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Stabilizing Lead-Free All-Inorganic Tin Halide Perovskites by Ion Exchange

Junke Jiang,† Chidozie K. Onwudinanti,‡ Ross A. Hatton,§ Peter A. Bobbert,¶ and Shuxia Tao∥

†Center for Computational Energy Research, Department of Applied Physics, and ¶Molecular Materials and Nano Systems, Department of Applied Physics, Eindhoven University of Technology, P.O. Box 513, 5600 MB Eindhoven, The Netherlands
‡Center for Computational Energy Research, DIFFER—Dutch Institute for Fundamental Energy Research, De Zaale 20, 5612 AJ Eindhoven, The Netherlands
§Department of Chemistry, University of Warwick, CV4 7AL Coventry, U.K.

ABSTRACT: Because of its thermal stability, lead-free composition, and nearly ideal optical and electronic properties, the orthorhombic CsSnI3 perovskite is considered promising as a light absorber for lead-free all-inorganic perovskite solar cells. However, the susceptibility of this three-dimensional perovskite toward oxidation in air has limited the development of solar cells based on this material. Here, we report the findings of a computational study which identifies promoting Rb, CsSn(BrIx)x perovskites for solar cell applications, prepared by substituting cations (Rb for Cs) and anions (Br for I) in CsSnI3. We show the evolution of the material electronic structure as well as its thermal and structural stabilities upon gradual substitution. Importantly, we demonstrate how the unwanted yellow phase can be suppressed by substituting Br for I in CsSn(BrIx)x with x ≥ 1/3. We predict that substitution of Rb for Cs results in a highly homogeneous solid solution and therefore an improved film quality and applicability in solar cell devices.

INTRODUCTION

Organic−inorganic hybrid halide perovskite solar cells (PSCs) have attracted strong attention in the past few years and are becoming one of the most promising types of emerging thin-film solar cells.1−5 In less than a decade, the power conversion efficiency (PCE) of PSCs has increased from 3.8% in 2009 to 22.7% in 2017.6−8 Despite the high efficiency of PSCs, two challenges currently hinder their upscaling toward practical applications. One issue is the long-term instability of PSCs, which is mainly caused by the intrinsic thermal instability of hybrid perovskite materials.8−13 Encouragingly, it has been demonstrated recently that mixing the cations or replacing the organic cation with an inorganic cation can improve thermal stability and photostability (e.g., substituting FA for MA in MAPbI3, Rb for Cs in CsSnI3, and Cs for MA in MAPbI3; MA stands for CH3NH3 and FA stands for NH2CHNH2).14,15 The other concern is the well-documented toxicity of lead (Pb), which is particularly problematic because lead halide perovskites decompose into lead compounds that have significant solubility in water.16 Consequently, an intensive research effort focused on finding air-stable lead-free perovskites suitable as the light-harvesting semiconductor in PSCs is now underway.9,17−20

Among the various alternatives to lead, tin (Sn) is regarded as a promising substitute because Sn-based hybrid perovskites have been shown to exhibit outstanding electrical and optical properties, including high charge carrier mobilities, high absorption coefficients, and low exciton binding energies,21−23 Theoretical predictions by Even et al.24 and Chiarella et al.25 also confirmed the promising properties of Sn perovskites, such as suitable band gaps and favorable effective mass. However, Sn-based perovskites also have drawbacks, which have limited their application in efficient PSCs.1,23,26−29 The primary challenge is the susceptibility of tin toward oxidation from the +2 to the +4 oxidation state upon exposure to ambient air, which, in the case of CsSnI3, ultimately results in the formation of CsSnI5 whose relatively weak light absorption across the visible spectrum is undesirable for a photoabsorber.22,30−33 Consequently, to date, there has been much less research effort directed at the advancement of tin halide PSCs than their lead analogues, and their PCE has remained below 10%.27,28 Recently, a PCE as high as 9.0% in PSCs was achieved using single-crystalline FASnI3 made by mixing a small amount of two-dimensional (2D) Sn perovskites with three-dimensional (3D) FASnI3 in which the organic FA molecules are oriented randomly,34 an approach that promises further improvement.

