Het voorspellen van de kniklast van een balk door analyse van de eigenfrequentie bij toenemende axiale belasting

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Abstract: In this paper we consider a system of globally, uniformly and linearly coupled Hindmarsh and Rose oscillators. This model is a reduction of the celebrated Hodgkin-Huxley equations, which are considered as the most physiologically realistic model of neural dynamics at the level of a single cell. Exploiting recently developed framework for analysis of synchronization phenomena - passivity-based approach (A. Pogromsky, H. Nijmeijer) - we derive sufficient conditions for global/local asymptotic synchronization in the system. Apart from simply showing a possibility of synchronization, we also try to estimate the least possible values for the coupling connections that are sufficient for convergence of the trajectories to the synchronization manifold.

Keywords: synchronization, passive systems, spiking neurons

1. INTRODUCTION

The Hindmarsh and Rose model (J.L. Hindmarsh, 1984) is a reduced version of the celebrated Hodgkin-Huxley equations for modelling spike initiation in the squid giant axon (A.L. Hodgkin, 1952). The model governs the dynamics of the current through the neuron depending on the membrane potential and internal currents in the cell. Despite that the membrane potential in the original model was described by PDE, in Hindmarsh and Rose model the equations were reduced to ODE under assumption that the axon is space-clamped.

Lacking the features of real neuron like dependence of the membrane potential on spacial distance from the soma along the axon, solutions of Hindmarsh and Rose equations capture such inherent property of the neuron like spiking in both periodic and bursting regimes depending on the external stimulation (see figure 1 for the illustration, where symbols x, y, z state for the membrane potential, recovery variable and adaptation current correspondingly). Furthermore, for a specific set of the parameters and input currents, the model can exhibit chaotic dynamics (Kaas-Petersen, 1987) which in turn is essential in the applications where human-like associative memory is required with the ability to retrieve more than one pattern simultaneously (A. Rafone, 2003).

It is suggested in (Malsburg, 1981; Malsburg, 1999) that a process of retrieval of the stored patterns is related to spontaneously occurring syn-
chrony in the arrays (or lattices) of the neurons.

For that reason investigation of the conditions for synchronization in the ensembles of nonlinear oscillators given by the model in (J.L. Hindmarsh, 1984) is relevant for both theoretical and experimental studies of human-like processing of information.

Most of the published results in the field are concentrated on numerical investigation of the phenomenon (see for example (D. Hansel, 1992; R. Huerta and Rabinovich, 1998)). According to our knowledge, no successful attempts have been made to attack the problem of synchronization in arrays of Hindmarsh and Rose oscillators analytically and especially from control-theoretic prospective. There are a few publications that try to apply control-theoretic analysis for the model (A.E. Milne, 2001). However, applicability of these and similar approaches is limited by assumptions on availability of internal variables for direct measurements and due to the requirements to apply control efforts to every single equation in the system. Therefore new theoretical framework is to be provided to analyze the conditions of synchronization in the system.

As a starting point for our theoretical analysis a recently suggested technique of passivity-based synchronization has been chosen (Pogromsky, 1998). Within this framework we aim to establish an analytical proof for synchronization in an ensemble of Hindmarsh and Rose models and derive estimates of the coupling strengths for which the synchronization is guaranteed.

The contribution of the paper is as follows: first we derive sufficient conditions for synchronization in a network of Hindmarsh and Rose oscillators. These conditions should neither depend on the bounds of the solutions nor should they result in the growing of the coupling parameter when the number of oscillators is increasing. Once the bound for the coupling parameter is defined, we proceed with a local analysis and provide the conditions for local stability of the synchronization manifold.

2. NOMENCLATURE AND PRELIMINARIES

In this section we specify the mathematical model of a Hindmarsh and Rose oscillator and introduce necessary notations.

A single Hindmarsh and Rose oscillator is defined by the following system:

\[
\begin{align*}
x &= -ax^3 + bx^2 + I + y - z + u \\
y &= c - dz^2 - y \\
z &= \varepsilon (s(x + x_0) - z),
\end{align*}
\]

where \(x\) is the membrane potential, \(y\) — recovery variable and \(z\) — adaptation variable. External stimulation is given by constant \(I\) and input \(u\). Variable \(x\) in (1) is usually considered as a natural output of the cell. Parameters \(a, b, c, d, s, x_0, \varepsilon\) are all positive constants. The values of these parameters are specified in Table 1.

