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Polarization-sensitive near-field investigation of photonic crystal microcavities

Silvia Vignolini,1,a Francesca Intonti,2 Francesco Riboli,1 Diederik S. Wiersma,1 Laurent Balet,2 Lianhe H. Li,3 Marco Francardi,4 Annamaria Gerardino,4 Andrea Fiore,5 and Massimo Gurioli2
1LENS and INFM-BEC, Via Nello Carrara 1, 50019 Sesto Fiorentino Italy
2CNISM, Unità di Ricerca di Firenze and Department of Physics, Via Sansone 1, 50019 Sesto Fiorentino, Italy
3EPFL Institute of Photonics and Quantum Electronics, Station 3, CH-1015 Lausanne, Switzerland
4Institute of Photonics and Nanotechnology, CNR, via del Cineto Romano 42, 00156 Roma, Italy
5COBRA Research Institute, Eindhoven University of Technology, 5600 MB Eindhoven, The Netherlands

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We report on polarization sensitive imaging of two-dimensional photonic crystal microcavity modes. By using a near-field scanning optical microscope with a polarization sensitive setup, it is possible to selectively map, with a resolution beyond the diffraction limit, each electric field component in the plane of the sample. In addition, the simultaneous analysis of photoluminescence maps in different polarization channels allowed us to obtain important insight on near-field microscopy detection mechanism. Finite difference time domain simulations confirm the experimental results.

In the past few years, photonic crystal microcavities (PC-MCs) have been intensively studied and improved both for their potential applications in the field of optoelectronics as well as for their fundamental physical properties.1,2 The geometry of the PC-MCs defines both the spatial distribution and the polarization properties of the optical modes, and the possibility to control and combine these two degrees of freedom is of the utmost interest for the realization of several applications, such as PC based lasers3 or photonic fibers.4

Up to now, the PC-MCs polarization properties have been studied mainly by far-field microscopy.5–7 In order to achieve a better spatial resolution, scanning near-field microscopy has already proved to be a powerful tool for addressing the optical modes of PC based devices.8–20 Moreover, it was recently demonstrated that the spectral shift of the cavity modes induced by the local probe directly maps the electric field intensity, providing an additional imaging option with higher spatial resolution.15–18 However, even if particular configurations of near-field optical microscopes allow to control the polarization properties of light19,20 a direct correlation between the spatial map and the polarization of the electromagnetic modes of PC-MCs has not yet been obtained.

In this paper we demonstrate, by using a scanning near field microscope with polarization control, that it is possible to completely retrace the vectorial maps of the electric field associated to the optical modes. By exploiting the advantage of the tip-induced spectral shift we are able to map the total electric field intensity associated to the mode. At the same time, the intensity maps of the photoluminescence (PL) signal of the mode, in the two orthogonal polarization channels, provide the imaging of the two electric field components in the plane of the PC-MC membrane.

The sample under consideration is a two dimensional PC-MC on a suspended membrane incorporating quantum dots (QDs) acting as local light sources. The sample consists of a GaAs based heterostructure: three layers of high-density InAs QDs emitting at 1300 nm are grown by molecular beam epitaxy at the center of a 320-nm-thick GaAs membrane. The membrane is grown on top of a 1500-nm-thick Al0.7Ga0.3As sacrificial layer.22 The studied structure consists of a two-dimensional triangular lattice of air holes with lattice parameter \(a=301 \text{ nm}\) and filling fraction \(f=35\%\), where the cavity is formed by four missing holes organized in a diamondlike geometry. An in-plane scanning electron microscope image of the PC-MC is reported in Fig. 1(a), together with the \((x, y, \theta)\) reference system, to which we will refer for the polarization configurations; the direction \(\theta=0^\circ\) (\(\theta=90^\circ\)) denotes the \(x\) (\(y\)) axis. A preliminary far-field characterization of the sample (via micro-PL at room temperature) is obtained using a 100× microscopy objective and a single-mode optical fiber with a 5 \(\mu\)m core diameter acting as confocal pin-hole. The polarization-dependent spectrum of the cavity mode signal is measured by using a linear polarizer and by rotating a half-wave plate in front of the collection. Figure 1(b) shows the \(\theta\) angular dependence of the PL peak intensity for the two main cavity modes, hereafter labeled M1 [black circles in Fig. 1(b)] and M2 [red squares in Fig. 1(b)], centered around 1263 and 1240.5 nm, respectively [see Fig. 1(c)]. The mode M1 is mainly polarized along the \(x\) direction (the ratio between opposite polarizations in the plane is 1:100), while the mode M2 is characterized by an elliptical polarization (the ratio between opposite polarizations in the plane is 1:2) along the \(y\) direction. The experimental spectral and polarization features of the two cavity peaks is well reproduced by calculations obtained by using a commercially available finite-difference time-domain (FDTD) solver package (Crystal Wave), as reported in Fig. 1(d), where the simulated emission spectra for the two in-plane components of the electric field are shown.

In order to map the spatial distribution of the electric field associated to these modes, we use a commercial scan-
nel. A pretty large extinction factor the mode M1 disappears embedded QDs is coupled to the same probe, passes through on the sample surface. A portion of the PL signal of the photonic structure.