As compared to hybrid organic−inorganic Sn perovskites, all-inorganic Sn perovskites could have the advantage of...
Marshall et al.22 achieved the highest PCE to date of 3.56%, used CsSnI3 to fabricate a Schottky contact solar cell, which variation of perovskites are scarce.7,48 A recent study35,36 investigated the mixing of cations and anions in all-inorganic Sn-based semiconductors with band gaps in the range 1.3–2.0 eV. Importantly, our results indicate that substitution of Br for I in CsSnI3 can prevent the unwanted γ-to-Y phase transition, evidenced by the favorable formation energies of the γ phase over the Y phase. In addition, calculations of the free energy of mixing and the prediction of phase diagram demonstrate that further substitution of Rb for Cs in CsSnI3 can improve the mixing thermodynamics, which is expected to improve the film-forming properties. Our predicted trends in the thermodynamic stability and band gaps provide a guideline to develop more efficient and stable lead-free all-inorganic perovskites for PSCs.

**COMPUTATIONAL METHODS AND STRUCTURAL MODELS**

The initial structure optimizations are performed using DFT as implemented in the Vienna ab initio simulation package (VASP).55,56 The Perdew, Burke, and Ernzerhof (PBE) functional within the generalized gradient approximation is used.77 The outermost s, p, and d (in the case of Sn) electrons are treated as valence electrons, whose interactions with the remaining ions are modeled by pseudopotentials generated within the projector-augmented wave method.58,59

Figure 1 shows the crystal structures and cells used in the DFT calculations. Unit cells with 20 atoms (four ASnX3 units) are treated as valence electrons, whose interactions with the remaining ions are modeled by pseudopotentials generated within the projector-augmented wave method.58,59

![Figure 1. Top (a–c) and side (d–f) views of the cubic (α, Pm3m), orthorhombic (γ, Pnma), and yellow phase (Y, Pnma) of ASnX3 (A = Cs and Rb and X = Br and I).](image-url)
subsequent electronic structure calculations were performed using an efficient approximate quasiparticle DFT method, namely, the DFT-1/2 method. The DFT-1/2 method stems from Slater’s proposal of an approximation for the excitation energy, a transition-state method,60,61 to reduce the band gap inaccuracy by introducing a half-electron/half-hole occupation.62 The computational effort is the same as for standard DFT, with a straightforward inclusion of SOC when coupled with VASP. In this work, we extend the use of the DFT-1/2 method with the same settings (CUT values of $8, 6, 8, 8, 8, 8$, and $0.1$ meV and $8, 6, 8, 8, 8, 8$, respectively) are used to achieve an energy and force convergence of $0.1$ meV and $8, 6, 8, 8, 8, 8$, respectively. The $8, 6, 8, 8, 8, 8$ and $8, 6, 8, 8, 8, 8$ configurations (2, 4, and 2 for $8 = 1/4, 1/2, 3/4$, respectively) are used to achieve an energy and force convergence of $0.1$ meV and $8, 6, 8, 8, 8, 8$, respectively. The $8, 6, 8, 8, 8, 8$ and $8, 6, 8, 8, 8, 8$ configurations (2, 4, and 2 for $8 = 1/4, 1/2, 3/4$, respectively) of substituting Rb for Cs were considered. Owing to the large number of possible configurations of substituting Br for I (2, 11, 22, 223, 223, and 222 possible configurations for $8 = 1/6, 1/3, 1/2, 3/2, 5/6, 6$, respectively), we have considered only two possible configurations for each concentration of Br, namely, the two extreme cases with most negative and least negative $8$. From Figure S2 and Table S1, the formation energy of configuration 3 is the most negative, whereas that of configuration 7 is the least negative. This indicates that the Br ions tend to sit as close as possible to each other and to form as many bonds as possible with Sn ions. We use this strategy to select two extreme configurations for all other Br–I alloys considered in this work.