A network of oscillators (1) can be described by the following system:

\[
\begin{align*}
\dot{x}_i &= -ax_i^3 + bx_i^2 + I + y_i - z_i + u_i \\
\dot{y}_i &= c - dz_i^2 - y_i \\
\dot{z}_i &= \varepsilon (s(x_i + x_0) - z_i)
\end{align*}
\]

where index \(i \in \{1, \ldots, n\}\) states for the number of each oscillator in the network, and \(u_i\) is a coupling function between the nodes.

Definition 1. Let coupling functions \(u_i : \mathbb{R}^n \to \mathbb{R}\) be given. Coupling is said to be symmetric iff \(u_i(v_{j}) = u_j(v_{i})\), where \(v \in \mathbb{R}, e_i = (\delta_{ik})_{k=1}^n, \delta_{ik} - \text{Kronecker delta}.\)

Definition 2. Coupling is said to be uniform iff it is symmetric, \(u_i(v_{j}) = u_j(v_{i})\) and \(k, j \neq i\).

Definition 3. Coupling is said to be preserving iff it vanishes on the synchronization manifold.

We restrict ourselves to a class of linear coupling functions:

\[y(\mathbf{z}) = -\Gamma \mathbf{z},\]

where \(y = \text{col}(u_1, \ldots, u_2), x = \text{col}(x_1, \ldots, x_n), \Gamma = (\gamma_{ij})_{n \times n}\) is an \(n \times n\) matrix. The \(i\)-th row of matrix \(\Gamma\) is denoted by symbol \(\Gamma_i\). It is clear that symmetric coupling corresponds to
symmetric matrices \( \Gamma \). In case of uniform coupling it is convenient to factorize matrix \( \Gamma \) as follows:

\[
\Gamma = \gamma (\alpha I_n + \Gamma_0), \quad \Gamma_0 = (1 - \delta_{ij}), \quad \gamma, \alpha \in \mathbb{R}, \tag{4}
\]

where \( I_n \) is the identity matrix of appropriate dimensions.

In our study we exploit passivity-based approach to synchronization. Therefore some additional notations are required for consistency. Consider the nonlinear time-invariant system:

\[
\begin{aligned}
\dot{x} &= f(x) + g(x)u \\
y &= h(x)
\end{aligned} \tag{5}
\]

where \( x(t) \in \mathbb{R}^k \) is the state vector, \( f : \mathbb{R}^k \rightarrow \mathbb{R}^k \), \( g : \mathbb{R}^k \times \mathbb{R}^{m \times k} \rightarrow \mathbb{R}^l \), \( h : \mathbb{R}^k \rightarrow \mathbb{R}^l \), \( f, g, h \in C^1 \); \( u(t) \in \mathbb{R}^m \) and \( y(t) \in \mathbb{R}^l \) are input and output vectors respectively.

**Definition 4.** (Pogromsky, 1998) System (5) is called \( C^r \)-semipassive if there exists a \( C^r \)-smooth, \( r \geq 0 \) nonnegative function \( V : \mathbb{R}^k \rightarrow \mathbb{R}^+ \) and a function \( H : \mathbb{R}^+ \rightarrow \mathbb{R}^+ \) such that for any initial conditions \( x(0) \) and any admissible input \( u \) the following dissipation inequality holds:

\[
V(x(t)) - V(x(0)) \leq \int_0^t (y(s))^T u(s) - H(x(s))ds
\]

for all \( 0 \leq t \leq T_{x,y}, \) where the function \( H \) is nonnegative outside some ball:

\[ \exists \rho > 0 \mid x \mid \geq \rho \Rightarrow H(x) \geq 0 \]

The rest of the paper is organized as follows. In Section 3 we show that system (1) is semipassive with radially unbounded storage function. This fact implies boundedness of the solutions of interconnected system (2) for a class of the coupling functions. Relying on these properties we derive sufficient conditions for global/local asymptotic synchronization in system (2). These are formulated in Sections 4 and 5 respectively. Section 6 concludes the paper.

