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A Compact Integrated Polarization Splitter/Converter in InGaAsP–InP

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Abstract—A novel design for an integrated passive polarization splitter/converter combination is presented. The device consists of a Mach–Zehnder interferometer with polarization converters in both arms. The device is analyzed using the transfer matrix method and fabricated in InGaAsP–InP. Measurement results show a splitting ratio of approximately 10 dB and a conversion of >90%. This device can be monolithically integrated with passive and active components.

Index Terms—Indium phosphide, integrated devices, polarization converter (PC), polarization splitter, wafer stepper.

I. INTRODUCTION

The polarization state of light is of ever greater importance in modern telecommunications networks. First of all, a lot of components in the network are highly polarization-dependent; furthermore, polarization-mode dispersion can degrade the transmission in an optical fiber. On the other hand, the polarization can be employed in, e.g., polarization multiplexing, polarization diversity, and polarization-based filtering [1]. In all these cases, polarization splitters and converters are key elements.

Passive polarization splitters and converters that are able to be integrated with both active and passive components are preferred. Passive polarization splitting can be achieved by loading a waveguide with metal [2], by mode-evolution [3], [4], or by modal birefringence [5], [6]. Splitters based on the latter have the advantage that they have low loss and show a high splitting ratio. A drawback is their length (1–3 mm) which is large compared to other components on the chip. Shorter splitters based on photonic crystal waveguides [7] are reported, but these have the disadvantage of higher losses and more complex processing. We present a compact 600-μm-long integrated splitter based on polarization converters (PCs); preliminary results were shown in [8].

II. PRINCIPLE

The device consists of a Mach–Zehnder interferometer (MZI) with PCs in both arms, as is depicted in Fig. 1. Light coupled into the input waveguide of the first multimode interference (MMI) coupler is split into the two branches with equal power and phase. In the upper branch, a PC is placed that rotates the polarization 90°, so after this, the orthogonal polarization propagates through this branch.

In the lower branch, the light in the original polarization propagates over a distance L, before being rotated in a PC. The birefringence in the waveguides causes a phase shift between light in the arms. This phase shift is equal in magnitude but opposite in sign for both polarizations. When both signals are combined in the output MMI, the phase difference causes one polarization to appear in one of the outputs while the opposite polarization goes to the other output. To achieve the desired splitting, the phase difference between the branches needs to be ±(π/2) radians. This is obtained when

\[ L = \pi / (2(\beta_{TE} - \beta_{TM})) \]

where \( \beta_{TE}, \beta_{TM} \) are the propagation constants for the two (transverse electric and transverse magnetic) polarizations.

The PC consists of a ridge waveguide with a straight and a slanted wall [9], [10].

III. ANALYSIS

The polarization splitter circuit is simulated by concatenating the transfer matrices of each of the sections in the device (Fig. 1): the input coupler (MMI1); a PC in the upper arm (PC1); straight waveguides of length L in both branches (WG); a PC in the lower arm (PC2); and the output coupler (MMI2). Simulation results as a function of the conversion of the PCs in the arms, with TE polarized light at the input, are shown in Fig. 2(a). The power of TE polarized light from the outputs is only zero if the conversion in the arms \( c_{PC} = 1 \). For a lower conversion, the nonconverted part is split equally over the two branches. The conversion of the PCs depends critically on the width (±50 nm needed for \( c_{PC} > 90\% \) [9]), so these devices are considered to be the limiting factor in the performance. The influence of the coupling coefficient of the couplers is less important as MMs can be made tolerant to width deviations [11].

With these, Fig. 3 shows the splitting ratio of the splitter for a coupling coefficient of 0.5 for MMI2, defined as

\[
\text{SR(TE} \text{in}) = 10 \log \left( \frac{P_{TEmax1} + P_{TEcont1}}{P_{TEmax2} + P_{TEcont2}} \right),
\]

Fig. 1. Schematic of the MZI polarization splitter/ converter.
For a splitting ratio larger than 13 dB, a conversion $c_{PC}$ above 90% is needed. The total conversion of the splitter circuit $c_{PS}$ at the wanted output port is larger than 95%, because the unconverted part is split equally over the outputs.

