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Thickness dependence of unidirectional spin-Hall magnetoresistance in metallic bilayers

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A nonlinear magnetoresistance—called unidirectional spin-Hall magnetoresistance—is recently experimentally discovered in metallic bilayers consisting of a heavy metal and a ferromagnetic metal. To study the fundamental mechanism of unidirectional spin-Hall magnetoresistance (USMR), both ferromagnetic and heavy metallic layer thickness dependence of the USMR are presented in a Pt/Co/AlOx trilayer at room temperature. To avoid ambiguities, second harmonic Hall measurements are used for separating spin-Hall and thermal contributions to the non-linear magnetoresistance. The experimental results are fitted by using a drift-diffusion theory, with parameters extracted from an analysis of longitudinal resistivity of the Co layer within the framework of the Fuchs-Sondheimer model. A good agreement with the theory is found, demonstrating that the USMR is governed by both the spin-Hall effect in the heavy metallic layer and the metallic diffusion process in the ferromagnetic layer. Published by AIP Publishing. https://doi.org/10.1063/1.5003725

In the field of Spintronics, a new way of spin control based on the spin-Hall effect (SHE) recently has attracted a great deal of attention. It originates from the spin-orbit (SO) interaction which converts a charge current into a net flow of spin angular momentum, exerting a SO torque on the magnetization. This leads to an energy-efficient way of writing information to magnetic memories by switching a magnetic entity via sending a charge current through a nearby nonmagnetic metal.1,2 Apart from writing, a possible way of reading the information to magnetic memories by switching a magnetic entity via sending a charge current through a nearby nonmagnetic metal.1,2 Apart from writing, a possible way of reading the

Very recently, a unidirectional contribution to magnetoresistance—called unidirectional spin-Hall magnetoresistance (USMR)—has been reported in a ferromagnetic/heavy metallic (FM/HM) bilayer structure.5–7 Being different from the ordinary SMR, the resistance changes by reversing the magnetization or the current direction, which could be potentially utilized for reading operation. Based on a drift-diffusion-relaxation theory,8 this nonlinear behavior is attributed to the dependence of electron mobility on spin-polarization, which is tuned by the spin-Hall effect induced spin accumulation. This spin accumulation is limited to a thin region at the FM/HM interface due to a finite spin diffusion length in both layers, leading to a non-trivial FM and HM thickness dependence of the USMR. So far, this particular dependence on thickness is not evidenced by any experiments. Thus, a systematic investigation of how USMR depends on the layer thickness is urgently needed not only for a better understanding of the origin of USMR but also for the enhancement of USMR in practical applications.

In this paper, we present the FM and HM layer thickness dependence of USMR in Pt/Co/AlOx trilayers at room temperature. The experimental results are fitted by using the aforementioned drift-diffusion-relaxation theory, with parameters extracted from an analysis of the longitudinal resistivity of the Co layer within the framework of the Fuchs-Sondheimer model.9 Furthermore, second harmonic Hall measurements enable us to disentangle spin-Hall and thermal gradient contributions to the non-linear magnetoresistance, allowing for a more precise fitting. Good agreement with the theory is found, demonstrating that the USMR depends on both the spin-Hall effect in the HM layer and the electron spin diffusion and relaxation in the FM layer.

For the measurement of magnetoresistance, the multi-layer structures are Pt(1–8 nm)/Co(4 nm)/AlOx(1.15 nm) (6 samples) and Pt(4 nm)/Co(1–50 nm)/AlOx(1.15 nm) (11 samples), where we vary either the Pt or Co thickness as indicated by the thickness range in the parentheses. These multilayers are then patterned in the form of a Hall bar, shown in Fig. 1(a), by using electron-beam lithography and lift-off. The length of the Hall bar is 100 μm, the lateral width is 5 μm, and the spacing between two Hall bars is 20 μm. The samples are deposited on Si/SiO2 substrates by DC magnetron sputtering. Pt was deposited at a rate of 0.08 nm/s, and Co was sputtered at a rate of 0.05 nm/s. After deposition, a 1.15 nm thick Al capping layer was finally deposited and further oxidized (by using plasma oxidation during 90 s at 1 × 10−1 mbar) on top of the Pt/Co stack, to prevent oxidation of the Co layer in air.

The magnetoresistance measurements presented in this work were performed at room temperature by using an AC current source with a current density of 1 × 107 A/cm2

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my fitted line with respect to the confirmed by a good agreement between the data and a measurement of magnetic domains during magnetization switching. In comparison with the angle dependent measurement, where is measured between the directions, plotted as a solid line in Fig. 1(b). Note that the transition for the PMA stack is significantly sharper compared to the angle dependent measurement of resistance between the directions. After that, reaches a minimum when the field is along the direction, and finally, it returns to its original value after a full rotation.

