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Resistive Intrinsic ZnO Films Deposited by Ultrafast Spatial ALD for PV Applications

Naomi Nandakumar, Bas Dielissen, Diana Garcia-Alonso, Zhe Liu, Roger Görtzen, Wilhelmus M. M. (Erwin) Kessels, Armin G. Aberle, and Bram Hoex

Abstract—Spatial atomic layer deposition (ALD) of intrinsic zinc oxide (*i*-ZnO) films is scaled up from the laboratory to the industrial level, and the film properties are investigated in detail. A high growth rate of 35 nm/min is achieved. The deposited films are transparent and have an unusually high resistivity of about 100 Ω · cm. This is attributed to the extremely short precursor exposure and purge duration of spatial ALD (~8 ms), as compared with temporal ALD (~1–10 s). The growth of highly crystalline and nearly stoichiometric *i*-ZnO films is achieved. This makes these *i*-ZnO layers ideal for applications as insulating window layers in Cu(In,Ga)Se₂ solar cells.

Index Terms—Atomic layer deposition (ALD), Cu(In, Ga)Se₂ (CIGS), photovoltaics, solar cells, spatial, thin films, transparent conductive oxides, zinc oxide.

I. INTRODUCTION

INTRINSIC zinc oxide (*i*-ZnO) is a widebandgap semiconductor with a variety of useful properties that have made it a material of increasing attention in numerous fields of research [1]. The high transparency of ZnO in the visible wavelength range combined with its tunable electrical conductivity enables its use in devices such as transparent electronics [2], gas sensors [3], and solar cells [4]. ZnO has a direct bandgap of 3.37 eV at room temperature [5]. Thin films of ZnO are used in the photovoltaic (PV) industry in the undoped, as well as doped form, for example, in record-efficiency Cu(In,Ga)Se₂ (CIGS) thin-film solar cells and in silicon heterojunction solar cells [6]–[8]. In a standard CIGS configuration, a ZnO bilayer is used, consisting of a resistive *i*-ZnO layer capped with a doped ZnO layer. This

thin *i*-ZnO layer has been an integral component of CIGS solar cells over the years, because it has been consistently shown to improve reproducibility and production yield for this type of solar cell [4]. This improvement is attributed to the role of resistive *i*-ZnO in preventing the formation of shunts between the highly doped ZnO front contact and the molybdenum rear contact through pin holes and crevices in the inhomogeneous absorber layer [9]. The presence of the resistive *i*-ZnO layer results in higher resistive losses and increased carrier recombination at the buffer/absorber interface due to an increased separation between the Fermi level and the conduction band of the absorber [10]. However, numerical modeling and experimental studies have shown that this high series resistance provided by the *i*-ZnO layer prevents spatial electrical inhomogeneities (e.g., low-performance areas of locally enhanced recombination) from dominating the open-circuit voltage of the entire device [10]–[12]. This explains the high solar cell efficiencies obtained even when ~100-nm *i*-ZnO layers are used [13].

The preferred deposition technique for *i*-ZnO films in the PV industry is direct current and reactive sputtering from ceramic and metal targets, respectively [4]. However, sputtering of *i*-ZnO films has been shown to cause significant surface damage to the underlying absorber layer due to highly energetic ions and other plasma-induced effects [14], [15]. Consequently, there is a need for a low-damage technique for the deposition of the *i*-ZnO layer.

Atomic layer deposition (ALD) is one such deposition technique. Due to its self-limiting chemical reactions at low substrate temperatures, uniform growth and precise thickness control can be achieved in addition to low surface damage. Conventional (also known as temporal) ALD has been a less popular choice in the PV industry due to its low deposition rate of <2 nm/min. One way to increase throughput is by using batch reactors. Alternatively, the recent development of the spatial ALD technique achieves high growth rates by separating the gaseous precursors in space rather than in time [16]. A moving substrate is sequentially exposed to each precursor as illustrated in Fig. 1, removing the requirement for separate pumping cycles between each exposure. This combines the advantages of conventional ALD (i.e., the growth of uniform, pinhole-free, and highly conformal films on large-area substrates) with high deposition rates, making it an attractive option for the PV industry.

