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Passive and Active Constant Force-Displacement Characteristics and Optimization of a Long-Stroke Linear Actuator

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Abstract—In applications such as vibration isolation, gravity compensation, pick and place machines, etc., the applicability of commercial low force passive devices is limited and hence would benefit from an improved and preferably variable force level. This paper presents and investigates such a long stroke constant-force versus displacement actuator topology, where analytical and equivalent circuit models are detailed and compared with comprehensive three dimensional (3D) finite element analyses. Furthermore, the optimization of such an actuator by means of aggressive output space mapping is considered. This technique employs a combination of a single analytical equation and a 3D finite element model. Specifically, the optimization is applied to obtain a passive and active force level of 200 N and 300 N, respectively, where the force-displacement response is constant over 60% to 90% of the total stroke.

I. INTRODUCTION

Passive constant-force actuators (PCFA) are commercially available for low force levels [1]. Measurements performed on such an actuator (Fig. 1) show indeed a constant level of the force output versus translator displacement for a certain stroke range. However, the internal topology of the measured actuator is unknown and thus the actuator topologies discussed in this paper (Figs. 2, 3 and 4) are independently derived. In applications such as vibration isolation, gravity compensation, pick and place machines, etc. or in applications where the PCFA is placed in parallel to a linear brushless actuator, the applicability of low force passive devices is limited and hence would benefit from an improved and preferably adjustable force level. The constant force characteristic is independent on the displacement of the PCFA translator and therefore excludes the position sensor, which is especially important for high production volume applications where low cost is important. A further enhancement would be achieved if the PCFA’s force output could be increased or decreased by means of DC excitation, with the resulting actuator configuration being denoted as an active constant-force actuator (ACFA). For example, in pick and place machines, where the load is constantly varying, the changing gravitational force could be counteracted by an ACFA instead of attaining the increased force from the linear brushless actuator [2]. Another example is given by fluid flow control applications that require a variable valve position with respect to the desired output flow [3]. In these applications solenoids are used, however although that these devices have a very simple structure they require, in order to produce force, a constant current excitation and, furthermore, the force-per-volume ratio is limited for applications with long-stroke valve motions [3]. This paper investigates such a long stroke constant force actuator topology with increased and adjustable force, where the corresponding topologies are shown in Figs. 3 and 4.

For the initial evaluation of the PCFA and ACFA concepts, analytical and lumped parameter models are derived in Sections II to IV. However, in order to verify their solution accuracy, the actuator performance is evaluated by means of the 3D finite element analysis (FEA) in Section V. This initial analysis is based on a set of arbitrary geometrical parameters and in order to maximize the static performance, the optimal design of a long stroke ACFA is considered in Section VI. A recently introduced optimization technique, i.e. the aggressive
output space-mapping (AOSM), is implemented for this purpose. Finally, the conclusions are given in Section VII.

II. SIMPLIFIED PASSIVE ANALYTICAL MODEL

A first estimate of the force capability can be derived from well established analytical expressions. This simplified analytical model is derived based on a number of simplifying assumptions: no leakage fluxes or fringing effects are considered, and the magnet and iron relative permeabilities are taken to be equal to 1 and infinity, respectively. For a magnet material having a linear demagnetization characteristic, with a working point that lies on the linear region, the flux density is given by:

$$B_m = B_r + \mu_r \mu_s H_m,$$

where, $B_m$ is the working flux density, $H_m$ is the corresponding magnetic field strength, $B_r$ is the remanent flux density and $\mu_r$ is the relative recoil permeability.

Starting from the general form of Ampere’s law:

$$\oint \mathbf{H} \cdot d\mathbf{l} = \int_s \mathbf{J} \cdot d\mathbf{A} = I_m,$$

and considering the magnetic field strength to be (piece-wise) constant on the integration path, the simplified expression is obtained:

$$H_{in} l_m + H_{in} l_m = 0. \quad (3)$$

Further, it is assumed that the airgap and magnet flux are equal:

$$\Phi_g = \Phi_m, \quad (4)$$

and, therefore, for the surface-mounted magnet rotor topology, which is the only topology considered in this paper, the average magnet flux density, from (1), is:

$$B_m = \frac{B_r}{1 + \frac{l_g S_m}{l_m S_m}} \quad (5)$$

where $l_g$ and $S_g$ are the airgap length and surface area, and $l_m$ and $S_m$ are the magnet thickness and the magnet area, respectively. The airgap flux density can be derived from (4) as:

$$B_g = B_m \frac{S_m}{S_g}. \quad (6)$$

The force capability of the PCFA is then derived from the rate of change of magnetic co-energy with respect to the translator displacement. The magnetic co-energy over the volume of the problem is given by:

$$W' = \int_V \left( \frac{1}{2} B \cdot dH \right) dV,$$

and, given the simplifying assumptions, it can be reduced to:

$$W' = \frac{B^2}{2\mu_0} S l. \quad (7)$$

The force output, for the direction of travel along the z-axis, is then calculated by:

$$F_z = -\frac{\partial W'}{\partial z}_{z=const}, \quad (9)$$

where, by substituting (5) and (6) in (8), the following expression, which is independent of the $z$-axis displacement due to the exclusion of axial leakage in the analytical model, is obtained for the force amplitude:

$$F_z = -\frac{B^2}{2\mu_0} c_m l_m \left[ \frac{1}{1 + \frac{l_g c_m}{l_m c_g}} - \frac{1}{1 + \frac{l_g c_m}{l_m c_g}} \right] \quad (10)$$

In this, $c_g$ and $c_m$ are the width of the airgap and magnet flux path, respectively, and the expressions for calculating the various considered parameters are summarized in Table I. In this, $l_{g1}$ and $c_{g1}$ are, respectively, the length and area of the airgap of the overlapping part, $l_{g2}$ and $c_{g2}$ are the length and airgap of the non-overlapping part and $l_m$ and $c_m$ are the length and area of the permanent magnets. The various dimensions of $l_1$ to $l_3$ and $g$ are shown in Fig. 11 and summarized in Table II. Further, in Table I, $n_p$ represents the number of poles, where the use of a low number of poles in PCFAs is common (typically 2-4) as it particularly offers advantages in terms of complexity and cost [1]. Although that, for an increased force level, the choice of pole-pair number is influenced largely by considerations of size and magnetic flux leakage.

The force capability is determined by the specific magnetic loading, i.e. airgap flux density. In surface mounted magnet configurations, scaling-up the magnet thickness can produce considerable increase in airgap flux density, albeit at a diminishing rate of return. The specific force capability could also be adjusted by adding a stator mmf using a meandered wound coil in between the magnets (Fig. 4) which can increase or decrease the amplitude of the constant force characteristic.
respectively. For the ACFA, and ACFA topologies.

III. SIMPLIFIED ACTIVE ANALYTICAL MODEL

When considering the active device of Fig. 4, the equations of Section II are somewhat changed, where the stator mmf has to be included, since

\[ H_{S_c} + H_m l_n = mmf_c, \]  

and, consequently, the average airgap flux density is given by:

\[ B_m = \frac{\mu_0} {l_m} \left( \frac{mmf_c + mmf_f} {l_m + \frac{1}{l_m} S_n} \right), \]

where the magnet and stator coil magnetomotive force are respectively \( mmf_c \) and \( mmf_f \) (the sign of the \( mmf_f \) source is negative if its direction opposes the magnet mmf). The permanent magnet and stator mmfs are given by:

\[ mmf_m = 2 B_m x_3 \quad \text{and} \quad mmf_f = 2I_c, \]

respectively. For the ACFA, \( I_c \) represents the current in the coil and, only for verification purposes, a value of 1300 A is defined here, which, considering the coil cross-sectional geometry calculated with the parameters given in Table II, corresponds to a current density of approximately 12.3 A/mm\(^2\). This current density is characteristic for a well cooled device or ACFA with a reduced duty-cycle. This relative high density is used to illustrate the ACFA force capability in Section V. A reduced current density could also be applied, albeit at a corresponding reduction in force amplitude.