The measurement evidences a resistance contribution that depends on the sign of , and the difference of between the +y and −y directions is defined as the USMR. This is further confirmed by a good agreement between the data and a fitted line with respect to obtained by an anomalous Hall effect measurement, plotted as a solid line in Fig. 1(b). Note that the transition for the PMA stack is significantly sharper than that for the in-plane sample, since a high field is needed to pull the magnetization in the plane for a stack with PMA.

To further investigate the USMR, we have measured while sweeping the external magnetic field along transverse (y) and longitudinal (x) directions. Figure 1(c) shows that is constant as a function of y field and reverses sign upon sweeping the field from the direction to the −y direction. Two spikes are observed near zero field due to the formation of magnetic domains during magnetization switching. In contrast to a field in the y direction, no difference is measured between the x and −x directions, indicating that, as expected, the USMR only exists in the transverse direction. Compared with the angle dependent measurement, where one has to ensure that the field is strong enough to saturate the sample in the z direction, the field dependent measurement serves as a more efficient way of quantifying the USMR. Thus, in the following, USMR will be obtained by sweeping the field. As a further test, Fig. 1(d) shows the current dependence of USMR measured in this way, which is linear with the injected current density and converges to zero for decreasing current, since the spin accumulation at the interface scales with the current density.

To verify the role of the interfacial spin accumulation due to the SHE, we examined the dependence of the USMR on the thickness of the HM and FM layers. Figures 2(a) and 2(b) show the absolute change of second harmonic resistance measured at constant current density as a function of the Co and Pt thickness. Both curves exhibit qualitatively similar behavior: An initial sharp increase and a gradual decrease as the layer becomes thicker. Apart from the USMR, thermal effects could also contribute to the . Thus, exclusion of these thermal effects is required before further analysis, which will be discussed in the following part.

Figure 2(c) shows the temperature profile in a line along the thickness direction of the nanowire by using simulation software suites Comsol multiphysics 5.2. The maximum temperature is found at the top owing to the fact that the heat dissipation is faster through the bottom substrate than through the top ambient air. A temperature gradient in the z direction will generate an electric current in the same direction, which interacts with the FM layer through the anomalous Hall effect and generates an electric field (∝ΔT × Msat) in the x direction. This will cause a resistance change in the x direction and possesses the same symmetry as that of the USMR. In order to separate the thermal contribution from the USMR, we measure the second harmonic...
Hall resistance to quantify the thermal resistance\(^{11}\) (see Supplementary Material for details). As plotted in the same figure of USMR, see the red dots in Fig. 2(a), the thermal contributions are found to be increased with the thickness. The maximal thermal resistance is observed at a Co thickness of 50 nm, which accounts for about 50% of the USMR and the ratio is smaller for thinner Co. In the following part, the thermal resistance will be subtracted from the USMR to achieve an accurate analysis.

In order to compare the experimental measurement of USMR with the model, we first convert the absolute USMR into normalized USMR, i.e., USMR divided by normal longitudinal resistance. For this purpose, the longitudinal resistance \(R_{xx}\) is measured and plotted versus Pt and Co thickness in Figs. 3(a) and 3(b). The plot reveals that the resistance monotonically decreases with thicknesses. The solid line represents the fit which utilizes the Fuchs-Sondheimer approach to extend the conventional \(\tau^{-1}\) resistance model by considering the scattering at the two Co/Pt interfaces.\(^9\) The fit describes the experimental data well and gives the bulk resistivities of \(\rho_{\text{Pt}} = 37.5 \mu\Omega \text{cm}\) and \(\rho_{\text{Co}} = 31.1 \mu\Omega \text{cm}\), which are comparable to the values in the literature.\(^{12,13}\)

Next, we examine the dependence of USMR on the Pt and Co thickness in Pt/Co/AlO\(_x\) samples. As shown in Fig. 3(c), the normalized USMR is the largest for a Pt thickness of about 5 nm and is reduced for a thicker or thinner Pt layer. Its strong thickness dependence shows that USMR in the structures is mainly influenced by the SHE in Pt. USMR decreases when the Pt layer is thinner than the spin diffusion length due to the reduced spin current caused by back reflection at the interface. On the other hand, for a thicker Pt layer, USMR is also reduced by current shunting. A similar behavior is found for USMR upon varying the Co thickness, as shown in Fig. 3(d), although the maximal USMR is now reached at a Co thickness of 10 nm. Above all, the qualitative behavior matches the prediction of the drift-diffusion-relaxation model. In addition, we also make a quantitative comparison of the experimentally observed values with the model, which describes the USMR as