### 3. BOUNDEDNESS OF THE SOLUTIONS FOR THE COUPLED SYSTEMS

**Proposition 5.** System (1) is semipassive with radially unbounded storage function.

**Proof of Proposition 5.** Consider the following positive-definite function:

\[
V(x,y,z) = \frac{1}{2}(c_1x^2 + c_2y^2 + c_3z^2) \tag{7}
\]

According to Definition 4 the proof is completed if we find nonnegative numbers \( c_1, c_2, c_3 \) such that inequality (6) holds for some nonnegative (outside a ball in \( \mathbb{R}^3 \)) function \( H(\cdot) \). It is easy to verify that the following domain of parameters \( c_i \) suffices these requirements: \( c_i = 1; \ c_2 < c_1 \varepsilon \lambda_i \), \( 0 < \lambda_i < 1, \ i = 1, 2, \ c_3 = \frac{15}{4} \). In particular inequality (6) is satisfied for \( c_1 = 1, c_2 = 0.01, c_3 = 125 \) with

\[
H(x,y,z) = 0.465 \left( \frac{4}{100} \right) (-1.65 + x)^4 - (8.98 + x)^2 - 0.1(-462.2 + x^2)^2 - 1.1(x^2 + 0.05y)^2 - 0.0025(-1.33 + y)^2 - 0.011(-100x + y)^2 - 0.00275(-159 + z)^2 - 0.272(2.02x + z)^2. \tag{8}
\]

The proposition is proven.

**Proposition 6.** Let system (2) be given. Let, in addition, coupling function be given by (3) with positive semi-definite \( \Gamma \). Then solutions of (2) are bounded for any initial conditions.

**Proof of Proposition 6.** According to Proposition 5 each \( i \)-th subsystem in (5) is semipassive with radially unbounded storage function \( V(x_i, y_i, z_i) \). The dissipation inequality for the \( i \)-th system in the ensemble can be written as

\[
V(x_i(t), y_i(t), z_i(t)) - V(x_i(0), y_i(0), z_i(0)) \leq \int_0^t (y_i(s))^T u_i(s) - H(x_i(s), y_i(s), z_i(s))ds
\]

Denoting \( W(x,y,z) = \sum_{i=1}^5 V(x_i, y_i, z_i) \) we obtain

\[
W(x,y,z) = \sum_{i=1}^5 (x_i(s))^T \Gamma_0 \xi_i(s) - H(x_i(s), y_i(s), z_i(s)))ds
\]

where \( H(x,y,z) \) is nonnegative outside a ball in the extended state space. The rest of the proof is straightforward. The proposition is proven.

### 4. GLOBAL UPPER BOUND GAIN

In this section we provide analytically calculated bounds for the coupling parameter which guarantees asymptotic synchronization of an ensemble of linear, uniform and preserving coupled Hindmarsh and Rose oscillators. The results are formulated in Proposition 7:

**Proposition 7.** Let system (2) be given with linear, uniform and preserving with respect to the manifold \( x_1 = x_2 = \cdots = x_n \) coupling with

\[ \gamma > \frac{4.5d^2 + b^2}{n} \]

Then all solutions of the system are bounded and
\[
\lim_{t \to \infty} x_i(t) - x_j(t) = 0, \\
\lim_{t \to \infty} y_i(t) - y_j(t) = 0, \\
\lim_{t \to \infty} z_i(t) - z_j(t) = 0
\]
for any \( i, j \in \{1, \ldots, n\} \).

**Proof of Proposition 7.** According to the conditions of the proposition, coupling function is linear, uniform and preserving with respect to the manifold \( x_1 = x_2 = \cdots = x_n \). Then \( \text{col}(1, \ldots, 1) \in \text{Ker}(\Gamma) \). The last automatically implies that \( \alpha = -(n-1) \) in decomposition (4). Hence according to Gershgorin’s circle theorem, matrix \( \Gamma \) is positive semi-definite. Therefore, it follows from Proposition 6 that solutions of system (2) are bounded.