The ACFA force response, assuming infinitely permeable iron and no leakage fluxes, is given by:

\[ F_c = \frac{-\mu_0} {2} c_a \left( \frac{mmf_c + mmf_f} {l_m} \right)^2 \ldots \]

\[ \left[ \frac{1} {1 + \frac{l_5 c_{15}} {l_m c_{15}}} \right] \]

where the expressions for calculating the various parameters considered for the force calculation of the specific dimensions are summarized in Table I for low pole number (LPN) and high pole number (HPN) PCFA and ACFA topologies.

The analytical expressions (10) and (14) provide fast means to determine the specific force amplitude for passive and active CFAs. However, by not taking into account the aforementioned leakages and flux defocusing, the calculated force response calculations are somewhat different from the force calculated by FEA, and consequently the resulting force response is independent of the \( Z \)-displacement. However, by considering the magnetic flux leakage paths this is true only for 80-90% of the actuator stroke \((z\)-displacement), as will be discussed in Section V. Therefore, the leakage fluxes have to be considered in order to approximate the force response behavior for the entire stroke and, thus, a more comprehensive magnetic equivalent model (MEC) model is considered in the next section.

IV. PASSIVE AND ACTIVE MAGNETIC EQUIVALENT CIRCUIT MODELS

This magnetic equivalent circuit for the LPN/HPN PCFA (Figs. 2 and 3) is built based on a set of simplifying assumptions: the iron cores are considered linear with a relative magnetic permeability of 1000, the 2\(^{nd}\) quadrant of PM BH-characteristic is considered linear with a remanent flux density of 1.23 T and its relative recoil permeability is approximated to 1. The axial cross-section, as shown by Fig. 5, shows the motivation for considering a two part MEC model to estimate the force-displacement characteristic. The first part is used for the overlapping part of the actuator, indicated by \( R_1 \) and \( R_2 \) in Fig. 5, and the second part is used to model the non-overlapping part and the end-leakages, indicated by \( R_8 \) and \( R_{10} \) in Fig. 5.

The various reluctance elements of these magnetic equivalent circuit models are calculated starting from the general expression:

\[ R = \frac{l S} {\mu_0 \mu_r}, \]

where a flux tube (flux path) is characterized by the length, \( l \), the vacuum and the material relative permeability, \( \mu_0 \) and \( \mu_r \), and the cross-sectional area, \( S \), which may vary with the length. Following the radial \((R-\theta)\) cross-section of the overlapping part of the PCFA, as shown in Fig. 6, the magnetic equivalent circuit from Fig. 6 can be recognized. A more detailed analysis of the MEC method and theory of the magnetic flux tubes can be found in [4] and [5].

In Fig. 7, the PM is modeled as an mmf source in series with a reluctance element, \( R_I \). The airgap and fringing magnetic flux distribution are modeled by the corresponding reluctance elements, \( R_2 \) and \( R_8 \). The stationary and translating back-iron parts are modeled by the circuit elements, \( R_9 \) and \( R_{10} \), respectively. The main magnetic flux path is indicated by \( \Phi_i \). The flux is determined by using a system of linear equations resulting from Kirchhoff’s laws:

\[ [A] \cdot [\Phi] = [f] \]

where \([A]\) is the coefficient matrix given by:
\[ A^{(1)} = \begin{bmatrix} -1 & -1 & 1 \\ R_1 & -R_8 & 0 \\ R_5 & 0 & R_4 \end{bmatrix}, \quad (17) \]

and

\[ R_{r1} = 2R_1 + 2R_2 + R_3. \quad (18) \]

The unknown magnetic fluxes are written in the column matrix:

\[ \begin{bmatrix} \Phi^{(1)} \end{bmatrix} = \begin{bmatrix} \Phi_1 \\ \Phi_2 \\ \Phi_3 \end{bmatrix}, \quad (19) \]

and the column vector \( f^{(1)} \) is:

\[ \begin{bmatrix} f^{(1)} \end{bmatrix} = \begin{bmatrix} 0 & 0 & mmf_{mmf} \end{bmatrix}. \quad (20) \]

The system is solved for the magnetic fluxes, and the flux densities are calculated with:

\[ B_i = \frac{\Phi_i}{A_i}. \quad (21) \]

The force in \( z \)-direction is represented by (9), where the co-energy is calculated with:

\[ W' = \sum B_i^2 \frac{1}{2\mu_0} S_i l_i. \quad (22) \]

