabilizable if there exists $m \in \mathbb{N}^{\mathbb{Z}^d}$ and $\xi \in \Omega = \{1, \ldots, 2d\}^{\mathbb{Z}^d}$ such that

$$\eta - \Delta m = \xi$$

(6.2)

It is clear that if $\eta$ is stabilizable, then it is weakly stabilizable and we can choose $m = \lim_{V \to \mathbb{Z}^d} m_V$. However it is not clear whether there exist unstable configurations which can be stabilized directly in infinite volume but which satisfy $\lim_{V \to \mathbb{Z}^d} m_V(0) = \infty$, i.e., the infinite volume toppling numbers are not obtained as the limit of toppling numbers in larger and larger volumes. In the following proposition we prove that a measure $\nu$ with $\nu(\eta(x)) > 2d$ for all $x \in \mathbb{Z}^d$ cannot be weakly stabilized. This means in words that mass cannot be “swept away” to infinity.

The following example shows that the opposite, importing mass from infinity, is not impossible. Consider $f : \mathbb{Z}^2 \to \mathbb{N}$:

$$f(x, y) = x^2 + y^2$$

then $\Delta f = -4$ and hence, for example,

$$\bar{\sigma} = \bar{2} - \Delta f$$

(6.3)

where $\sigma$ (resp. $\bar{2}$) denotes the configuration with height 6 (resp. 2) at every site. However,

$$\lim_{V \to \mathbb{Z}^2} \Delta_V^V(\sigma - \bar{2}) = \infty$$

so this “infinite volume toppling” cannot be obtained as a limit of finite volume topplings. Notice however that “toppling” according to $f$ is not “legal” in the following sense: we cannot find an order of topplings such, performed in this order, only unstable sites topple and at the end every site has toppled $x^2 + y^2$ times. The point of (6.3) is that the equality

$$\eta = \xi - \Delta m$$

(6.4)

for some $m \in \mathbb{N}^{\mathbb{Z}^d}$ does not imply that the densities of $\eta$ and $\xi$ are equal. However as we will see later, the equality in (6.4) does imply that the density of $\eta$ is larger or equal than that of $\xi$.

In the following proposition we show point 3 of the conjecture for weak stabilizability.

Proposition 6.5. Let $\nu$ be stationary ergodic such that $\nu(\eta(0)) > 2d$. Then $\nu$ is not weakly stabilizable.
Proof. Suppose that there exist \( m \in \mathbb{N}^{2d} \) such that
\[
\eta - \Delta m = \xi
\]
with \( \xi \) stable and \( \eta \) a sample from \( \nu \). Let \( X_n \) be the position of simple random walk starting at the origin at time \( n \). From (6.6) it follows that
\[
m(X_n) - \frac{1}{2d} \sum_{k=1}^{n-1} (\xi(X_k) - \eta(X_k))
\]
is a martingale (w.r.t. the filtration \( \mathcal{F}_n = \sigma(X_r : 0 \leq r \leq n) \), so \( \eta \) and \( \xi \) are fixed here). Therefore taking expectations w.r.t. the random walk
\[
\frac{1}{n} \mathbb{E}_0 m(X_n) = \frac{1}{n} \mathbb{E}_0 \left( \frac{1}{2d} \sum_{k=1}^{n-1} (\xi(X_k) - \eta(X_k)) \right)
\]
Since \( \nu(\eta(0)) = 2d + \delta \), it follows by stability of \( \xi \) that
\[
\frac{1}{n} \left( \sum_{k=1}^{n-1} (\xi(X_k) - \eta(X_k)) \right) \leq 2d - \frac{1}{n} \sum_{k=1}^{n-1} (\eta(X_k))
\]
Therefore, using dominated convergence
\[
0 \leq \limsup_{n \to \infty} \frac{1}{n} \mathbb{E}_0 m(X_n) = \limsup_{n \to \infty} \frac{1}{n} \mathbb{E}_0 \left( \frac{1}{2d} \sum_{k=1}^{n-1} (\xi(X_k) - \eta(X_k)) \right)
\leq 1 - \mathbb{E}_0 \left( \frac{1}{2d} \lim_{n \to \infty} \frac{1}{n} \sum_{k=1}^{n-1} \eta(X_k) \right) < -\delta
\]
which is a contradiction.

\[\square\]