**RESULTS AND DISCUSSION**

Before studying the mixing of A cations and X anions in $8_{13}$, we first performed calculations for the four pure compounds: CsSnI$_3$, CsSnBr$_3$, RbSnI$_3$, and RbSnBr$_3$. The calculated lattice parameters of orthorhombic ($8$) $8_{13}$ are shown in Table 1. Those of other polymorphs including cubic $8$, tetragonal $8$, and $8$ phase structures are listed in Table S2. The optimized lattice parameters are in good agreement with experiments, with a slight overestimation of lattice parameters by about $1\%$, and with other theoretical results (differences within $0.1\%$).$7,38,49,65,66$ It should be noted here that the predicted lattice parameters of $\alpha$-CsSnI$_3$ and $\gamma$-CsSnI$_3$ in our previous work are smaller because of the use of LDA, which slightly underestimates the lattice parameters.$62$ In this work, PBE is used, resulting in a slight overestimation of lattice parameters. Consequently, the predicted band gap of $\gamma$-CsSnI$_3$ (1.36 eV) in this work (will be discussed in the next paragraph) is slightly higher compared to that of previous work (1.34 eV).$62$

<table>
<thead>
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<th>material</th>
<th>lattice constants (this work)</th>
<th>lattice constants (experimental)</th>
<th>lattice constants (other theoretical work)</th>
<th>$E_g$ (DFT-1/2 + SOC)</th>
<th>$E_g$ (DFT-1/2)</th>
<th>$E_g$ (DFT-1/2 + SOC)</th>
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</thead>
<tbody>
<tr>
<td>$\gamma$-CsSnI$_3$</td>
<td>$8.99, 12.52, 8.63$</td>
<td>$8.69, 12.38, 8.64$</td>
<td>$8.94, 12.52, 8.69$</td>
<td>$1.36$</td>
<td>$1.27$</td>
<td>$1.34$</td>
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<tr>
<td>$\gamma$-RbSnI$_3$</td>
<td>$8.91, 12.28, 8.47$</td>
<td>$8.93, 12.28, 8.47$</td>
<td>$1.55$</td>
<td>$1.13$</td>
<td>$1.83$</td>
<td></td>
</tr>
<tr>
<td>$\gamma$-CsSnBr$_3$</td>
<td>$8.36, 11.79, 8.22$</td>
<td>$8.38, 11.55, 7.98$</td>
<td>$1.72$</td>
<td>$2.01$</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

$\alpha$-CsSnI$_3$ and $\gamma$-CsSnI$_3$ are 1.36 and 1.72 eV, respectively, in excellent agreement with reported experimental measurements$69$ (1.27 and 1.75 eV) and $GW_0$ calculations$68$ (1.34 and 1.83 eV). There are no experimental reports known to us of the band gap of either RbSnI$_3$ or RbSnBr$_3$. Only theoretical results from HSE06 for $\gamma$-RbSnI$_3$ and PBE for $\alpha$-RbSnBr$_3$ are found to be 1.41 and 0.57 eV, respectively.$69$ Our predicted band gap for $\gamma$-RbSnI$_3$ is 1.55 eV. Substituting Br for I in $\gamma$-RbSnI$_3$ further increases the band gap to 2.01 eV.

It is worth noting that although CsSnBr$_3$ is reported to have the $\alpha$ structure at room temperature, the actual atomic arrangement at finite temperature (due to the dynamic
disorder of the ions in the lattice)\(^{38}\) resembles that of the \(\gamma\) phase. Consequently, it is not surprising that predicted band gaps using \(\alpha\) structures are always significantly smaller than those measured experimentally.\(^{7,24,71,72}\) Therefore, in this work, we always report band gaps calculated using \(\gamma\) phases. \(\text{RbSnI}_3\) has been reported to exist in a nonperovskite 2D \(y\) phase structure owing to the small cationic size of \(\text{Rb}^+\).\(^{38,66}\)

Nevertheless, for comparison with the alloys \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\), the band gaps of \(\text{RbSnI}_3\) in a 3D \(\gamma\) phase are also predicted. In addition, all band gaps of \(\alpha\) structures are also provided in the Supporting Information in Table S2 and Figure S3 for comparison. The calculated effective masses of the electrons and holes at the \(G\) point for \(\gamma\)-\(\text{ASnI}_3\) are given in Table S3.