Moreover, since the selection of the polarization pressure and rotation to the fiber. By changing the configuration, the sample is excited with light from a diode laser (780 nm) coupled into a chemically etched, uncoated near-field fiber probe, that is raster scanned at a constant height on the sample surface. A portion of the PL signal of the embedded QDs is coupled to the same probe, passes through a linear polarizer, is dispersed by a spectrometer (resolution 0.1 nm), and finally it is detected by a cooled InGaAs array. All the data reported in this letter refer to room temperature.

We control the polarization of the light in our SNOM experiment by using a non-polarization maintaining fiber and a polarization compensator acting on the fiber. Each photon, collected by the tip, changes its polarization state during the propagation along the fiber. This unknown polarization change is, however, identical for all the collected photons and therefore can be totally compensated. The polarization compensation is obtained using a system based on the Babinet–Soleil compensator that permit to apply a controlled pressure and rotation to the fiber. By changing the configuration of the polarization compensator and fixing the angle of the polarizer in front of the spectrometer, we can selectively collect only one polarization channel in the plane of the membrane. Moreover, since the selection of the polarization channel takes place at the end of the fiber, it does not play a role in the interaction between the near-field probe and the photonic structure.

A test of the polarization control in SNOM spectra is reported in Fig. 1(c), where the near-field spectra (averaged in a region of $2 \times 2 \ \mu m^2$) for two different polarization channels are provided. The red curve in Fig. 1(c) refers to the spectrum for the $E_y$ polarization channel (where the signal of the mode M1 disappears) while the black line in Fig. 1(c) refers to the spectrum in the orthogonal $E_x$ polarization channel. A pretty large extinction factor (1:100) is obtained in the SNOM spectra, denoting the sensitivity of our setup. Since the mode M1 is prevalently characterized by only one component of the electric field, in the following we will focus the discussion on the results for the mode M2.

In order to obtain a high resolution image of the electric field intensity associated to the mode and to discuss the possible degenerate nature of M2, it is convenient to study the tip-induced spectral shift of the optical mode. In particular, we want to address the point whether the resonance M2 is composed by two orthogonal modes at the same wavelength or it is a single nondegenerate mode with both x and y components of the electric field. The tip-induced spectral shift depends on both the geometry of the tip and on the electric field intensity associated to the mode itself. By assuming that the tip has a circular geometry, we expect that its interaction does not depend on the electric field orientation. However, in the case of degenerate modes with different electric field spatial distribution and orthogonal polarization, a local perturbation acts differently for the two modes, possibly breaking the degeneracy. As a consequence, for such kind of modes, we expect to observe different spatial maps of the tip-induced spectral shift in the two orthogonal polarization channels. On the contrary, in the case of nondegenerate modes, the spectral shift maps are independent on the polarization configuration in the detection. Figure 2 shows (in a blue-to-white color scale) the results of the measured tip induced spectral shift for the mode M2 in two perpendicular polarization configurations ($x$ and $y$, respectively), as compared with the theoretical calculation of the electric field intensity, Fig. 2(e). The comparison between experimental and simulated data demonstrates that there is a direct correspondence between the frequency shift induced by the tip and the unperturbed electric field profile of the mode, accordingly to the findings in Refs. 15 and 16. More relevant, for the topic addressed here is the experimental observation that the spectral shift maps are also independent on the polarization configuration in the detection, including the absolute value of the mode shift and therefore we can conclude that the mode M2 is a single nondegenerate mode. The tip-induced spectral shift map gives a direct and high fidelity experimental imaging of the electromagnetic local density of states but, at least for nondegenerate modes, it does not provide selective information on the different electric field components.

Figure 3 shows, in the last column, the spatial distribution of the PL signal associated to mode M2 for three different polarization configurations ($\theta=0^\circ, 90^\circ, 45^\circ$) compared with the calculated intensity of the electric field components at different distances $d$ from the membrane surface (first four
The procedure of using an effective distance maps fitting parameter is the distance maps are well reproduced by the calculated intensity associated simulated PL maps clearly shows that all three experimental sample surface. The value of map measured by uncoated tips can be retraced by calculating this problem it is commonly assumed that the near-field PL intensity distribution in the sample surface. The value of the SNOM PL intensity map differs from the electric field axis of the sample weighted by the tip geometry. Therefore, the electric field intensity distribution averaged over the normal direction of the sample surface is related to the collected signal is related to the electric field intensity distribution averaged over the normal axis of the sample weighted by the tip geometry. Therefore, the SNOM PL intensity map differs from the electric field intensity distribution in the sample surface. The value of $d$ does not represent the real height of the tip but it is an effective free parameter that deals with the fact that the collected signal is related to the electric field intensity distribution averaged over the normal axis of the sample weighted by the tip geometry. Therefore, the SNOM PL intensity map differs from the electric field intensity distribution in the sample surface.

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In conclusion, we demonstrate that it is possible to obtain polarization sensitive maps of PC cavity modes. In addition, the simultaneous analysis of different PL maps with the polarization control allowed us to obtain important insights on SNOM detection mechanisms.

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