The resultant force will have only a \( z \)-component since the radial component is canceled giving the axial symmetry of the model.

\[ \begin{bmatrix} 1 & -1 & -1 & 0 & 0 \\ 0 & 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} A^{(2)} \\ \Phi^{(2)} \end{bmatrix} = \begin{bmatrix} f^{(2)} \end{bmatrix}, \quad (23) \]

where the coefficient matrix is:

\[ A = \begin{bmatrix} 2R_6 & 0 & R_3 & R_5 & 0 \\ 0 & 0 & R_3 & 0 & -R_11 \\ 2R_8 & R_3 & 0 & R_3 & 0 \end{bmatrix}, \quad (24) \]

for the second MEC model presented in Fig. 9. The reluctance \( R_{r2} \) in this model represents the summation of two times \( R_{10} \) and \( R_p \). The unknown magnetic fluxes are written in the column matrix:

\[ \begin{bmatrix} \Phi^{(2)} \end{bmatrix} = \begin{bmatrix} \Phi_4 \\ \Phi_5 \\ \Phi_6 \\ \Phi_7 \\ \Phi_8 \end{bmatrix}, \quad (25) \]

and the column vector \( f^{(2)} \) is:

\[ \begin{bmatrix} f^{(2)} \end{bmatrix} = \begin{bmatrix} 0 & 0 & mmf_{mmf} & 0 & mmf_{mmf} \end{bmatrix}. \quad (26) \]

The derivation of the ACFA’s MEC model follows the same pattern, however the reluctance elements representing the stator back-iron and magnets are adjusted to the ACFA topology and the equations (20) and (26) are updated to include the stator mmf.
\[
\begin{bmatrix}
0 & 0 & \text{mmf}_m + \text{mmf}_f
\end{bmatrix},
\]
for the simplified model of Fig. 7 and:
\[
\begin{bmatrix}
0 & 0 & \text{mmf}_m + \text{mmf}_f \\
& & \ldots \\
0 & \text{mmf}_m + \text{mmf}_f
\end{bmatrix}
\]
for the model of Fig. 9.

In these models all the various reluctances, are determined from (15) and the parameters and dimensions indicated in Tables I and II. The total force is derived using (9), where the total magnetic energy (22) is given by the sum of the energies of both parts of the MEC model. In order to illustrate the respective accuracy of the analytical and the MEC model, they are verified versus finite element models in the next section.

V. FINITE ELEMENT ANALYSIS AND RESULTS COMPARISON

The PCFA and ACFA are numerically modeled by means of non-linear 3D FEA software. The electromagnetic force produced on the translator is obtained using the virtual work approach. The 3D magnetostatic solver of the Maxwell 3D (Ansoft Co.) software package is used to perform the analysis.

Around the PCFA/ACFA an air domain is considered, which extends to four times the largest dimension of the actuator in the R and Z direction, respectively. As an example, in the case of the HPN ACFA, the mesh is composed of 55000 tetrahedral elements for an axial length of 250 mm. A reduced model of the actuator (Fig. 10) can be used by exploiting the model’s periodicity. The advantage of this reduced model is the significant improvement in calculation time. Moreover, due to the constant output force characteristic of the actuator over an extended part of the stroke, the model can be further restricted to a single position which brings a significant reduction in calculation time for the optimization routine from Section VI. The evaluated dimensions and geometry for the LPN and HPN are given in Table II and Fig. 11. It needs noting that these dimensions represent by no means optimized designs and are considered solely to investigate the proof of principle. They are used to illustrate the potential of the proposed actuator configurations. The permanent magnets are assumed to be a sintered NdFeB with a remanence of 1.23T, where in the finite element analysis the standard non-linear BH-curve for mild steel, Al 1010, is used for both stationary and translating back-iron.

The force-displacement characteristic of the LPN PCFA (Fig. 2) is calculated by means of the analytical model of Section II, the MEC model of Section IV and the finite element analysis, where 0 mm (initial position) corresponds to the stator with magnets and translator being fully aligned. Fig. 12 shows the results, for both parallel and radial magnetization patterns in the FE analysis, where a near constant force versus displacement is achieved for approximately 90% of the Z-displacement.

