Plasma Catalysis: Distinguishing between Thermal and Chemical Effects

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Abstract: The goal of this study is to develop a method to distinguish between plasma chemistry and thermal effects in a Dielectric Barrier Discharge nonequilibrium plasma containing a packed bed of porous particles. Decomposition of CaCO₃ in Ar plasma is used as a model reaction and CaCO₃ samples were prepared with different external surface area, via the particle size, as well as with different internal surface area, via pore morphology. Also, the effect of the CO₂ in gas phase on the formation of products during plasma enhanced decomposition is measured. The internal surface area is not exposed to plasma and relates to thermal effect only, whereas both plasma and thermal effects occur at the external surface area. Decomposition rates were in our case found to be influenced by internal surface changes only and thermal decomposition is concluded to dominate. This is further supported by the slow response in the CO₂ concentration at a timescale of typically 1 minute upon changes in discharge power. The thermal effect is estimated based on the kinetics of the CaCO₃ decomposition, resulting in a temperature increase within 80 °C for plasma power from 0 to 6 W. In contrast, CO₂ dissociation to CO and O₂ is controlled by plasma chemistry as this reaction is thermodynamically impossible without plasma, in agreement with fast response within a few seconds of the CO concentration when changing plasma power. CO forms exclusively via consecutive dissociation of CO₂ in the gas phase and not directly from CaCO₃. In ongoing work, this methodology is used to distinguish between thermal effects and plasma–chemical effects in more reactive plasma, containing, e.g., H₂.

Keywords: nonequilibrium plasma; plasma catalysis; gas temperature; calcium carbonate decomposition

1. Introduction

Plasma catalysis is receiving more and more attention in the last few years, since the specific interactions between plasma and catalyst surface may lead to synergistic effects [1–4]. One of the earliest plasma catalytic applications is the abatement of volatile organic compounds (VOC) [5,6], while in the last decade research has been focused more on CO₂ conversion [7–9], conversion of hydrocarbons via reforming, and coupling [10–12], as well as activation of N₂ [13,14]. Reforming of hydrocarbons is an example of an endothermic reaction, where plasma catalysis holds promise because of activation of hydrocarbons at low temperature, but also because electrical energy would be used to generate the required heat. Methane coupling and CO₂ dissociation are examples of thermodynamically hill-up reactions, which are clearly more challenging.

Nonequilibrium plasma, e.g., microwave of Dielectric Barrier Discharge (DBD) plasma, is especially attractive because it operates at relatively low temperatures [15–17]. Consequently, catalyst sintering is prevented. Moreover, low temperatures are a necessity to enable catalysis in the first place,
facilitating initial adsorption that decreases entropy. Starting and stopping plasma reactors is much faster than usual thermal reactors, which is an advantage when fast capacity changes are required, e.g., in connection with intermittent energy supply and storage.

DBD plasma is frequently used for studying plasma catalytic conversion. The high AC voltages applied at relatively low frequency (50 to $10^5$ Hz) produces a nonequilibrium plasma with very high electron temperatures (1–10 eV equal to $10^4$–$10^5$ K), rather high vibrational temperatures ($10^3$ K), and rather low rotational and translational temperatures in the plasma zone, typically in the order or smaller than 100 K [6,15–17]. The fact that energy would be directed directly to bond breaking, without the need to heat up the gas mixture completely, is very attractive as heat exchangers to recover the heat would become redundant. The presence of a dielectric between the two electrodes prevents the formation of hot plasma in a single spark, forming instead several microfilaments, resulting in a more uniform plasma. The low gas temperature allows the application of a catalyst directly in the plasma generation zone without fast deactivation, maximizing the interaction between active species and the catalytic phase. Furthermore, DBD plasma can be generated at atmospheric pressure, which is interesting from the application point of view. Very promising results were presented in the last years on several topics, such as CO$_2$ conversion [9,18–27] and CH$_4$ reforming [12,13,28–34], in terms of high conversion and selectivity. However, the main issue of DBD remains the low energy efficiency achieved, i.e., the ratio between chemical energy stored in the produced molecules and electrical energy applied, which rarely surpasses 10%. This is explained by dissociative excitation by electron impact, involving a large activation barrier, dominating over vibrational excitation [7,35–37]. A more suitable technique for vibrational excitation is microwave plasma, which uses GHz frequencies [38–40]. However, the temperature increase is more pronounced in microwave plasma, limiting the opportunities for plasma catalysis.

Interaction between plasma and catalyst can proceed in many ways [1–4]. Obviously, the plasma will introduce new chemical species including activated species, radicals, and ions, which may all adsorb on the catalyst opening new reaction pathways and influencing the products distribution. Plasma can also induce photocatalytic effects by UV irradiation, impingement of charged particles and thermal fluctuations. The surface and subsurface of a catalyst can be modified by plasma via poisoning, implantation, sputtering, and etching. The presence of a catalyst influences the plasma by changing the electrical field distribution, but also modifying the free volume and the residence time in the plasma zone. Plasma also affects the temperature of the system, obviously influencing reaction rates of chemical reactions.

Unfortunately, it is not possible to measure the temperature in a DBD plasma catalytic reactor directly. Application of a thermocouple inside the plasma is not possible due to the high electric fields present. Nevertheless, thermocouples were used a few millimeters outside the plasma zone inside the reactor tube, or outside the reactor tube just alongside the plasma zone [41–44]. Furthermore, many attempts have been done to measure temperatures indirectly, for instance by emission spectroscopy of UV–Vis radiation probing electronic transitions in nitrogen and hydroxyl groups [45–48], by UV absorption spectroscopy [47], or by infrared emission [23,44]. Unfortunately, these methods have serious limitation depending on the reactor material properties as well as the packed bed properties.