6.2 Activated random walkers system and stabilizability at low density

Dickman proposes in [5] the following mechanism of stabilization. Consider a configuration \( \eta \in \mathbb{N}^{2d} \). To each \( x \in \mathbb{Z}^d \) is associated a Poisson process \( N^x_t \), for \( x \neq y \) these processes are independent. On the event times of \( N^x_t \) a site topples if it is unstable (the random walkers are “activated”), otherwise nothing happens. This means that after time \( t \) the configuration \( \eta \) evolves towards \( \eta_t \) according to a Markov process with generator
\[
Lf(\eta) = \sum_{x \in \mathbb{Z}^d} \mathbb{1}_{\eta(x)>2d} (f(\eta - \Delta x,.) - f(\eta))
\]
One says now that the configuration is stabilizable by this process if for any \( x \in \mathbb{Z}^d \) the value \( \eta_t(x) \) jumps only a finite number of times. One can write the configuration \( \eta_t \) as

\[
\eta_t = \eta_0 - \Delta n^t_\eta
\]

where \( n^t_\eta \) is the vector collecting at each site the number of topplings at \( x \) in \([0, t]\). If \( \eta \) is distributed according to a translation invariant probability measure \( \nu \) on \( \mathbb{N}^{\mathbb{Z}^d} \) then so is \( n^t_\eta \) under the joint measure \( \nu \times \mathbb{P} \) where \( \mathbb{P} \) is the distribution of the Poisson processes \( N^t_\xi, x \in \mathbb{Z}^d \). Moreover \( n^t_\eta(x) \leq N^t_\xi \) by definition and hence

\[
\mathbb{E}(\eta_t(x)) = \mathbb{E}(\eta_0(x))
\]

i.e., this process conserves the density.

**Lemma 6.12.** A configuration is stabilizable by the process with generator \( L \) if and only if it is stabilizable (in the sense of definition 2.7).

**Proof.** Suppose \( \eta \) is stabilizable by the process with generator \( L \). Consider then the generator

\[
LVf(\eta) = \sum_{x \in V} I_{\eta(x) > 2d} (f(\eta - \Delta^V_x, x) - f(\eta))
\]

corresponding to toppling inside \( V \) only, according to the finite volume toppling matrix, on the event times of the Poisson process \( N^t_\xi \). Call \( n^t_V(x, \eta) \) the number of updates of \( x \) in \([0, t]\). It is easy to see that \( n^t_V(x, \eta) \uparrow n^t_\eta(x) \) as \( V \uparrow \mathbb{Z}^d \), where \( n^t_\eta \) is defined above. By assumption, for any \( x \) there exists \( t^V_x(\eta) \) such that for any \( t \geq t^V_x(\eta) \) \( n^t_V(x, \eta) = n^t_V(x, \eta) \), and moreover \( t^V_x(\eta) \uparrow t_\xi(\eta) \) for which \( n^t_\eta(x) = n^t_x(\eta) \) for any \( t > t_\xi(\eta) \). Therefore in definition 2.7 we can identify

\[
m^V(x) = n^t_V(x, \eta)
\]

and

\[
m(x) = n^t_x(\eta)
\]

Suppose that \( \eta \) is stabilizable in the sense of definition 2.7. Then, clearly, in finite volume we have the equality

\[
m^V(x) = n^t_V(x, \eta)
\]

Since \( m^V \uparrow m \) we have

\[
\sup_{V \subset \mathbb{Z}^d} n^t_V(x, \eta) < \infty
\]

and hence

\[
t_\xi(\eta) = \sup_{V \subset \mathbb{Z}^d} t^V_x(\eta)
\]

is finite \( \mathbb{P} \) almost surely. Now pick \( t > t_\xi(\eta) \). Then

\[
n_t(x, \eta) = \lim_{V \uparrow \mathbb{Z}^d} n^t_V(x, \eta) = \lim_{V \uparrow \mathbb{Z}^d} n^t_V(x, \eta)
\]

(6.13)
where in the second step we used that the processes with generator $L_V$ converge to the process with generator $L$ weakly on path space. Indeed for any local function
\[
\lim_{V \uparrow \mathbb{Z}^d} L_V(f) = L(f)
\]
So the convergence of the processes follows from the Trotter-Kurtz theorem. The right hand side of (6.13) does not depend on $t$ anymore. Hence $\eta$ is stabilizable by the process with generator $L$.