![Figure 10. Reduced 3D finite element model of the CCFA with boundary conditions.](image)

![Figure 11. Schematic of the dimensions of both the PCFA (upper part) and ACFA (lower part), where \(\alpha\) and \(x_5\) determine the coil area.](image)

<table>
<thead>
<tr>
<th>(x_1) (mm)</th>
<th>LPN PCFA</th>
<th>HPN PCFA</th>
<th>HPN ACFA</th>
</tr>
</thead>
<tbody>
<tr>
<td>(x_2) (mm)</td>
<td>5.0</td>
<td>9.0</td>
<td>9.0</td>
</tr>
<tr>
<td>(g) (mm)</td>
<td>8.0</td>
<td>10.0</td>
<td>10.0</td>
</tr>
<tr>
<td>(x_3) (mm)</td>
<td>3.0</td>
<td>5.0</td>
<td>5.0</td>
</tr>
<tr>
<td>(x_4) (mm)</td>
<td>5.0</td>
<td>15.0</td>
<td>15.0</td>
</tr>
<tr>
<td>(x_5) (mm)</td>
<td>N/A</td>
<td>N/A</td>
<td>15.0</td>
</tr>
<tr>
<td>(\alpha) (°)</td>
<td>N/A</td>
<td>N/A</td>
<td>6.0</td>
</tr>
</tbody>
</table>

From Fig. 12 it can be concluded that a relatively large discrepancy is apparent between the analytical and the MEC models. This is due to the non-inclusion of the various leakages. Including these elements, such as in the MEC model shows that a reasonable force level estimate is achieved, albeit that still some 20% overestimate is visible of the force-displacement compared to the parallel magnetization and underestimate for the radial magnetization. For this PCFA using radial magnetized magnets clearly brings an advantage of 50% over parallel magnetization. Therefore, only the radial magnetization is used in the models of the HPN PCFA and ACFA. However, this relative increased force level of the radial magnetization will be likely diminished at an ever increasing number of poles. The available stroke of the actuators determined by the axial length, for these examples an axial length of 50 mm is selected to minimize the computational requirement.
However, very long strokes can be achieved and this is visible in Fig. 12 where the constant force region is about 90% of the stroke at 50 mm.

It needs noting that an increase in axial length does not result in an increase in force level in the rise and drop-off regions of the force-displacement characteristic, hence, an increased percentage of the stroke exhibits constant force. For example, a PCFA with a length of 250 mm has a constant force response over a stroke of 235 mm, respectively 94%.

Further, by varying the airgap diameter of the device and the number of poles, as with the HPN PCFA shown by Fig. 3, an increased force response can be reached. The increase in the number of poles reduces the length of the mean flux tubes, flux lines, hence, the back-iron thicknesses can be reduced for the same flux density level. The achieved force-displacement characteristic, for radial magnetization and the dimensions summarized in Table II for HPN PCFA, is shown in Fig. 13. For a more clear presentation and since the MEC model provides a better approximation, the analytical response is not included in Fig. 14.

The analytical equations (10) and (14) result in the passive, high active and low active force levels of respectively 274 N, 439 N and 147 N. This shows that the relative differences between the models are the same, respectively: an increase of 60% for analytical, 63% for MEC and 66% for the FEA in force level for current aiding to the magnet mmf and a decrease of 44%, 44% and 50% in force level for a current opposing the magnet mmf. Further, this shows the applicability of an ACFA and the better, compared to analytical, suitability of the MEC model to estimate the amplitude of the force. The MEC model does estimate the force to within approximately 5% and captures the behavior of the passive and active force displacement characteristic even in the initial non-linear region. It needs noting that current aiding the magnet mmf provides an increased added force to the passive characteristic, which is due to the square relation between flux density and force. Effectively, for a passive force of 177 N, a coil current, with a density of 12.3 A/mm², aiding to the magnet mmf provides 116 N increase in force amplitude. This would allow for the compensation of a mass variation of approximately 12 kg.