This study proposes a method to distinguish between thermal effect and plasma chemistry effects in fixed bed DBD plasma reactors. The decomposition of calcium carbonate is used as a model system for packed beds containing porous particles. It is well known that thermal decomposition results in formation of exclusively CaO and CO$_2$ [49], whereas formation of CO would indicate that plasma chemistry is involved. A pure thermal effect is likely when using an Ar plasma, since no chemistry is expected between activated argon species and CaCO$_3$. The choice for CaCO$_3$ as model system in combination with a DBD reactor is inspired by its relevance for CO$_2$ separation. The calcium looping cycle consists of carbonation of CaO for capturing followed by calcination of CaCO$_3$ in order to recycle calcium oxide and to produce pure CO$_2$ [49]. Bottleneck is the calcination reaction that requires high temperatures in order to achieve high CO$_2$ concentrations in the outlet, i.e., at least 950 °C to achieve
1 bar CO$_2$ [50]. Such temperatures result in sintering, decreasing the CO$_2$ capture capacity when calcium oxide is recycled multiple times [51–53]. Using a DBD plasma during the calcium carbonate decomposition might circumvent the need for such high temperatures, and in addition CO$_2$ will be converted by plasma into CO, converting electrical energy into chemical energy and producing an added-value product.

The method to distinguish between thermal and plasma chemistry is based on the fact that plasma cannot exist in the pores inside particles if they are smaller than a few micrometers, as can be understood from Paschen’s Law, which is generally accepted [54,55]. It was recently reported in a theoretical study that penetration of plasma in pores is possible to some extent [56]; however, we will discuss that under our conditions the plasma is limited to the interparticle volume and the external surface area is exposed under the conditions applied. The internal surface area, caused by the presence of small pores in the material is not exposed to the plasma, but would be influenced by any thermal effect. The theory is explained in detail in the section Methods in Appendix B and shown in Figure A13. The method then consists of two approaches: first, the effects of both the internal surface area as well as external surface area will be explored, and second, the dynamics of the performance on changing plasma power will be evaluated. Decomposition rate and eventual further reactions of the carbon dioxide product will be assessed. Argon plasma is used as a reference for the method to be developed as to distinguish between thermal effects and plasma–chemical effects.

2. Results

X-ray fluorescence (XRF) measurements confirmed the purity of CaO (99.12%) containing some minor impurities, i.e., SiO$_2$ (0.16%), MgO (0.12%), and Al$_2$O$_3$ (0.095%).

Table 1 in the Materials and Methods Section presents the surface area and particles sizes of the five prepared samples. The surface areas reported for samples A and B are well reproducible, but it should be noted that systematic errors may be larger, given the relatively low value of the surface areas. The surface area of batch III (samples C–E) is below the detection limit of the N$_2$ physisorption equipment, from which we deduce that the surface area is below 0.5 m$^2$/g. In any case, the total surface of the samples increases in the order C<B<A. Remarkably, the surface areas of the parent oxides are much higher, confirming the theory that formation of a carbonate layer induces closure of small pores, due to the lower density of CaCO$_3$ (2.71 g cm$^{-3}$) compared to CaO (3.35 g cm$^{-3}$). However, the order in surface area remains the same, reassuring that the surface area of the samples is in the order A>B>C>D>E. Figure 1 shows the pore size distribution measured with mercury porosimetry for the carbonated samples synthesized from calcium ascorbate (sample B, Figure 2a) and from calcium carbonate (sample C, Figure 2b). Sintering at 900 °C for 24 h (sample C) causes formation of large pores of typically 400 nm, compared to sample A which was treated with CO$_2$ only at 630 °C. Remarkably, the pore volume of the sintered sample is ~20%, much larger than sample B (~5%), indicating that smaller pores collapsed favoring enlargement of the bigger pores.

Table 1. Characteristics of the samples prepared from three different precursors.

<table>
<thead>
<tr>
<th>Sample Code</th>
<th>Precursor (Batch #)</th>
<th>Carbonation Time at 630 °C (h)</th>
<th>Sintering Time at 900 °C (h)</th>
<th>CaO S.S.A. (m$^2$ g$^{-1}$)</th>
<th>CaCO$_3$ S.S.A. (m$^2$ g$^{-1}$)</th>
<th>Particles Diameter (µm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>Ca Gluconate (I)</td>
<td>5</td>
<td>0</td>
<td>46.2</td>
<td>1.7 ± 0.1</td>
<td>250–300</td>
</tr>
<tr>
<td>B</td>
<td>Ca Ascorbate (II)</td>
<td>4</td>
<td>0</td>
<td>23.2</td>
<td>0.8 ± 0.1</td>
<td>250–300</td>
</tr>
<tr>
<td>C</td>
<td>CaCO$_3$ (III)</td>
<td>5</td>
<td>24</td>
<td>10.1</td>
<td>&lt;0.5</td>
<td>100–125</td>
</tr>
<tr>
<td>D</td>
<td>CaCO$_3$ (III)</td>
<td>5</td>
<td>24</td>
<td>10.1</td>
<td>&lt;0.5</td>
<td>38–45</td>
</tr>
<tr>
<td>E</td>
<td>CaCO$_3$ (III)</td>
<td>5</td>
<td>24</td>
<td>10.1</td>
<td>&lt;0.5</td>
<td>38–45</td>
</tr>
</tbody>
</table>
Changing the plasma power causes a fast response of the CO concentration in the order A–C, showing that the CO and CO kinetics, instead of thermodynamics, and thus also by the surface area.