**IV. FABRICATION**

The waveguides used in the splitter are 2 $\mu$m wide, and deeply etched into a layer stack having a 300-nm InP top cladding, and a 500-nm $Q(1.25)$ waveguide layer on an InP substrate. This yields a $\Delta \beta = \beta_{TE} - \beta_{TM}$ of 0.03 $\mu$m$^{-1}$, so for this device an offset $L$ between the converters of 52 $\mu$m is needed. The total length of the device, including input and output MMIs, is about 600 $\mu$m. The device is coupled to 1.8-mm-long shallow waveguides (etched 100 nm into the waveguide layer). The PCs consist of an asymmetric waveguide with a straight and a slanted sidewall.

On one chip, separate PCs and the integrated splitter/converter are fabricated (Figs. 3 and 4). The processing of the polarization splitter/converter is similar to the process described earlier [9]. All waveguides are defined by lift-off of Ti on top of silicon nitride. The lithographical definition is made in an ASML PAS5500/250 5 reduction wafer stepper. This allows a very accurate width control, better than 20 nm on an 800-nm line. This optical lithography is advantageous as compared to electron beam lithography, because it has a large writing field, better uniformity, and is suited for mass production.

The etching of the waveguides and the straight side of the PC is done in a CH$_4$–H$_2$ reactive ion etching. The slanted side is etched in a Br$-$methanol solution. This etchant etches both InP and InGaAsP anisotropically with an angle of 54° with respect to the surface, as shown in Fig. 3.

**V. CHARACTERIZATION**

Both the integrated splitter/converter and the separate PCs on the same chip are examined with the setup shown in Fig. 5. The devices are exited using an erbium-doped fiber amplifier as a source and a 2.5-nm-wide bandpass filter, set to a central wavelength of 1555 nm. This signal is chopped and the polarization is fixed using a polarizer. The light is coupled into the chip and the output is coupled through a polarizer to determine the output polarization. It is detected with a photodiode connected to a transimpedance amplifier and a lock-in amplifier. The separate PCs are measured first. The conversion as a function of the width of the device is examined. The results are shown in Fig. 6.

A conversion of 95% can be achieved for a width of 0.75 $\mu$m. The PCs used in the splitter are 0.8 $\mu$m wide, as the simulated maximum conversion of 99% would be for this width. Due to
inaccuracies in the model, the actual maximum conversion occurs at a different width. According to the measurements, the converter used in the splitter will have a conversion of 88%. A higher conversion is expected for narrower waveguides. The full splitter/converter is measured at the same wavelength. The results are stated in Table I.

The conversion $c_{PC}$ of the converters in the branches for TE (for TM) can be calculated by dividing the output power in TM (TE) in both outputs by the total power from both outputs

$$c_{PC}(TE_{in}) = \frac{P_{TM_{out1}} + P_{TM_{out2}}}{P_{TE_{out1}} + P_{TE_{out2}} + P_{TM_{out1}} + P_{TM_{out2}}}.$$  

The conversion equals 87%. The net conversion of the device $c_{PS}$ is 91% for both polarizations. This is less than the expected net conversion from Fig. 2(b). This is probably caused by an imperfect output coupler.

The resulting splitting ratio, defined in (1), is 9.1 dB for $TE_{in}$ and 11.4 dB for $TM_{in}$. The low splitting ratio is caused by a deviation from the actual $L$ to the optimal $L = \pi/(2(\beta_{TE} - \beta_{TM}))$, caused by imperfections in the model to calculate the propagation constants of the modes.

The difference in splitting ratio for TE and TM is caused by the polarization dependence of the output coupler. The input coupler is a 1 $\times$ 2 coupler and its splitting is symmetric, inherent to the design and thus polarization-independent.

The excess losses are $5,0 \pm 0.1$ dB compared to a straight $3\mu$m shallow waveguide. The PCs have a loss of $1,5 \pm 0.3$ dB; both the input and output MMI have an expected loss of 1 dB.

The additional losses are most probably caused by the waveguide roughness which is visible in Fig. 4. The narrow, deeply etched waveguides in the splitter suffer more from the roughness than the shallow connecting waveguides. The roughness is mainly caused by a nonoptimal lift-off process used for the masking of the waveguides. Improved fabrication and adjustments to the design will lead to lower losses and better performance.

VI. CONCLUSION

A novel type of interference-based integrated polarization splitter and converter is presented. The device is simulated using the transfer matrix method. Splitting ratios larger than 95% are expected for conversion ratios of the converters of more than 90%. The PCs are the limiting factor in the fabrication process. The device is fabricated and measurements show a splitting of 9.1 dB for TE, 11.4 dB for TM, and a conversion of 91%. This device has the potential to be a short splitter that can be integrated with active and passive components, easier than a shorter photonic-crystal-based splitter.

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REFERENCES