Figure 2 shows the computed band gaps for the \(\gamma\) phases of \(\text{ASn(Br}_{1-x}\text{I}_x)\) and \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) perovskites, whereas the band gaps of the other structures are shown in Figure S4. In general, the band gap increases with increasing percentage of \(\text{Br}\) in \(\text{ASn(Br}_{1-x}\text{I}_x)\) and \(\text{Rb}\) in \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\). From Figure 2a and Table S4, generally, the band gaps change because of the variations in both volume and lattice distortion. However, the changes in cell volume have more pronounced effects on the band gaps than the changes in lattice distortion, that is, octahedral tilting. The reduction of the cell volume is responsible for the widening of the band gap in \(\text{ASn(Br}_{1-x}\text{I}_x)\) or \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) solid solutions with an increased \(\text{Br}\) or \(\text{Rb}\) percentage. For the band gap variations with the same \(\text{Br}\) or \(\text{Rb}\) percentage in \(\text{ASn(Br}_{1-x}\text{I}_x)\) or \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) solid solutions, there is no certain relationship found between the degree of lattice distortion (i.e., the degree of octahedral tilting, which is the tilting angle difference \(\Delta\theta\)) and band gaps.

It is well-known that for single-junction and multijunction solar cells, the Shockley–Queisser limit sets the optimal band gap range of 0.9–1.6 and 1.6–2.0 eV, respectively, for achieving a maximum PCE.\(^{25,73}\) The band gaps of \(\text{CsSn(Br}_{1-x}\text{I}_x)\) are almost completely in the optimal range (1.30–1.55 eV) for single-junction PSCs. When substituting \(\text{Rb}\) for \(\text{Cs}\) in \(\text{CsSn(Br}_{1-x}\text{I}_x)\), the band gaps of \(\text{RbSn(Br}_{1-x}\text{I}_x)\) increase by 0.2–0.3 eV as compared to their \(\text{Cs}\) counterparts, making \(\text{RbSn(Br}_{1-x}\text{I}_x)\) (\(x > 1/3\)) ideal as a wide-band gap material for tandem solar cells in conjunction with narrow-band gap semiconductors such as \(\text{Si}\) or \(\text{CsSn(Br}_{1-x}\text{I}_x)\). The changes in band gap when mixing \(\text{Rb}\) and \(\text{Cs}\) cations are much smaller than those when mixing \(\text{I}\) and \(\text{Br}\). This is true for all values \(0 \leq y \leq 1\) and also for different structures with a fixed \(y\). The band gaps of \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{Br}_{\text{y}}\text{SnI}_{3-(1-\text{y})}\) (1.71–2.01 eV) and \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) (1.36–1.55 eV) are in the ideal range for tandem and single-junction solar cells, respectively.

Because of the different crystal structure and electronic properties of the \(\gamma\) phase (i.e., a 2D structure and an indirect band gap of 2.6 eV), the unwanted phase transition from \(\gamma\) to \(\text{Y}\) can considerably decrease the efficiency of a solar cell.\(^{14,22,75–77}\) In addition, the \(\text{Y}\) phase will spontaneously react with \(\text{O}_{2}\) when exposed to air, resulting in \(\text{Cs}_2\text{SnI}_6\) with a face-centered cubic structure.\(^{30,32}\) Therefore, we focus here on the evolution of the stability of the \(\gamma\) and \(\text{Y}\) phases upon gradual substitution of \(\text{Br}\) for \(\text{I}\) and \(\text{Rb}\) for \(\text{Cs}\). The results of our calculations for other structures are given in Figure S5.