Let us derive synchronization conditions for (2). Consider the following nonnegative function:

\[
V = 0.5 \sum_{i=1}^{n-1} \left( C_x(x_i - x_{i+1})^2 + C_y(y_i - y_{i+1})^2 + C_z(z_i - z_{i+1})^2 \right), \tag{9}
\]
where \( C_x, C_y > 0 \) are to be defined and \( C_z = C_x/(sa) \). Its time-derivative can be expressed as follows:

\[
\dot{V} = \sum_{i=1}^{n-1} \left( -C_x(x_i - x_{i+1})^2 \left( \frac{x_i^2}{2} + \frac{x_{i+1}^2}{2} \right) + b \left( x_i + x_{i+1} \right) + \gamma \right) + \\
C_y(y_i - y_{i+1})(x_i - x_{i+1}) - C_y d(x_i - x_{i+1}) \\
- C_y d(x_i - x_{i+1})^2 \\
- C_y d(z_i - z_{i+1})^2 \tag{10}
\]

Consider the following term in (10):

\[
C_y(y_i - y_{i+1})(x_i - x_{i+1}) - C_y d(x_i - x_{i+1}) \\
C_y d(x_i - x_{i+1})^2 - C_y d(z_i - z_{i+1})^2 \\
d \left( y_i(t) - y_{i+1}(t) \right) - \left( y_i(t) - y_{i+1}(t) \right) + \left( x_i(t) - x_{i+1}(t) \right) + \left( z_i(t) - z_{i+1}(t) \right)
\]

It can be written as follows:

\[
\begin{align*}
&\frac{C_x}{4C_x \Delta} (x_i - x_{i+1})^2 - \left( \frac{C_x}{4C_x \Delta} \right)^{0.5} (x_i - x_{i+1}) \\
&\Delta C_y^{0.5} (y_i - y_{i+1})^2 + C_y d^2 \left( x_i - x_{i+1} \right)^2 \\
&\frac{d}{d} \left( y_i(t) - y_{i+1}(t) \right) - \left( y_i(t) - y_{i+1}(t) \right) + \left( x_i(t) - x_{i+1}(t) \right) + \left( z_i(t) - z_{i+1}(t) \right) \\
&\left( 1 - \Delta \right)^{0.5} (y_i - y_{i+1})^2, \quad \Delta \in (0, 1)
\end{align*}
\]

Taking this into account one can rewrite (10) as:

\[
\dot{V} \leq \sum_{i=1}^{n-1} \left( -C_x(x_i - x_{i+1})^2 \left( \frac{x_i^2}{2} + \frac{x_{i+1}^2}{2} \right) + b(x_i + x_{i+1}) + \gamma y_i \right) + \\
C_y d (x_i - x_{i+1})^2 - C_y e (z_i - z_{i+1})^2 \tag{11}
\]

Let \( \frac{Cy}{C_x} = \frac{2(1 - \Delta)}{d^2} \).

Then

\[
\dot{V} \leq \sum_{i=1}^{n-1} \left( -C_x (x_i - x_{i+1})^2 \left( \frac{x_i^2}{2} + \frac{x_{i+1}^2}{2} \right) + b(x_i + x_{i+1}) + \gamma y_i \right) + \\
\gamma n - \frac{d^2}{8(1 - \Delta)} - C_y e (z_i - z_{i+1})^2 \tag{12}
\]

Let \( \gamma > \left( \frac{d^2}{8(1 - \Delta)} + b^2 \right)/n \). Then we have that

\[
\lim_{t \to \infty} \int_0^t \left( x_i(r) - x_{i+1}(r) \right)^2 dr < \infty \\
\lim_{t \to \infty} \int_0^t \left( z_i(r) - z_{i+1}(r) \right)^2 dr < \infty
\]

Furthermore, the system trajectories are bounded and the system right-hand side is continuous. Hence according to Barbalat’s lemma we can conclude that

\[
\lim_{t \to \infty} (x_i(t) - x_j(t)) = 0, \quad \lim_{t \to \infty} (z_i(t) - z_j(t)) = 0
\]

To show that differences \( y_i(t) - y_j(t) \) tend to zero as \( t \to 0 \) it is sufficient to notice that

\[
-d \left( y_i(t) - y_{i+1}(t) \right) - \left( y_i(t) - y_{i+1}(t) \right) + \left( x_i(t) - x_{i+1}(t) \right) + \left( z_i(t) - z_{i+1}(t) \right) \to 0 \quad \text{as} \quad t \to \infty
\]

The lowest admissible bound \( \gamma \) for \( \gamma (d, b, n, \Delta) \) with respect to \( \Delta \) can be defined by \( \gamma = 0.5d^2 + b^2/n \).

The proposition is proven.

The proposition provides bounds for \( \gamma \) which are independent of the initial conditions, the excitation parameter \( I \) in the model and the parameter \( c \) which regulates the dynamics of the spikes.