VI. ACTUATOR OPTIMIZATION

In order to minimize the volumetric space envelope and costs, different optimization routines can be used, where, most commonly, the design optimization is undertaken by
means of some analytical or quasi-static numerical calculations. The disadvantage is that these methods require simplifying assumptions that can form a problem, where, most commonly, optimization results are analyzed by FEA. Furthermore, the required 3D FEA for this case is very computationally expensive and thus not the preferred method to be directly used for optimization purposes. Therefore, it can be envisaged that an optimization routine which uses the advantage of an iterative approach, involving both analytical and finite element models, will provide for a relatively fast optimization routine. In this paper an aggressive output space-mapping (AOSM) optimization technique is considered [6], since it exploits the advantages of the analysis methods of Section III, i.e. the very low computational effort of the analytical expressions and the accuracy of the FEA from Section V. Following from the space-mapping (SM) terminology, i.e. the surrogate based family of optimization techniques from which AOSM is derived, the analytical and finite element (FE) models are denoted as coarse, \(c(x)\), and fine models, \(f(x)\), respectively. It needs noting that the analytical model is chosen as a coarse model, instead of the more accurate MEC model, with the purpose of showing that SM techniques can successfully exploit very coarse models. An extra equation,

\[ T_s = P(h_s S_t) + T_c, \]

is considered in the coarse model for calculating the stator surface temperature, based on the assumption that at steady-state the temperature is evenly distributed within the actuator volume. In this equation, \(P\) is the total copper loss, \(S_t\) is the outer stator surface, \(h_s\) is the convection coefficient and \(T_c\) represents the ambient temperature.

Various approximation methods are used in optimization to derive surrogate models for computationally expensive functions. The particularity of the SM techniques is that a mathematical approximation is used to correct a coarse model which inaccurately represents the physical phenomenon, where a linear response mapping is considered in AOSM. The mapping is constructed using an approximated Jacobian, \(B_k\), obtained from a rank-one Broyden update.

The core of the AOSM algorithm is structured as follows:
1. Initialize \(k = 1\), \(B_1 = I\) (identity matrix), \(y_1 = y\) and \(\varepsilon_0\).
2. Compute a fine model iterate:
   \[ x_k = \underset{x \in X}{\arg \min} \| c(x) - y_k \|. \]
3. Evaluate \(f(x_k)\) and the error residuals.
4. Stop and set \(x^* = x_k\) if an appropriate error criterion is satisfied.
5. If \(k \geq 2\) update \(B_k\):
   \[ B_k = B_{k-1} + \frac{c(x_{k-1}) - c(x_k) - B_{k-1} h_{k-1} y_k}{h_{k-1}^T h_{k-1}}, \]
   where
   \[ h_{k-1} = f(x_{k-1}) - f(x_{k-1}). \]

6. Set \(y_{k+1} = c(x_k) + B_k (y - f(x_k))\), \(k = k + 1\) and go to step 2.

As it can be observed from the step 2 of the algorithm, the expensive fine model is replaced by the coarse one and the mapping is employed to alter, in step 6, the design specification, \(y_k\), in respect to which the coarse model is re-optimized. This approach brings the advantage that the mapping is used only once per iteration.

The HPN ACFA configuration investigated in Section V, provides a passive constant force amplitude of approximately 200 N, i.e. with an airgap diameter of 59.5 mm, magnet thickness of 5 mm, over 60% of the stroke. This actuator design is optimized in this section to achieve the 200 N passive response level and 50% adjustability considering a current density of 12.5 A/mm² with a packing factor of approximately 0.6, as summarized in Table III. The relatively high filling factor, for the meander wound coil, can be applied since the design does not include any tooth tips, as shown in Fig. 11.

The design variables \((x_{1...5}\) and \(\alpha\) ) are presented in Fig. 11. The optimization objective is the minimization of the translator mass and actuator outer radius for the above specified force outputs, where, giving the fixed axial length, the objective translates into a force maximization problem. The accomplished optimized long-stroke constant-force displacement characteristic of the ACFA, for the design of Fig. 4, is shown in Fig. 15, which is achieved with a translator mass of 0.66 kg, axial length, \(L_{ax}\), of 50 mm and 16 permanent magnet poles, where the design specifications, constraints and the various dimensions are summarized in Table III and IV.