Equilibrium CO experiments, between 900 to 1700 ppm, which is significantly lower than the 7000 ppm thermodynamic typically 20 min, after an induction time of a few minutes. In general, this is observed when decomposition is limited to max 50% of CaCO₃ present initially. The amounts of CO₂ desorbed from the samples A, B, and C are 4, 3.8, and 2.8 mg, respectively, equivalent with 0.66, 0.61, and 0.39 gCO₂/gCaO, respectively. The amounts of CO₂ decrease in the same order as surface area. It can be estimated that the thickness of the CaCO₃ layer is in the order of 30 nm for all three sample. This is consistent with the observation that the maximum CO₂ concentration during decomposition experiments, between 900 to 1700 ppm, which is significantly lower than the 7000 ppm thermodynamic equilibrium CO₂ concentration at 630 °C [50]. Consequently, the CO₂ concentration is determined by kinetics, instead of thermodynamics, and thus also by the surface area.

Figure 2 presents the results of isothermal decomposition at 630 °C of the carbonated samples A–C, showing that the CO₂ concentration generated via decomposition is within 10% constant during typically 20 min, after an induction time of a few minutes. In general, this is observed when decomposition is limited to max 50% of CaCO₃ present initially. The amounts of CO₂ desorbed from the samples A, B, and C are 4, 3.8, and 2.8 mg, respectively, equivalent with 0.66, 0.61, and 0.39 gCO₂/gCaO, respectively. The amounts of CO₂ decrease in the same order as surface area. It can be estimated that the thickness of the CaCO₃ layer is in the order of 30 nm for all three sample. This is consistent with the observation that the maximum CO₂ concentration during decomposition experiments, between 900 to 1700 ppm, which is significantly lower than the 7000 ppm thermodynamic equilibrium CO₂ concentration at 630 °C [50]. Consequently, the CO₂ concentration is determined by kinetics, instead of thermodynamics, and thus also by the surface area.

Figure 3 presents the effect of plasma power, measured with a Lissajous plot (see Appendix B), on the decomposition of carbonated sample C. No plasma was applied during the first two minutes and CO₂ is the only product observed, while in the presence of plasma CO and O₂ are also produced, next to CO₂. Every two minutes a different voltage is applied, and the power is measured after ca. 1 minute. Changing the plasma power causes a fast response of the CO concentration in the order
of seconds, while the CO2 concentration needs typically a minute to stabilize. The O2 concentration shows a delay; this effect is not understood at this time, but it may be speculated that interaction with CaO is responsible. Figure 3 shows that steady-state decomposition was achieved for the two lower plasma power values. The highest power setting of 9.6 W caused exhaustion so that the product concentrations are likely to be underestimated. Additional experiments were performed at constant maximum power as shown in Figure 4. Two experiments were performed at 2.1 W and demonstrate a reproducibility within ±5%, even though plasma was applied with a 1.5 min delay in the second experiment (Figure 4b). Figure 4c is performed at 4.4 W, showing a higher decomposition rate and shorter steady state duration. Therefore, the power was not further increased for this sample and in general the maximum power was limited to 5.1 W.

![Figure 3](image_url)

**Figure 3.** Concentration of the products of CaCO3 decomposition (sample C) as function of time at 630 °C in pure argon, flow rate of 30 mL/min, plasma power is changed every 2 min with values of 0, 0.4, 3.2, and 9.6 W.

![Figure 4](image_url)

**Figure 4.** Concentration of the products of CaCO3 decomposition (sample C) as function of time at 630 °C in pure argon, flow rate of 30 mL/min, plasma power is 2.1 W (a,b) and 4.4 W (c). Plasma is turned on at the beginning of decomposition (a,c) or after 1.5 min (b).

Figure 5 shows three typical results on the effect of the specific surface area on the decomposition at 2.1 W plasma power, by comparing the samples A (ex calcium gluconate, Figure 6a), B (ex calcium ascorbate, Figure 6b), and C (ex calcium carbonate Figure 6c), keeping particle size constant (250–300 μm). The total decomposition rate, as calculated based on the sum of the rates of formation of CO and CO2, seems to increase with increasing specific surface area, as also observed in the absence of plasma. During the decomposition of sample B, the power has been turned off after 9 min, resulting in a rapid decrease in the CO and CO + CO2 concentrations. Figure 6 presents all data on
the decomposition rate measured on the three samples when changing the plasma power, showing that the rate of decomposition at the same power is significantly lower for the sample with the lowest specific surface area. The difference between sample A and B is not larger than experimental scatter, although the data suggest a slightly higher rate for sample A.

![Figure 5](image-url) Concentration of the products of decomposition of sample A (a), sample B (b), and sample C (c), with a temperature of 630 °C in pure argon and flow rate of 30 mL/min; plasma power is 2.1 ± 0.1 W.

![Figure 6](image-url) Sum of CO₂ and CO concentrations plotted as function of power obtained during decomposition of CaCO₃ samples A, B, and C with different surface areas; all the experiments performed at 630 °C, in pure argon, and flow rate 30 mL/min.

Figure 7 shows the influence of particles size on CaCO₃ decomposition by comparing sample C (250–300 μm, Figure 7a), D (100–125 μm, Figure 7b), and E (38–45 μm, Figure 7c) at plasma powers varying between 1.3 and 2.1 W. Figure 5d shows details on the response time after switching on the plasma for the three samples. The time to reach steady state is 50, 40, and 10 seconds for samples C, D, and E, respectively. These times are in reasonable agreement with the Fourier times of CaCO₃ particles of these sizes. The different powers do not allow direct comparison of the decomposition rates. Instead, the effect of plasma power on the decomposition rate for samples C, D, and E is presented in Figure 8, showing that the particle size has no effect on the decomposition rate within experimental error. All the experiments addressed in Figures 6 and 8 are shown in detail in Figures A1–A6 in Appendix A.