\[
\rho(E) - \rho(-E) = \frac{6\theta L_F L_H (p_\sigma - p_N) \sigma_F \text{tanh} \left( \frac{d_F}{L_E} \right) \text{tanh} \left( \frac{d_H}{2L_H} \right)}{\epsilon_F \left( d_I \sigma_F + d_H \sigma_H \right) \frac{L_H (1 - p_\sigma^2) \text{tanh} \left( \frac{d_F}{L_E} \right) \text{coth} \left( \frac{d_H}{L_H} \right) + 1 \right) },
\]

where \(d_F\) (\(d_H\)) is the thickness of the FM (HM), \(L_E\) (\(L_H\)) is the spin diffusion length of the FM (HM), \(\sigma_F\) (\(\sigma_H\)) is the conductivity of the FM (HM), \(\theta\) is the spin Hall angle of the HM, \(\epsilon\) is the electric field in FM, \(\epsilon_F\) is the Fermi energy, \(p_\sigma\) is the conductivity spin asymmetry, and \(p_N\) is the difference of density of states at Fermi energy. In our sample, the Fermi energy \(\epsilon_F = 5\) eV and spin asymmetry \(p_\sigma - p_N = 0.5\).\(^8\)

By fitting the thickness dependence of the normalized USMR to Eq. (1) (red line in Fig. 3), it can be seen that the specific behavior of the data is in line with the drift-diffusion-relaxation model based on the spin-Hall effect, and a spin Hall angle of 0.3 for Pt and spin diffusion lengths of 18 nm and 2.2 nm for Co and Pt, respectively, can be extracted. The spin diffusion lengths are similar to the literature values.\(^{14,15}\) This implies that the USMR in the Pt/Co system is governed by (1) spin-Hall effect in the Pt layer and (2) electron spin diffusion and relaxation in the Co layer. In a recent work\(^16\) which uses a similar structure (Py/Co), a unidirectional contribution was found in the first harmonic resistance by applying a high current density (\(\sim 10^8\) A/cm\(^2\)). In this experiment, magnon excitation, instead of an electronic diffusion-relaxation process, is claimed to attribute to this first harmonic USMR. It is also found that the magnon-induced USMR increases with increasing temperature. We do not intend to conduct a temperature dependence measurement here, due to the complex variation of all transport parameters with temperature\(^{17-19}\) (polarization, spin diffusion length, spin Hall angle, and conductivity) in the drift-diffusion-relaxation model [Eq. (1)], which will make it extremely difficult to draw pertinent conclusions. Moreover, in the magnon experiment,\(^16\) a much
higher current density is used compared to ours, and only the first harmonic in the resistance is addressed, which further complicates a meaningful comparison.

To disentangle electronic and magnonic contribution, a measurement of the temperature dependence of the USMR needs to be performed, which is beyond the scope of this paper. For the electron contribution described before [Eq. (1)], taking into account the temperature variation of all transport parameters (polarization, spin diffusion length, spin Hall angle, and conductivity) predicts that the USMR decreases with increasing temperature, whereas the magnon-induced USMR behaves oppositely.\textsuperscript{16}

The extracted room-temperature spin Hall angle in Pt appears to be higher than the value \(0.1\) measured in other work.\textsuperscript{20,21} We still think that this model captures the essential physics of the observed effect, although full quantitative agreement cannot be reached due to various reasons. One reason is the simplifications of the model by assuming spherical Fermi surfaces and constant density of states at the Fermi energy,\textsuperscript{8} which underestimates the magnitude of USMR. Moreover, the model\textsuperscript{8} includes only spin-dependent scattering in the bulk of the ferromagnetic layer. Like in the giant magnetoresistance effect, however, both bulk and interface scattering can contribute to the USMR.\textsuperscript{5} The underestimation would be more if the spin-mixing conductance is incorporated, since the Pt/Co interface is regarded as fully transparent in the model, i.e., the spin-mixing conductance is infinite. Finally, additional charge-spin conversion can take place at either the Pt/Co or Co/AlO\textsubscript{x} interface, which may lead to a larger spin-Hall effect.\textsuperscript{22,23}

In conclusion, USMR is observed in Pt/Co/AlO\textsubscript{x} systems and we have shown that the dependence of the USMR on the thickness of both the HM and FM layers agrees qualitatively with the theory based on the electron spin drift-diffusion-relaxation model. We believe that this result provides a better understanding of the physical origin of the USMR and is of importance for its possible applications in spintronic devices.

See supplementary material for the quantification of thermal contributions to the unidirectional spin-Hall magnetoresistance by measuring the second harmonic Hall resistance.

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\textsuperscript{22}V. P. Amin and M. D. Stiles, Phys. Rev. B \textbf{94}, 104419 (2016).