Furthermore, it is necessary to point out that the value for \( \gamma \) is decreasing with the rate of \( O(1/n) \) if the number of interconnected oscillators is increasing. This observation is similar to the results in (Pogromsky, 1998) except, however, the fact that the bound for \( \gamma \) in our case is defined explicitly (and only) by the parameters of the model itself.

One question, however, is still open: whether the bound for \( \gamma \) can be lowered? In order to answer this question we should notice that the results formulated in Proposition 7 are global and are independent of the initial conditions. Therefore it is natural to expect that there is room for further
improvements if we assume that only initial conditions in a neighborhood of the synchronization manifold are allowed. The analysis for this case is given in the next section.

5. LOCAL UPPER BOUND

The main idea behind our approach is first to define a neighborhood of the synchronization manifold and then design a Lyapunov candidate with non-positive derivative in the same domain of the system state space. The estimates of the coupling parameter are expected to depend on the size of the domain of admissible initial conditions. The results of this local analysis are formulated in Proposition 8.

Proposition 8. Let system (2) be given, coupling function \(\gamma(x)\) be linear, uniform and preserving with respect to the manifold \(x_1 = x_2 = \cdots = x_n\). Let, in addition,

\[
(x_i(0) - x_j(0))^2 + \frac{3}{4d^2 + 2[b - \frac{3}{4}d]d} (y_i(0) - y_j(0))^2 + \frac{1}{\delta^2} (z_i(0) - z_j(0))^2 < \delta^2
\]

for some \(\delta > 0\). Furthermore let \(\gamma > \gamma_0\), where

\[
\gamma_0 = \frac{1}{n} \left( \left[ \frac{b - \frac{3}{4}d}{2d^2} + \frac{\frac{1}{4}}{d^2 + 2[b - \frac{3}{4}d]d} \right] \delta^2 + \left[ (b - \frac{3}{4}d)^2 \right] B_4 \delta \right)
\]

Then

\[
\lim_{t \to \infty} x_i(t) - x_j(t) = 0 \\
\lim_{t \to \infty} y_i(t) - y_j(t) = 0 \\
\lim_{t \to \infty} z_i(t) - z_j(t) = 0
\]

Proof of Proposition 8. To prove the proposition consider the Lyapunov candidate given by equation (9). Its time-derivative is defined by (10). Rewrite it as follows:

\[
\dot{V} = \sum_{i=1}^{n-1} -C_x(x_i - x_{i+1})^2 \left( x_i^2 + x_{i+1}^2 + x_i x_{i+1} \right) - b(x_i + x_{i+1}) - \left( \frac{C_x}{C_y} - d(x_i + x_{i+1}) \right)^2 \frac{C_y}{4C_x} - C_y \left( y_i - y_{i+1} + \frac{C_x}{C_y} - d(x_i + x_{i+1}) \right) \times \left( x_i - x_{i+1} \right)^2 - C_x \left( z_i - z_{i+1} \right)^2
\]

Moreover, assume that \(|x_i - x_{i+1}| \leq \delta\). This automatically implies that \(x_{i+1}(t) = x_i(t) + \mu(t)\), where \(|\mu(t)| \leq \delta\). Denote \(\alpha = \frac{\delta}{\delta^2}\), then:

\[
\dot{V} \leq \sum_{i=1}^{n-1} -C_x(x_i - x_{i+1})^2 \left( \left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - C_x (x_i - x_{i+1})^2 \left( \left( \frac{1}{4d^2} \right) \mu^2(t) - \left( \frac{3}{4d^2} \right) x_i \mu(t) - C_y \left( y_i - y_{i+1} \right) \right) - \left( \frac{C_x}{C_y} - d(x_i + x_{i+1}) \right) \left( x_i - x_{i+1} \right)^2
\]

Consider term \(1 - \frac{d^2}{4d}\) \(\mu^2(t) - \left( \frac{1}{3} \right) \mu(t) + \left( \frac{3}{4d^2} \right) x_i \mu(t)\) in (14). Given that the solutions of the system are bounded \(|\mu(t)| < \delta\) and \(|\mu(t)| < \delta\), we can estimate it as follows:

\[
\dot{V} \leq \sum_{i=1}^{n-1} -C_x(x_i - x_{i+1})^2 \left( \left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - \left( \frac{1}{4d^2} \right) \delta^2 - \left( \left( \frac{1}{4d^2} \right) \mu^2(t) + \left( \frac{3}{4d^2} \right) x_i \mu(t) \right) B_4 \delta
\]