Figure 9 shows the results of a series of plasma enhanced decomposition experiments in the presence CO₂ in the feed gas, measured on sample C. All experiments were done with one sample by performing 22 carbonation and decomposition cycles. The stability of the sample was verified by repeating decomposition measurements in the absence of plasma, demonstrating invariable results within 5% as shown in Figure A15 in Appendix A.
Figure 7. Concentration of the products of decomposition of sample C (a), sample D (b), and sample E (c) as function of time with a temperature of 630 °C in pure argon and flow rate of 30 mL/min; plasma power is 2.1 W for sample C, 1.3 W for sample D, and 1.7 W for sample E. (d) Initial response of CO + CO₂ concentration after activating the plasma for all three samples.

Figure 8. Sum of CO₂ and CO concentrations as function of power during decomposition of CaCO₃ samples with different particles size (samples C, D, and E); all the experiments were performed at 630 °C, in pure argon, and at a flow rate 30 mL/min.
Figure 9. CO + CO\textsubscript{2} (solid line), CO\textsubscript{2} (dashed line), CO (dash-dotted line), and O\textsubscript{2} concentrations monitored by mass spectrometry during CaCO\textsubscript{3} decomposition in different CO\textsubscript{2} feed concentrations in Ar and different plasma power. The amount of CaO is 6 mg, temperature is 630 °C, and the flow rate is 90 mL min\textsuperscript{−1}. The fed CO\textsubscript{2} concentration is 0 ppm (a), 1000 ppm (b), 2000 ppm (c), and 3200 ppm (d). The plasma power is 1.5 W (a), 1.3 W (b), 1.3 W (c), and 1.5 W (d).

Figure 9 shows the effect of CO\textsubscript{2} in the feed on plasma-enhanced decomposition at similar plasma powers, i.e., between 1.3 and 1.5 W. The reaction rate substantially decreases in presence of extra CO\textsubscript{2}, as can be estimated based on the sum of the concentrations of CO and CO\textsubscript{2} minus the CO\textsubscript{2} concentration at the inlet. Furthermore, the CO concentration increases with increasing CO\textsubscript{2} concentration.

In order to correct for the differences in power, additional experiments at other plasma powers were performed as shown in Figures A6–A9 in Appendix A and the results allowed interpolation of the CO and CO\textsubscript{2} values to 1.5 W plasma power, resulting in Figure 10. As the CO\textsubscript{2} concentration varies along the axis of the reactor, an averaged value is calculated according to Equation (1):

\[
[\text{CO}_2]_{ave} = \frac{[\text{CO}_2]_{in} + [\text{CO}_2]_{out}}{2}
\]  

(1)
well as the window of the CO₂ decomposition previously shown. The spread in the CO₂ concentration provides the window of concentrations occurring along the axis of the reactor; note this is not an error margin. Figure 11 confirms that the CO concentration indeed increases with increasing CO₂ concentration in the gas phase and this is observed for both CaCO₃ as well as the blank experiment with α-Al₂O₃ only, which is shown in Figure A10 in Appendix A.

Figure 10. CO outlet concentration vs. CO₂ average concentration obtained in a packed bed of 10 mg CaCO₃ and 90 mg Al₂O₃ at 630 °C with different CO₂ concentrations in Ar and a flow rate of 90 mL/min (squares); all results interpolated to 1.5 W plasma power. The same experiment was also performed with 100 mg Al₂O₃ only (circles). Error margins of the CO concentration are provided, as well as the window of the CO₂ concentrations (note these are not error margins).

The spread in the CO₂ concentration provides the window of concentrations occurring along the axis of the reactor; note this is not an error margin. Figure 11 confirms that the CO concentration indeed increases with increasing CO₂ concentration in the gas phase and this is observed for both CaCO₃ as well as the blank experiment with α-Al₂O₃ only, which is shown in Figure A10 in Appendix A.

Figure 11. Average increase of temperature in the plasma zone as function of power during the CaCO₃ decomposition previously shown.

3. Discussion

Scheme 1 presents a simple reaction scheme describing CaCO₃ decomposition, both in the absence and presence of plasma. Thermal decomposition without plasma involves exclusively R₁ (solid arrow), whereas in the presence of plasma both R₂ and R₃ might contribute. The observed enhancement of decomposition by plasma might be caused by plasma chemistry according R₂ and/or by enhancing R₁ via increasing temperature, as will be discussed below. Furthermore, we will discuss whether R₂ or R₃ is responsible for the formation of CO.

\[
\text{CaCO}_3(s) \xrightarrow{R_1} \text{CaO}(s) + \text{CO}_2(g) \]

\[
\text{CaO}(s) + \frac{1}{2} \text{O}_2(g) \xrightarrow{R_3} \text{CaO}_2(s)
\]

Scheme 1. CaCO₃ decomposition in Ar plasma.
3.1. Formation of CO

Figure 10 shows that the CO formation rate is not affected by the presence of CaCO$_3$ in the plasma at low CO$_2$ feed concentrations. It is enhanced by feeding additional CO$_2$. Hence, we conclude that CO formation occurs only in the gas phase (R$_3$) and not on the CaCO$_3$ surface (R$_2$). The trend in Figure 11 also indicates that the order of CO formation in the CO$_2$ concentration is clearly smaller than one. This is qualitatively in line with the results of Ramakers et al. [25] and Butterworth et al. [27], reporting that Ar can enhance the activation of CO$_2$, which is attributed to the fact that the electron density is enhanced because Ar is much easier ionized than CO$_2$. On the other hand, one may also speculate that higher CO$_2$ concentration increases the probability of recombination of CO and O, quenching the reaction.