The formation energy of \(\text{ASn}_x\) is defined as \(\Delta H = E_{\text{ASn}_x} - E_{\text{AX}} - E_{\text{SnX}_2}\) where \(E_{\text{ASn}_x}\), \(E_{\text{AX}}\), and \(E_{\text{SnX}_2}\) are the total energies of \(\text{ASn}_x\), \(\text{AX}\), and \(\text{SnX}_2\), respectively. Here, a negative value of \(\Delta H\) represents favorable formation of \(\text{ASn}_x\) perovskites. The formation energy of \(\text{ASnX}_3\) in the \(\gamma\) phase is shown in Figure 3. For all values \(0 \leq y \leq 1\), the lowest energy \(\text{ASnX}_3\) solid solutions are located in the ideal range for tandem and single-junction solar cells, respectively.

![Figure 2](https://example.com/f2.png)

**Figure 2.** Calculated band gaps of (a) \(\text{ASn(Br}_{1-x}\text{I}_x)\) and (b) \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) perovskites. The dashed lines are guides to the eye.

![Figure 3](https://example.com/f3.png)

**Figure 3.** Formation energy (\(\Delta H\)) of (a) \(\text{ASn(Br}_{1-x}\text{I}_x)\) and (b) \(\text{Rb}_{\text{x}}\text{Cs}_{1-\text{x}}\text{SnI}_3\) perovskites for the \(\gamma\) and \(\text{Y}\) phases. Because of the large number of possible configurations for each substitution concentration \(x\) of \(\text{Br}\) in \(\text{ASn(Br}_{1-x}\text{I}_x)\), we only show in (a) the results for the two configurations with the most negative and least negative \(\Delta H\).
more negative $\Delta H$, the more stable the corresponding structure. It can be clearly seen in Figure 3 that all perovskites considered exhibit good thermal stability, with large negative $\Delta H$ values. Figure 3 also shows the effect of ion mixing on the stability of the $\gamma$ phase with respect to the Y phase.

For CsSnI$_3$, the formation energies of the $\gamma$ and Y phases are the same. However, with the increase of Rb concentration, the structural instability of Rb$_x$Cs$_{1-x}$SnI$_3$ becomes increasingly pronounced. Consequently, upon exposure to air, the rate at which the perovskite oxidizes is predicated to increase with an increasing Rb concentration.$^7$ The substitution of Rb for Cs seems to facilitate the formation of the Y phase, as evidenced by the more negative formation energy of the Y phase than the $\gamma$ phase. It should be mentioned that the formation energies are both negative for RbSnI$_3$ in both $\gamma$ and Y phases. However, the formation energy of the Y phase is relatively more negative, indicating that the Y phase is more stable than the $\gamma$ phase. Indeed, the instability of the $\gamma$ phase is in agreement with the experimental observation of RbSnI$_3$ only existing in a 2D yellow phase.$^7$ On the contrary, for CsSnBr$_3$, the formation energy of the Y phase is much less negative than that of the $\gamma$ phase (by 0.06 eV), indicating that the $\gamma$ phase is more stable than the Y phase. The substitution of Rb for Cs results in a slight decrease (to 0.04 eV) in the energy difference between the two phases, with the $\gamma$ phase still being favored.

The formation energies of mixing Cs and Rb in Rb$_x$Cs$_{1-x}$SnI$_3$ or Rb$_x$Cs$_{1-x}$SnBr$_3$ follow a perfect linear relation (Figure 3b), indicating favorable mixing thermodynamics. However, substitution of Br for I (Figure 3a) shows an unusual trend as a function of $x$: the curves show first a decrease and then an increase, with a valley point at $x = 1/3$ in both CsSn(Br$_{1-x}$I$_x$)$_3$ and RbSn(Br$_{1-x}$I$_x$)$_3$. When $x < 1/3$, the most negative $\Delta H$ of the $\gamma$ phase for each concentration is relatively more positive than or nearly equal to the most negative $\Delta H$ of the Y phase, which indicates that the Y phase is favored over the $\gamma$ phase. When $x = 1/3$, the most negative $\Delta H$ of the $\gamma$ phase is clearly more negative than the most negative $\Delta H$ of the Y phase, whereas the least negative $\Delta H$ of the $\gamma$ phase is almost equal to the most negative $\Delta H$ of the Y phase. When $x > 1/3$, all $\Delta H$ of the $\gamma$ phase for each concentration are more negative than those of the Y phase, which means that the $\gamma$ phase is stabilized. For Rb$_x$Cs$_{1-x}$SnI$_3$ or Rb$_x$Cs$_{1-x}$SnBr$_3$, mixing Cs and Rb does not change the stability of the $\gamma$ phase with respect to the Y phase (Figure 3b). For Rb$_x$Cs$_{1-x}$SnI$_3$, the Y phase is always favored when mixing Rb and Cs, whereas for Rb$_x$Cs$_{1-x}$SnBr$_3$, the opposite is true. We conclude that the addition of Br to Rb$_x$Cs$_{1-x}$SnI$_3$ tends to stabilize the favorable $\gamma$ phase and suppress the transformation to the Y phase. The critical Br concentration is about one-third. This prediction calls for an experimental validation.