Taking this estimate and inequality (14) into account we can derive the following:

\[
\dot{V} \leq \sum_{i=1}^{n-1} \left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - \left( \frac{1}{4d^2} \right) \delta^2 - \left( \frac{3}{4d^2} \right) x_i \mu(t) \right) B_4 \delta
\]

It is clear that \(\dot{V}\) will be nonpositive as long as

\[
\left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - \left( \frac{3}{4d^2} \right) x_i \mu(t) \right) B_4 \delta
\]

is strictly positive. Consider

\[
\left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - \left( \frac{3}{4d^2} \right) x_i \mu(t) \right) B_4 \delta
\]

in (15) and let \(3 - \frac{C_x}{C_y} > 0\). Then this term can be written as

\[
\left( \frac{3}{4d^2} \right) x^2 - 2(b - \frac{3}{4}d) x_i - \frac{\alpha}{4} + \gamma \right) - \left( \frac{3}{4d^2} \right) x_i \mu(t) \right) B_4 \delta
\]
Notice that the minimal value of 
\[ \frac{(b - \frac{1}{2}d)^2 \alpha}{3\alpha - \frac{1}{3}d^2} + \frac{\alpha}{4} \]
for \( \alpha \in (\frac{d^2}{3}, \infty) \) is equal to 
\[ \left( \frac{3}{2} \right)^2 + \frac{2|b - \frac{1}{2}d|d}{3} \]
for 
\[ \alpha = \frac{d^2 + 2|b - \frac{1}{2}d|d}{3} \]
Hence for any 
\[ \gamma > 0 \left( \left( \frac{b - \frac{1}{2}d}{2|d|} + \frac{1}{4} \right)^2 + \frac{2|b - \frac{1}{2}d|d|}{3} + \| \right| - \frac{3d^2}{4d^2 + 2|d||b - \frac{1}{2}d|} \right) \]
we get that \( \dot{V} \leq 0 \). This fact in turn implies that function \( V(\cdot) \) is not increasing as soon as \( |x_i - x_j| < \delta \). The last inequality can be satisfied by the choice of initial conditions as follows:
\[ C_p(x_i(0) - x_j(0))^2 + C_p(y_i(0) - y_j(0))^2 + C_p(x_i(0) - x_j(0))^2 < C_\delta \delta^2. \]
Taking into account that \( C_p/C_\Delta = 1/\alpha \) and that \( C_\delta/C_\Delta = 1/(\varepsilon) \) we can rewrite inequality (16) as
\[ (x_i(0) - x_j(0))^2 + \frac{3}{d^2 + 2|d||b - \frac{1}{2}d|} (y_i(0) - y_j(0))^2 + \left( \frac{1}{\varepsilon^2} (z_i(0) - z_j(0))^2 < \delta^2 \right. \]
The rest of the proof is analogous to that of Proposition 7 and follows explicitly from Barbalat’s lemma. The proposition is proven.

It is desirable to notice that the estimate (13) for \( b = 3, d = 5 \) and \( n = 2 \) results in the limit \( \delta \to 0 \), in the following inequality: \( \gamma > 1.5 \). This estimate is much closer to the bounds for \( \gamma \) obtained in our computer experiments.

6. CONCLUSION

In this paper we formulated sufficient conditions for asymptotic synchronization in the ensembles of global, linear and uniformly coupled Hindmarsh and Rose oscillators. We have shown that local stability conditions result in significantly smaller value for the coupling parameter \( \gamma \) in comparison to that derived for the arbitrary initial conditions. One of the explicit applications of this result is in determining the domain for the values of the coupling parameters for which the on-off intermittency (N. Platt, 1993) effects are more likely to appear given the specific connections and set of parameters.

We have also shown that sufficient conditions for asymptotic synchronization of linear, preserving and uniformly coupled nodes can be derived as a function of the system parameters which is not explicitly dependent on the bounds of the system solutions. On the other hand, the coupling gain ensuring asymptotic synchronization is decreasing at least as \( O(1/n) \) when the number of interconnected systems is growing.

These results, however, are restricted to very specific classes of coupling functions. The more realistic cases would be diffusive and nonlinear couplings between the elements of the network. These are topics for our future study.

REFERENCES