The formation of CO responds extremely fast to switching the plasma on, as can be seen in Figures 3–5 and Figure 7, much faster than the CO$_2$ response as will be discussed below. This is in line with the conclusion that CO formation is plasma controlled. In cases when CO$_2$ concentration increase slowly in time (Figures 4 and 5), it can be seen that the CO concentration follows, which is in line with the conclusion that the consecutive pathway R$_3$ is dominant over R$_2$.

3.2. Thermal Effect or Plasma Chemistry?

The decomposition rate at a fixed power depends on the total surface area, as can be observed from Figure 6. On the other hand, the decomposition rate remains constant within experimental error when varying the particles size and consequently the external surface area, as shown in Figure 8. Therefore, the external surface area has no influence on the rate of decomposition at any power. According to the Paschen law, typically plasma cannot exist in pores smaller than 6 µm, implying that any plasma chemistry on the surface of CaCO$_3$ can exclusively contribute at the external surface of the particles. Recent work [56] revealed that penetration into relatively large pores is possible; however, the pores in CaCO$_3$ are much smaller, the electrical field is two orders of magnitude lower compared to [56], and the calculated penetration depth is limited to 5 µm, which is much smaller than the particle size used. Therefore, it seems reasonable to assume that penetration of plasma in this study is negligible. As the external surface area has no significant effect, we conclude that R$_2$ does not contribute. The enhancing effect of plasma power on the decomposition rate, as well as the observation that increasing the specific surface area increases the rate of decomposition (both in Figure 6), both suggest that plasma induced temperature increase is responsible for the increase in the decomposition rate. In fact, this is a result that can be expected operating with an Ar plasma because a chemical reaction between Ar ions and CaCO$_3$ would not be expected. Work on DBD plasmas containing H$_2$ and H$_2$O is ongoing in which plasma induced chemical reactions are much more likely, in which we will use the methodology developed in this work.

Figure 11 shows the apparent temperature increase, as estimated based on the temperature that would be required to account for the increase in decomposition rate, based on the kinetics of the decomposition reaction [57]. Remarkably, all observations converge to a single line independent of both surface area and particle size. The order of magnitude of the temperature increase is quite similar to results reported in literature. Typical temperatures estimated in DBD plasmas range up to typically 200 K [17,23,41–48]. It should be noted that determination of the temperature is cumbersome, e.g., the temperature of the exiting gas provides only a minimum value because of rapid heat exchange between small reactors and environment, whereas infrared cameras and UV–Vis spectroscopy measurements have limited accuracy. Nevertheless, the order of magnitude agrees well with our observations. In short, although experimental details vary, a temperature increase of 50 °C due to 4 W plasma power input is concluded.

The temperature regulation of the oven played an important role in our study. It stabilized the temperature at a few millimeters outside the low voltage electrode at 630 °C. Therefore, any power input from the plasma will result in a decrease of the electrical power to the oven. Hence, the
temperature effect of the plasma is actually larger than estimated above, in contrast to all experiments performed at room temperature without any kind of temperature control.

The conclusion that thermal effects are dominant is further supported by the fact that the typical response time of the decomposition rate is in the order of one minute for large particles and somewhat faster for small particles (Figures 4 and 8d). The order of magnitude agrees well with the Fourier time of CaCO$_3$ particles of the sizes used. In any case, the response times are longer than response times observed for CO formation, as discussed above, which is in line with the conclusion that decomposition is thermally controlled, whereas CO$_2$ dissociation is obviously plasma controlled.

The conclusion is also reinforced by the fact that the sum of CO$_2$ and CO concentrations obtained by decomposition of sample C at 630 °C and 3.2 W plasma power corresponds to the CO$_2$ concentration obtained by decomposition at 680 °C (i.e., increasing the temperature of 50 °C as calculated in Figure 11) without plasma of the same sample, as observed by comparing Figures A3a and A11 in Appendix A.

The proposed method is validated, since it enables to distinguish between thermal and plasma chemistry effect on the decomposition rate. This method is applicable to other systems in which a plasma is in contact with a fixed bed of porous particles, including supported catalysts. It should be noted that in general the increase in gas temperature in a DBD plasma cannot be neglected, as done frequently in many studies. Second, only the external surface of catalysts particles interacts with plasma: plasma–catalyst synergy is therefore maximal for nonporous catalytic materials. Further work is currently ongoing to apply this method for CaCO$_3$ decomposition in more reactive plasma, e.g., H$_2$ plasma, as well as for plasma catalytic conversions.

4. Materials and Methods

4.1. Plasma Reactor

Figure 12 shows a schematic representation of the equipment used to measure plasma enhanced decomposition of CaCO$_3$. The fixed bed reactor is fed with either pure Ar, or a mixture of Ar containing 5% CO$_2$. The temperature of the oven is controlled with a Euro-Therm controller with an accuracy of ±0.5 °C between room temperature and 1000 °C. The isothermal zone is 8 cm long at 900 °C, defined as the part of the reactor with less than ±1 °C temperature variation. A Quadrupole Mass Spectrometer (MS) measures the composition of the gas downstream of the reactor. The MS signal for CO$_2$ (44 m/e) is calibrated between 0.16% and 5% CO$_2$, resulting in a linear relationship as shown in Figure A14. The reactor is a 4 mm inner- and 6 mm outer diameter quartz tube. The inner electrode is a stainless-steel rod of 1 mm diameter placed coaxially in the center of the reactor section. The outer electrode is a 1 cm long stainless-steel tube with 6 mm inner diameter, enclosing a plasma zone of 0.035 cm$^3$ in volume. The amount of CaCO$_3$ sample was limited to 10.5 ± 0.3 mg in order to prevent CO$_2$ concentrations approaching thermodynamic equilibrium, thus minimizing reabsorption of CO$_2$. The 10 mg sample was mixed with 90 mg α-Al$_2$O$_3$, filling the plasma zone completely and preventing any bypassing. An AC voltage of up to 10 kV peak to peak was applied to the inner electrode with a frequency of 23.5 kHz using a PMV 500–4000 power supply, while the outer electrode is connected to the ground via a probe capacitor of capacity 4 nF. The power of the plasma was calculated using the Lissajous method by measuring the voltage on the inner electrode with a Tektronix P6015A high voltage probe and on the outer electrode with a TT–HV 250 voltage probe, as described in the literature [58]. A sample of Lissajous plot is shown in Figure A12 in Appendix B.
The precursors were calcined in 20% O₂ in N₂ at atmospheric pressure, heating the sample to 900 °C (heating rate 15 °C/min), and keeping the temperature at 900 °C for 3 h. The calcined products were pelletized (pressure 160 bar), crushed and sieved in different particle size range: 250–300 μm, 100–125 μm, and 38–45 μm.