As shown in Figure 3, very different trends are observed for mixing of cations (Rb and Cs) and anions (I and Br) in $\gamma$-CsSnI$_3$. For a deeper insight, we have investigated the different mixing thermodynamics by calculating the Helmholtz free energy of mixing. Details of the calculations can be found in the Computational Methods and Structural Models section. Results of these calculations are shown in Figure S6. On the basis of the Helmholtz free energies, we plot the phase diagram for $\gamma$-ASnX$_3$ by using the GQCA$^{65}$ code, as shown in Figure 4.

For CsSn(Br$_{1-x}$I$_x$)$_3$, the critical temperature is 291 K (see Figure 4a), indicating that the mixing of anions (I and Br) is favorable at room temperature (300 K). However, for RbSn(Br$_{1-x}$I$_x$)$_3$ at 300 K, a miscibility gap is found in the composition region between $x_1 = 0.33$ and $x_2 = 0.70$ (see Figure 4b). The pure compounds RbSnI$_3$ and RbSnBr$_3$ are not miscible inside the miscibility gap under equilibrium conditions, leading to the formation of two phases with Br concentrations $x_1$ and $x_2$. Meanwhile, the alloy has spinodal points at the compositions $x_1^{\prime} = 0.40$ and $x_2^{\prime} = 0.62$ at room temperature. Thus, in the intervals $x_1 < x < x_1^{\prime}$ and $x_2^{\prime} < x < x_2$, a metastable phase can occur, showing small fluctuations in composition. The predicted critical temperature (the temperature above which the solid solution is stable for any
composition) is 312 K, which is significantly lower than the critical temperature of 343 K predicted for the MAPb(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} perovskite.\textsuperscript{63} This indicates that, although mixing of Br and I is not favored slightly below (for CsSn(BrI\textsubscript{1−x}I\textsubscript{x})) or around (for RbSn(BrI\textsubscript{1−x}I\textsubscript{x})) room temperature, the phase segregation in these alloys is less significant than that in MAPb(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} perovskites.

A uniform mixture can be synthesized either through control of the deposition kinetics or by annealing above the critical temperature. The uniform mixture tends to segregate below the critical temperature, but this segregation is a very slow process.\textsuperscript{63} The inclusion of smaller cations often provides an improvement, overcoming kinetic barriers and changing the local critical temperature. For example, smaller cations such as Cs and Rb were introduced in (FA/MA)PbI\textsubscript{3} and were shown to have a positive effect on the structural and photostability of state-of-the-art PSCs.\textsuperscript{8,10,13,78} Indeed, we predict that mixing of Rb and Cs in Rb\textsubscript{1−x}Cs\textsubscript{x}SnI\textsubscript{3} is very favorable at room temperature. For Rb\textsubscript{1−x}Cs\textsubscript{x}SnBr\textsubscript{3} and Rb\textsubscript{1−x}Cs\textsubscript{x}SnBr\textsubscript{3}, the phase diagrams show that mixing of cations (Rb and Cs) is favorable at temperatures above 118 and 137 K, respectively (see Figure 4c,d). Our prediction of the critical miscibility temperature of 137 K of Rb\textsubscript{1−x}Cs\textsubscript{x}SnBr\textsubscript{3} is in good agreement with the result of 140 K calculated by Jung et al.\textsuperscript{7}

