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Waveguide superconducting single-photon detectors for integrated quantum photonic circuits

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The monolithic integration of single-photon sources, passive optical circuits, and single-photon detectors enables complex and scalable quantum photonic integrated circuits, for application in linear-optics quantum computing and quantum communications. Here, we demonstrate a key component of such a circuit, a waveguide single-photon detector. Our detectors, based on superconducting nanowires on GaAs ridge waveguides, provide high efficiency (~20%) at telecom wavelengths, high timing accuracy (~60 ps), and response time in the ns range and are fully compatible with the integration of single-photon sources, passive networks, and modulators. 

The combination of single-photon sources, passive optical circuits, and single-photon detectors enables important functionalities in quantum communications, such as quantum repeaters and qubit amplifiers, and also forms the basis of all-optical quantum gates. However, present implementations are limited to few qubits, due to the large number of optical components required and the corresponding complexity and cost of experimental set-ups. The monolithic integration of quantum photonic components and circuits on a chip is absolutely required to scale implementations of optical quantum information processing to meaningful numbers of qubits. The integration of passive circuits has been demonstrated in waveguides based on silica-on-silicon and on laser-micromachined glass, but a platform for the simultaneous integration of sources, detectors, and passive circuitry is still missing. The integration of detectors is particularly challenging, as the complex device structures associated to avalanche photodiodes are not easily compatible with the integration with low-loss waveguides and even less with sources. Transition-edge sensors may be suited for integration, but they are plagued by very slow response times (leading to maximum counting rates in the tens of kHz range) and require cooling down to <100 mK temperatures. Here, we report a simple approach to the realization of single-photon detectors on optical waveguides in the GaAs/AlGaAs material system. It enables the demonstration of efficient waveguide single-photon detectors (WSPDs) and can provide ultrahigh sensitivity at telecommunication wavelengths, high counting rates, broad spectral response, and high temporal resolution due to low jitter values. In our design (see Fig. 1), the wires are deposited and patterned on top of a GaAs ridge waveguide, in order to sense the evanescent field on the surface. Four NbN nanowires (4 nm-thick, 100 nm wide, and spaced by 150 nm) are placed on top of a GaAs (300 nm)/Al0.75Ga0.25As waveguide, and a 1.85 μm-wide, 250 nm-deep ridge is etched to provide 2D confinement. We assume that a 100 nm-thick SiOx layer is left on top of the wires as a residue of the hydrogen silsesquioxane (HSQ) mask used for single-photon detectors (SSPDs) and can provide ultrahigh sensitivity at telecommunication wavelengths, high counting rates, broad spectral response, and high temporal resolution due to low jitter values. In our design (see Fig. 1), the wires are deposited and patterned on top of a GaAs ridge waveguide, in order to sense the evanescent field on the surface. Four NbN nanowires (4 nm-thick, 100 nm wide, and spaced by 150 nm) are placed on top of a GaAs (300 nm)/Al0.75Ga0.25As waveguide, and a 1.85 μm-wide, 250 nm-deep ridge is etched to provide 2D confinement. We assume that a 100 nm-thick SiOx layer is left on top of the wires as a residue of the hydrogen silsesquioxane (HSQ) mask used for

FIG. 1. (Color online) (a) Schematic view and (b) Contour and vector plot of the amplitude [V/m] and direction of the electric field for the fundamental mode (λ = 1500 nm) of the waveguide superconducting single-photon detector.

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The waveguides were cleaved leaving a 1 mm-long passive ridge waveguide between the cleaved facet and the WSPD. The WSPDs were characterized by end-fire coupling light from a polarization-maintaining lensed fiber (producing a spot with nominal diameter of 2.5 ± 0.5 μm) into the waveguides mounted on the cold finger of a continuous flow helium cryostat. Both the lensed fiber and the contact probes are mounted on piezoelectric positioners, which are thermally anchored to the cold plate to minimize the thermal load to the detector, resulting in an operating temperature <4 K. The inset of Fig. 3 displays the current-voltage characteristic measured for a 50 μm-long WSPD, showing a critical current (Ic) of 16.9 μA. The electro-optical response was measured by end-fire coupling a continuous wave 1300 nm diode laser through the lensed fiber in the TE polarization. The detector count rate was observed to be extremely sensitive to the fiber-waveguide alignment and to their distance, confirming that the detector responds to guided photons and not to stray light propagating along the surface or in the substrate. The count rate was measured to be proportional to the laser power (Fig. 3), proving operation in the single-photon regime. The inset in Fig. 4 shows a single WSPD output pulse, showing a pulse duration (full-width-half-maximum) of 3.2 ns and a 1/e decay time of 3.6 ns, which corresponds very well to the expected time constant τ = Lkin/R = 3.6 ns, where Lkin = 180 nH is the wire kinetic inductance (as calculated from the kinetic inductance per square for similar NbN wires, L = 90 pH/μm (Ref. 12)).