### Table III

<table>
<thead>
<tr>
<th>Type</th>
<th>Values</th>
</tr>
</thead>
<tbody>
<tr>
<td>Design specifications</td>
<td>(F_{passive} = 200, N, F_{active} = 300, N), (B_{mech} \leq 1.6, T, T_{ref} \leq 130^\circ C)</td>
</tr>
<tr>
<td>Other data</td>
<td>(J = 12.5, A/mm^2, \beta = 0.6) (filling factor)</td>
</tr>
<tr>
<td>(\rho_{eff} = 7650, kg/m^3, \rho_{air} = 8900, kg/m^3), (\rho_{sat} = 7700, kg/m^3, B_0 = 1.23, T)</td>
<td></td>
</tr>
<tr>
<td>(h = 20, W/m^2/K)</td>
<td></td>
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</tbody>
</table>

### Table IV

<table>
<thead>
<tr>
<th>Geometrical parameters of the ACFA</th>
<th>Non-optimized</th>
<th>Optimized</th>
</tr>
</thead>
<tbody>
<tr>
<td>(x_1) (mm)</td>
<td>49.0</td>
<td>56.6</td>
</tr>
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<td>(x_2) (mm)</td>
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<td>(x_3) (mm)</td>
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<td>1.0</td>
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<td>(x_4) (mm)</td>
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<td>(x_5) (mm)</td>
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<td>(x_6) (mm)</td>
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</tr>
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<td>(\alpha) (°)</td>
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<td>8.7</td>
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<td>(R_{outer}) (mm)</td>
<td>95.0</td>
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<td>(L_{ax}) (mm)</td>
<td>50</td>
<td>50</td>
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<td>Translator mass (kg)</td>
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<td>Total mass (kg)</td>
<td>5.13</td>
<td>3.62</td>
</tr>
</tbody>
</table>
It needs noting that the design solution is obtained within ten AOSM iterations which results in the necessity of performing only ten FE evaluations. The optimization procedure of the ACFA provides for a significant improvement. For example, the outer radius is reduced from 95.0 mm to 85.1 mm and the inner radius is increased from 49.0 mm to 56.6 mm.

In the above analysis the volume force density level it does not really have a meaning, since increasing the axial length does not provide for an increase in the force level, however does provide an increased stroke and constant force versus displacement. Therefore, the axial length can be reduced to have a small linear region, as shown in Fig. 16 for an axial length of 20 mm. In this case the force density is maximized and for the optimized ACFA this provides for: 788 kN/m$^3$ passively and 1182 kN/m$^3$ actively, however this reduces to, using the characteristic of Fig. 15 with an axial length of 50 mm: 315 kN/m$^3$ and 473 kN/m$^3$, respectively. When considering the axial length of 250 mm this reduces further to 63 kN/m$^3$ and 95 kN/m$^3$.

However, since the force amplitude is not influenced by the axial length, it might be advantageous to consider the force per area, i.e. radial cross-section, instead of volume. This approach illustrates the optimization effectiveness, i.e., a force density of 15.7 kN/m$^2$ is obtained for the optimized ACFA, compared with 9.6 kN/m$^2$ in the initial configuration.

VII. CONCLUSION
Passive and active long stroke constant-force actuator topologies are investigated in this paper. This type of actuators exhibits a constant force characteristic for approximately 90% of the stroke, and, furthermore, relatively very long strokes can be achieved. Although that it is also interesting to note that implementing some compensating rings, as discussed in [7], can increase the initial force to be nearly constant, hence, an almost constant force-displacement characteristic, i.e. for approximately 95% of the axial length, could be achieved even for relative short axial lengths. This paper also shows that an active version of this actuator obtained, for example, by means of implementing a meandered wound coil in between the permanent magnets, can adjust the amplitude of the force-displacement response by approximately 50%, i.e. adding to or subtracting from the passive characteristic.

The paper presents the means to determine the force-displacement characteristics for constant-force actuators. Furthermore, an optimized ACFA design is investigated. The actuator provides a constant force-displacement output characteristic and the translator mass has been minimized for the specified force levels. The optimized ACFA is designed considering both electromagnetic and thermal aspects, for which the AOSM technique, i.e. an optimization routine based on the SM, is successfully implemented. The results show that a significant improvement of the initial design is obtained.

REFERENCES