The slight difference of the predicted critical temperature could be caused by the small variation in energies per cell due to the differences in computational settings (energy cutoff value, k-point grid, version, and implementation of VASP codes) in DFT calculations. The critical temperatures of mixing of Rb and Cs in Rb\textsubscript{1−x}Cs\textsubscript{x}SnX\textsubscript{3} are much lower than those of mixing of Br and I in ASn(BrI\textsubscript{1−x}I\textsubscript{x}). Therefore, additional mixing of Cs and Rb in ASn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} is predicted to bring down the critical temperature for mixing of Br and I below room temperature, suppressing phase segregation and resulting in better material quality for PV applications.

### Conclusions

In summary, the effects of cation (Cs and Rb) and anion (I and Br) mixing in all-inorganic tin halide perovskites have been investigated with DFT-based calculations. Using standard DFT for structure optimization and the DFT-1/2 method with SOC for band structure calculations, we studied the evolution of the structural, thermodynamic, and electronic properties as a function of the extent of substitution of Rb for Cs and Br for I. We predict that Cs\textsubscript{x}Rb\textsubscript{1−x}Sn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} perovskites have direct band gaps in the range of 1.3–2.0 eV. The alloys with high I and Cs concentrations are well suited for highly efficient single-junction PSCs, whereas those with high Br and Rb concentrations are suitable as wide-band gap materials for tandem PSCs. Importantly, we found that substitution of Br for I can suppress the unwanted γ-to-Y phase transition. The critical concentration for stabilization of the γ phase with respect to the Y phase in Cs\textsubscript{x}Rb\textsubscript{1−x}Sn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} is \(x = 1/3\). Furthermore, phase diagrams based on the free energy of mixing show that a solid solution of Br and I is thermodynamically possible around and slightly above room temperature for CsSn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} and RbSn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3}, respectively. Finally, substitution of Rb for Cs to ASn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} is predicted to decrease the critical temperature to well below room temperature, enabling the formation of highly homogeneous solid solutions for improved solar cell performance. Our predictions regarding the stabilization of the γ phase and the use of five elements in Rb\textsubscript{1−x}Cs\textsubscript{x}Sn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} as an efficient and stable light absorber for PSCs call for experimental exploration.

### Associated Content

#### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.8b04013.

Physical insights of why Sn p and halide I or Br p orbitals are both half-ionized; illustration of the selection strategy of ASn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} to reduce computational effort; summary of lattice constants and band gaps (\(E_g\)) of polymorphs of ASnX\textsubscript{3} (CsSn\textsubscript{x}, RbSn\textsubscript{x}, CsSnBr\textsubscript{y} and RbSnBr\textsubscript{y}); band structures for CsSn\textsubscript{x}, CsSnBr\textsubscript{y}, RbSn\textsubscript{x}, and RbSn\textsubscript{y} with the γ phase and the α phase; effective masses for γ-CsSn\textsubscript{x}, γ-CsSnBr\textsubscript{y}, γ-RbSn\textsubscript{x}, and γ-RbSn\textsubscript{y}; average tilting angles and cell volumes of γ-CsSn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3}, γ-RbSn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3}, γ-RbCsSnI\textsubscript{x}, and γ-RbCsSnBr\textsubscript{y} band gaps and formation energies of ASn(BrI\textsubscript{1−x}I\textsubscript{x})\textsubscript{3} and Rb\textsubscript{1−x}Cs\textsubscript{x}SnX\textsubscript{3} perovskites, including the α phase; and Helmholtz free energy of mixing of α and γ polymorphs (PDF)

### Author Information

**Corresponding Author**

*E-mail: s.x.tao@tue.nl.*

**ORCID**

Junke Jiang: 0000-0003-2962-766X

Ross A. Hatton: 0000-0002-8851-1280

Shuxia Tao: 0000-0002-3658-8497

**Notes**

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### References

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