4.3. Carbonation

Five CaCO₃ samples have been produced via carbonation of CaO. The oxide synthesized from calcium ascorbate (batch II) has been treated in situ with 5% CO₂ in Ar at 630 °C for 4 h (heating rate 15 °C/min), the other two batches I and III were treated in a calcination oven with 20% CO₂ in N₂ at 630 °C for 5 h (heating rate 15 °C/min). The oxide synthesized from calcium carbonate (III) was consecutively sintered in pure CO₂ at 900 °C for 24 h (heating rate 15 °C/min), as summarized in Table 1. The resulting three samples (A, B, and C) were crushed and sieved, obtaining particles sizes in the range between 250 and 300 μm. The material made from CaCO₃ was also obtained with smaller particles, between 100 and 125 μm (sample D) and between 38 and 45 μm (sample E), respectively.

4.4. Characterization

The specific surface area, pore volume, and pore size distribution of the samples were measured both in CaO form as well as in CaCO₃ form, after carbonation. The samples were first degassed at 300 °C in vacuum for 3 h. The surface area was calculated based on the BET isotherm for N₂ adsorption at –196 °C in a Tristar 3000 analyzer (Micromeritics, Norcross, GA, USA). The pore size distribution was measured by Hg porosimetry. The chemical composition was determined with X-ray fluorescence analysis in a S8 Tiger (Bruker, Billerica, MA, USA).

4.5. Experimental Procedure

The carbonated samples (ex situ) were heated up in 5% CO₂ in Ar to the temperature at which decomposition is to be measured, in order to prevent any premature decomposition. The decomposition reaction is initiated by switching the gas composition from 5% CO₂ to pure Ar, at a constant flow rate of 30 mL/min. Isothermal decomposition experiments have been performed at different plasma powers by varying the applied voltage. The plasma power was varied during the experiment in case of low decomposition rates, allowing observations of steady state CO₂ concentrations for each plasma power. In case CaCO₃ is exhausted too fast, only one single power was
applied. The rate of decomposition is calculated based on the sum of CO₂ and CO concentrations in the exit of the reactor as measured with MS.

Sample C (Table 1, ex calcium carbonate, 250–300 µm) was measured by performing 20 carbonation-decomposition cycles. The sample was recarbonated, after a decomposition experiment, in the reactor (in situ) by CO₂ absorption at 630 °C in 5% CO₂ in Ar for 30 min in a constant flow of 90 mL/min. The carbonated samples were decomposed using a constant power plasma in the presence of a relatively low CO₂ concentration in the feed, varied between 0 to 3200 ppm. During these 20 cycles, blank experiments were done every few cycles by decomposing in the absence of plasma, in order to ensure that the sample did not change in the course of the experiments.

Blank experiments were performed with 100mg of α-Al₂O₃ with particles size of 250 to 300 µm in the absence of any CaCO₃, operating with low CO₂ concentrations in the feed, i.e., 1000, 2000, and 3200 ppm. The plasma power was varied between 0 and 10 W and the responses of the CO and CO₂ concentrations were measured with MS.

5. Conclusions

The effect of argon plasma on calcium carbonate decomposition was herein assessed by means of a comparative method which allowed us to distinguish between thermal effects and plasma chemistry, based on reaction rates and dynamics. It represents a systematic method to distinguish between thermal effects versus plasma chemistry effect in fixed DBD plasma applications. Application of a DBD Ar plasma causes two effects when decomposing CaCO₃.

First, the rate of CaCO₃ decomposition increases. We conclude that this effect is purely a thermal effect, based on the fact that the rate of decomposition is enhanced when the total surface area is increased, whereas the external surface area has no influence. If the contact of plasma with CaCO₃ would dominate, the opposite would be expected. Furthermore, the dynamics of CaCO₃ decomposition follow the dynamics of heat transfer in CaCO₃ particles.

Second, plasma induces formation of CO. We conclude that this occurs via decomposition of CO₂ in the gas phase, based on the observation that the rate of CO formation is ruled by the CO₂ concentration as well as the observation that dynamic changes are very fast, as expected for plasma effect.


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Conflicts of Interest: The authors declare no conflict0s of interest.

Appendix A.