FIG. 2. (Color online) (a) Scanning electron microscope (SEM) micrograph of Ti/Au electric contacts; (b) Collection of three SEM micrographs taken in different regions of a 30 μm long WSPD, the nanowires are still covered by the HSQ etching mask; (c) Atomic force microscopy (AFM) image of the 1.85 μm wide and 30 μm long HSQ mask used for the etching of the waveguide aligned on top of the NbN nanowires; and (d) AFM Enlarged view of the waveguide HSQ etching mask showing a realignment accuracy better than 100 nm.

FIG. 3. (Color online) Count rate as a function of laser power (λ = 1300 nm, TE polarization, Ith = 9.9 μA), showing a linear behavior and hence operation in the single-photon regime. Inset: Current-voltage characteristic of the WSPD, showing a critical current of 16.9 μA, the relaxation-oscillation region, and the beginning of the hot-spot plateau.
and $R = 50 \, \Omega$ is the load resistance. Considering that it takes a time of $\approx 3 \tau$ to recover 95% of the bias current after detection, we estimate a maximum count rate close to 100 MHz. By illuminating the device with a pulsed diode laser, a total jitter of 73 ps was measured on the WSPD output pulse, corresponding to a 61 ps intrinsic detector jitter after correcting for the 40 ps jitter from the laser pulse.

The measured QE (1300 nm, TE polarization) is plotted in Fig. 4 (left axis) as a function of the normalized bias current $I_b/I_C$. The system quantum efficiency (SQE, open symbols), defined as the number of counts divided by the average photon number in the fiber at the input of the cryostat, reaches 3.4% for a 50 $\mu$m-long device. For determining the number of photons coupled into the waveguide, transmission measurements were performed with a tunable laser on a sample containing 3 mm-long ridge waveguides, but without NbN wires and contact pads. From the measured Fabry-Perot fringes, and particularly from the maximum and minimum transmission (in TE polarization), $T_{\text{max}} = 6.1\%$ and $T_{\text{min}} = 1.8\%$, we deduce that the propagation loss over a 3 mm waveguide length is negligible. Assuming symmetric input/output coupling and using the standard expression for the Fabry-Perot transmission,\textsuperscript{13} we derive a coupling efficiency (from fiber input) of 17.4%. The corresponding device quantum efficiency (DQE), defined with respect to the number of photons coupled into the waveguide, is plotted as closed symbols in Fig. 4 and reaches 19.7%. This value is still lower than the calculated absorptance (90% in the 50 $\mu$m-long WSPD), which we mainly attribute to a limited internal quantum efficiency (detection probability upon absorption of a photon), and further improvements of film quality and wire etching process may result in notably improved values. Another potential cause for limited efficiency may be extrinsic loss (e.g., scattering) due to the nanowires, which is however believed to be small as compared to nanowire absorption. The dark count rate was measured in another cryostat without optical windows at 1.2 K and is presented in Fig. 4 (right axis), showing the usual exponential dependence as a function of the bias current.

In conclusion, we have demonstrated integrated waveguide single-photon detectors based on superconducting nanowires on GaAs ridge waveguides. They provide system (device) quantum efficiencies of 3.4% (20%) at 1300 nm, a timing resolution of 60 ps, and dead times of few ns. Further optimization of film deposition and device fabrication may result in efficiencies approaching 100% due to the high absorptance allowed by the waveguide geometry. Higher system QE and polarization-independence can be obtained by a waveguide design providing a more extended and symmetric mode profile and by integrating a tapered coupler.\textsuperscript{14} Integrated photon-correlation devices\textsuperscript{15} and photon-number-resolving detectors\textsuperscript{16} are straightforward to realize by integrating several wires on the same waveguide. Furthermore, this technology is fully compatible with the fabrication of passive quantum circuits on GaAs waveguides, and with single-photon sources based on InAs quantum dots in waveguides, and, therefore, opens the way to fully integrated quantum photonics circuits including sources and detectors.

Note added in proof: Two alternative approaches to waveguide single-photon detection, based on transition edge sensors\textsuperscript{8} and on superconducting nanowires on Si/SiO$_2$ waveguides,\textsuperscript{17} have been reported during the submission process of this manuscript.

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