In [11] the authors prove that there exists $\rho'_c > 0$ such that if $\nu(\eta(0)) \leq \rho'_c$, then $\nu$ is stabilizable by the process with generator $L$. $\rho'_c$ is the density of “minimal recurrent configurations”. The following theorem is then an immediate consequence.

Theorem 6.14. There exists $\rho'_c > 0$ such that if $\nu$ is any stationary ergodic probability measure on $\mathbb{N}^2$, then $\nu$ is stabilizable.


7 Metastable measures

Suppose that $\nu \geq \mu$, and $\nu(\eta(0)) > \mu(\eta(0))$, i.e., $\nu$ has a strictly higher density than $\mu$, and stochastically dominates $\mu$. In that situation $\nu$ can still concentrate on stable configurations, and hence be stabilizable. One feels however that such a measure is “on the brink of non-stabilizability”. This is formalized in the following definition.

Definition 7.1. A measure $\nu$ is called metastable (MS) if it is stabilizable and if $\nu \oplus \delta_0$ is not stabilizable with non zero probability, i.e., $\nu \oplus \delta_0(S) < 1$.

In words this means that upon stabilizing $\nu \oplus \delta_0$ in volume $V$, with positive probability the number of topplings $m_V(0)$ diverges as $V \uparrow \mathbb{Z}^d$. The simplest example of a metastable measure is the measure concentrating on the maximal stable configuration $\nu = \delta_{2d}$. The following theorem shows that there are other non-trivial metastable measures.

Theorem 7.2. Suppose that $\nu$ is a stationary and ergodic probability measure on $\Omega$, concentrating on the set of recurrent configurations $R$. Define $I_\nu(x) = 1_{\nu(x) = 2d}$ and call $\tilde{\nu}$ the distribution of $I_\nu$. Suppose that $\tilde{\nu}$ dominates a bernoulli measure $\mathbb{P}_p$ with $p$ sufficiently close to one such that the 1’s percolate and the zeros do not percolate. Then $\nu$ is metastable.

Proof. The idea of the proof is the following. Suppose we have a “sea” of height $2d$ and “islands” of other heights, and such that the configuration is recurrent. Suppose the origin belongs to the sea, and we add a grain at the origin. Start the first wave, i.e, topple the origin once and then all other unstable sites except the origin again. The toppled sites form the first wave. The first wave must be a simply connected subset of $\mathbb{Z}^d$ because the configuration is recurrent. It is clear that the “sea” of height
2d is part of the wave, and therefore every site is contained in the wave (because if an island is not contained then the wave would not be simply connected). So in the first wave every site topples exactly once, but this implies that the resulting configuration is exactly the same. Hence we have infinitely many waves.

Let us now formalize this. For a given configuration \( \eta \) (distributed according to \( \nu \)) a volume \( V \subset \mathbb{Z}^d \) is called "a lake with islands" if all the boundary sites of \( V \) have height \( 2d \), and if from the origin there is a path along sites having height \( 2d \) to the boundary. From the fact that the zeros do not percolate, and 1's do percolate it follows that with positive probability, the origin is in infinitely many nested lakes, i.e. \( V_1 \subset V_2, \ldots, V_n, \ldots, \) with for \( i \neq j \), \( \partial V_i \cap \partial V_j = \emptyset \). Consider a configuration from that event, consider a volume \( V \supseteq V_n \), add a grain at the origin and stabilize the configuration in \( V \). In the first wave all sites will topple once because islands not contained in the wave would contain forbidden subconfigurations, which is impossible since the configuration is recurrent. After the first wave, the only sites that change height are on the boundary of \( V_n \). Therefore, the origin is still unstable and a second wave must start. The sites included in this wave will contain the set of sites included in the first wave needed to stabilize \( \eta_{V_{n-1}} \), but this set, with the same argument, is at least \( V_{n-1} \). Continuing like this, one sees that at least \( n \) waves are needed for the stabilization of \( \eta \) inside \( V \). Since with positive probability we find infinitely many lakes containing the origin, the number of topplings is diverging in the limit \( V \uparrow \mathbb{Z}^d \) with positive probability, which is what we wanted to prove.

Remark 7.3. We believe but cannot prove that any translation invariant probability measure \( \nu \in \mathcal{P}(\mathcal{H}) \) with \( \nu(\eta(0)) \in (\rho_c, 2d) \) is either metastable or not stabilizable.

Acknowledgments. We thank Corrie Quant and Ronald Meester for interesting discussions and suggestions.

References