Appendix A.1. Decomposition in Argon Plasma of All the Samples

Figures A1–A5 show the results of all the CaCO₃ decomposition experiments on samples A–E, respectively, in presence of plasma and without recycling. The results with power higher than 6 W and further than the second step of power have been rejected due to exhaustion.
Figure A1. Concentration of the products of CaCO$_3$ decomposition (sample A) as function of time at 630 °C in pure argon and flow rate of 30 mL/min. Plasma power is 0.3, 1.4, and 8.4 W (a); 1.4, 3.5, and 6.4 W (b); 5.1 W (c); 3.9 W (d); 2.2 W (e); and 3.7 W (f). Plasma is turned on after 2 min and changed every 2 min (a,b), after 1.5 min (c), or at the beginning of the decomposition (d–f).
Figure A2. Concentration of the products of CaCO$_3$ decomposition (sample B) as function of time at 630 °C in pure argon and flow rate of 30 mL/min. Plasma power is 1, 4.1, and 5.8 W (a); 0.5, 2.1, and 8.5 W (b); and 2.1 W (c,d). Plasma is turned on after 2 min and changed every 2 min (a,b) or at the beginning of the decomposition (c,d).

Figure A3. Cont.
Figure A3. Concentration of the products of CaCO$_3$ decomposition (sample C) as function of time at 630 °C in pure argon and flow rate of 30 mL/min. Plasma power is 0.4, 3.2, and 9.6 W (a); 4.4 W (b); and 2.1 W (c,d). Plasma is turned on after 2 min and changed every 2 min (a) or at the beginning of the decomposition (b,c) or after 1.5 min (d).

Figure A4. Concentration of the products of CaCO$_3$ decomposition (sample D) as function of time at 630 °C in pure argon and flow rate of 30 mL/min. Plasma power is 0.2, 2.3, and 8.1 W (a) and 1.7 W (b). Plasma is turned on after 2 min and changed every 2 min (a) or after 1.5 min (b).

Figure A5. Concentration of the products of CaCO$_3$ decomposition (sample E) as function of time at 630 °C in pure argon and flow rate of 30 mL/min. Plasma power is 4.4, 1, 4.9, and 1 W (a) and 1.7 W (b). Plasma is turned on after 2 min and changed every 2 min (a) or after 1.5 min (b).

Figures A6–A9 show the results of all the CaCO$_3$ decomposition experiments on sample C in presence of plasma when the sample is recycled 20 times, including experiments with extra CO$_2$ added.
Figures A6–A9 show the results of all the CaCO₃ decomposition experiments on sample C in presence of plasma when the sample is recycled 20 times, including experiments with extra CO₂ added.

Figure A6. Cont.
Figure A6. Concentration of the products of CaCO₃ decomposition (sample C) as function of time at 630 °C, flow rate of 90 mL/min. The fed CO₂ concentration is 0 ppm and the plasma power is 1.5 W (a), 2.1 W (b), 1.7 W (c), 2.2 W (d), 1.5 W (e), 2 W (f), and 1.1 W (g). The plasma is applied at the beginning of the reaction. The decomposition cycles are in order #20, #21.

Figure A7. Concentration of the products of CaCO₃ decomposition (sample C) as function of time at 630 °C, flow rate of 90 mL/min. The fed CO₂ concentration is 1000 ppm and the plasma power is 1.3 W (a) and 1.5 W (b). The plasma is applied at the beginning of the reaction. The decomposition cycles are in order #20, #21.

Figure A8. Concentration of the products of CaCO₃ decomposition (sample C) as function of time at 630 °C, flow rate of 90 mL/min. The fed CO₂ concentration is 2000 ppm and the plasma power is 1.4 W (a) and 1.9 W (b). The plasma is applied at the beginning of the reaction. The decomposition cycles are in order #13, #14.
Figure A9. Concentration of the products of CaCO₃ decomposition (sample C) as function of time at 630 °C, flow rate of 90 mL/min. The fed CO₂ concentration is 3200 ppm and the plasma power is 1.4 W (a) and 1.5 W (b). The plasma is applied at the beginning of the reaction. The decomposition cycles are in order #16, #17.

Appendix A.2. CO₂ Dissociation with Only α-Al₂O₃

Figure A10 shows the result of CO₂ dissociation experiments on a sample consisting of 100 mg α-Al₂O₃ with particle size of 250–300 μm. The applied gas composition is respectively 1000 ppm (Figure A10a), 2000 ppm (Figure A10b), and 3200 ppm (Figure A10c) in Argon. The CO concentrations at 1.5 W plasma power are obtained by interpolation and showed in Figure 11.

Figure A10. Concentration of the products of CO₂ dissociation on 100mg α-Al₂O₃, particle size of 250–300 μm, as function of time at 630 °C, flow rate of 90 mL/min. The fed CO₂ concentration is 1000ppm (a), 2000ppm (b), and 3200ppm (c); the plasma powers are 0.6, 1.6, 3.1, 4.8, and 7.2 W (a); 0.7, 1.8, 4.1, and 5.6 W (b); and 0.5, 1.5, 3.1, and 4.9 W (c). The plasma is applied for a few minutes and then turned off.
Figure A11 shows the result of a CaCO3 decomposition experiment at 680 °C without plasma on sample C. It is observed that the CO2 concentration at the steady state corresponds within the error margins to the sum of CO and CO2 concentrations during decomposition at 630 °C with 3.2 W plasma power, shown in Figure A3a. The comparison supports the estimation of a 50 °C increase for 3.2 W, as shown in Figure 11.

![Graph showing CO2 concentration over time](image)

**Figure A11.** Concentration of CO2 during CaCO3 decomposition (sample C) as function of time at 680 °C and a flow rate of 30 mL/min.

**Appendix B.**

**Appendix B.1. Determination of Plasma Power**

The Lissajous plot measured during decomposition of sample C in Ar with 3200 ppm extra CO2 and plasma power 1.5 W, indicating plasma charge as function of voltage for one cycle, is shown in Figure A12. The hysteresis has a shape of parallelogram, indicating that a plasma is generated inside the reactor. The slopes of the two sides indicate respectively a cell capacity of 2.8 pF with plasma off and an effective capacity of 8 pF corresponding to the dielectric capacity and indicating that the plasma occupies the whole empty volume within the electrodes. The power of the plasma is calculated based on the area of the parallelogram:

\[
P = \frac{1}{T} \int_{0}^{T} VdQ
\]

(2)

![Lissajous plot diagram](image)

**Figure A12.** Lissajous plot of 1.5 W plasma applied on the CaCO3 and Al2O3 mixture with 3200 ppm of CO2 in Ar.
Appendix B.2. Method

The assumptions made in order to evaluate the argon plasma effect on calcium carbonate decomposition are the following:

- According to the Paschen’s Law the electron filaments and active species cannot exist inside the pores of the carbonate particles, this is explained in the following paragraph.
- Temperature gradients between particles outer surface and inner core are minimized within much smaller times than the duration of the experiment, i.e., within 2 min, according to the Fourier time calculated in the specific case.

Figure A13 shows the voltage needed to generate a plasma as function of the distance between the charges at an argon pressure of 1 atm, where the parameters $A$, $B$, and $\gamma$ are determined experimentally [59]. The physical mechanism can be explained for the 2 branches of the curve:

- $d_{\text{gap}} > 6\mu m$: In order to generate a discharge (that is produced by an avalanche of ionizations of Argon atoms by high energy electrons), we need a minimum electric field that is able to accelerate electrons above the ionization energy of Ar atoms in the mean free path of electrons at that pressure (ca. 0.5 $\mu$m). Since the electric field is inversely proportional to $d_{\text{gap}}$ at constant voltage, $V_b$ increases linearly with $d_{\text{gap}}$.
- $d_{\text{gap}} < 6\mu m$: the gap distance is comparable with the mean free path, so the probability of collisions that produce ionizations decreases. In order to have enough collisions to produce a discharge, the voltage has to be increased.

It follows that plasma cannot form in the pores of CaCO$_3$ particles within the plasma zone, since their diameter is below 2 $\mu$m.

![Figure A13. Burning voltage in function of the gap distance for argon at 1 atm, as calculated by [59].](image)

In order to discriminate whether the gas temperature increment or the active species drives the CaCO$_3$ decomposition in presence of plasma, we compared the effect of plasma on the reaction rate on samples with different specific surface area and perform a similar comparison for samples with different particles size. By choosing different particles sizes or different specific surface areas, we can vary selectively the external surface, exposed to the active species and temperature increase, and the internal surface, exposed only to a temperature increase.

Appendix B.3. CO$_2$ and CO and Calibration

Figure A14 shows the CO$_2$ signal elaborated by MS (m/e = 44) for CO$_2$ concentrations up to 5% in Ar and in the inset CO$_2$ concentrations up to 1000 ppm are expanded. In this range the signal is
linear with the CO\textsubscript{2} concentration, with a R\textsuperscript{2} coefficient of 0.9996. The CO signal (m/e = 28) for the CO concentration of 5% in Ar is 1.15 ± 0.05 times larger than the CO\textsubscript{2} signal (m/e = 44) at the same CO\textsubscript{2} concentration.

![Graph showing the relationship between CO\textsubscript{2} concentration and MS signal for m/e = 44](image)

**Figure A14.** Mass spectrometer signal m/e = 44 for different CO\textsubscript{2} concentrations in Ar from 0 to 5%; inset: enhancement of the CO\textsubscript{2} concentration range of 0 to 1000 ppm.

**Appendix B.4. Sample Stability during Absorption–Desorption Cycles**

Figure A15 shows the aging of the sample synthesized from calcium carbonate with the largest particles size, within 22 absorption–desorption cycles: the concentration of CO\textsubscript{2} during the initial stage of decomposition does not show significant change during the 21 absorption–desorption cycles at 630°C in absence of extra CO\textsubscript{2} and plasma, as expected for a sample already sintered at 900 °C for 24 h. The experiments in presence of extra CO\textsubscript{2} and plasma were all performed from cycle 9 to cycle 22, where the sample shows the highest stability.

![Graph showing CO\textsubscript{2} concentration over cycles](image)

**Figure A15.** CO\textsubscript{2} concentration at the first stage of decomposition of 10 mg CaCO\textsubscript{3} for different cumulative number of cycles. The decompositions are performed at standard conditions, i.e., 630 °C without plasma in absence of CO\textsubscript{2} in the gas mixture.

**Appendix B.5. Heat Transfer in the Packed Bed**

Plasma is considered as a uniform heat source at temperature T\textsubscript{∞} around the particles which are assumed spherical and have an initial temperature T\textsubscript{0}. We calculated the time needed to increase the temperature of the center of the particles of the 90% of the initial temperature difference T\textsubscript{∞}−T\textsubscript{0} as follows [60]

\[
t = \frac{\rho C_p V}{h A} \ln \left( \frac{T_0 - T_\infty}{T - T_\infty} \right)
\]  

(3)
The physical constants in the equation are the following: \( \rho \) is the density of CaCO\(_3\) (2710 kg/m\(^3\)); \( C_p \) is the heat capacity of CaCO\(_3\) at 630 °C (1.26 J g\(^{-1}\) K\(^{-1}\)) \[61\]; V and A are the volume and section of the CaCO\(_3\) particles, respectively, and h is the convective heat transfer coefficient of a dry gas in a packed bed reactor, approximated to 15 W m\(^{-2}\) K\(^{-1}\) \[60\]. As a result, the times needed to heat up the particles are respectively 4, 11, and 24 seconds for the smallest (38–45 µm), medium (100–125 µm), and largest (250–300 µm) particles, respectively. Equation (1) is an approximation for spherical particles without porosity, so we assume that more accurate calculation will give slightly longer times.

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