In 2011, Illiberi *et al.* demonstrated the deposition of *i*-ZnO films by spatial ALD on a prototype lab-scale system consisting of a circular reactor head with inlets for each precursor [17]. The substrate was placed on a circular table that rotated underneath the reactor head. In the present work, the deposition

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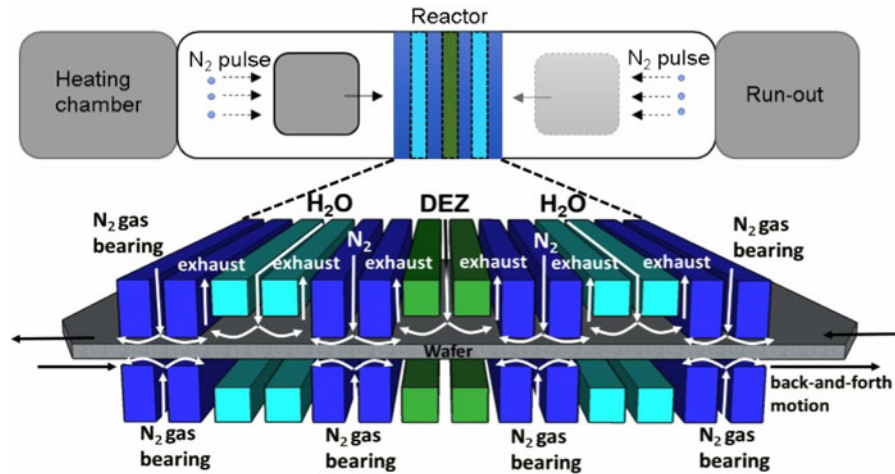


Fig. 1. Schematic illustration of the working principle of spatial ALD. The gaseous precursors (DEZ and H₂O) are separated in space rather than in time, as in the case of temporal ALD. The substrate is moved back and forth by N₂ pulses on either side of the reactor. Each pass results in one ALD cycle.

of *i*-ZnO has been successfully scaled up using a commercial spatial ALD system currently used in the PV industry. Highly uniform *i*-ZnO films with precisely controlled thickness, high transparency (>90% in the visible wavelength range), and exceptionally high resistivity ($\sim 100 \Omega \cdot \text{cm}$) were deposited at growth rates up to $\sim 35 \text{ nm/min}$ at low deposition temperatures ($< 250 \text{ }^\circ\text{C}$) and at atmospheric pressure. In earlier work, we showed that ultrafast ALD can be used to deposit uniform *i*-ZnO films at low temperatures and with low precursor consumption [18]. In this paper, we build on that work by investigating the properties of the *i*-ZnO films in detail and furthermore studying the effect of the short exposure and purge durations of spatial ALD on its electrical and structural properties.

II. EXPERIMENTAL DETAILS

i-ZnO films were deposited using SoLayTec's InPassion LAB spatial ALD system. A schematic of the system is shown in Fig. 1. Diethyl zinc (DEZn SSG, AkzoNobel) and deionized water vapor (H₂O) were used as the ALD precursors. Nitrogen (N₂ > 99.999%, Air Liquide) was used as the carrier gas for both precursors and also as a gaseous separation curtain for the precursor gases in the reactor. All precursors were transported to the reactor through heated lines to prevent condensation. Both DEZ and H₂O were further heated to the deposition temperature in a heat exchanger before being injected into the reactor through slits. Each slit is flanked by exhaust channels that carry away excess gas on each side, with two adjacent slits sharing one exhaust as shown in Fig. 1. During deposition, the sample is supported by N₂ gas bearings from the top and bottom which keep the sample floating in the reactor throughout deposition. The sample is also moved steadily back and forth inside the reactor at a fixed wafer speed, using pulses of N₂ at both ends of the reactor. During each pass, the wafer is exposed to one DEZ slit and two H₂O slits for each ALD cycle. Hence, in such a spatial ALD system, the pulse (exposure to DEZ and H₂O) and purge durations are fixed for a certain wafer speed. Unlike in temporal ALD, here, the pulse for each reactant and the purge time in between pulses cannot be independently varied.

Typical DEZ-based *i*-ZnO deposition temperatures for ALD range from 100 to 250 °C [19]. In the present work, this same temperature range was used for the reactor and precursor gases. Prior to the ALD process, the wafer is heated from room temperature to the deposition temperature in the run-in heating chamber and in an inert N₂ atmosphere.

i-ZnO films were deposited on pseudosquare 125 mm \times 125 mm p-type Czochralski *c*-Si substrates with a (1 0 0) crystal orientation. However, 156 mm \times 156 mm wafers can also be easily processed in the same system. All substrates were cleaned using a standard Radio Corporation of America clean [20], followed by a dip in 1% hydrofluoric acid with a final rinse in water before drying. In order to enable spectroscopic ellipsometry (SE) measurements, the films were deposited on single-side polished *c*-Si substrates. The SE substrates were prepared by chemical-mechanical polishing, resulting in a surface roughness of less than 5 Å. For electrical measurements, the films were deposited on as-procured polished *c*-Si samples with a 150-nm insulating layer of thermally grown SiO₂.

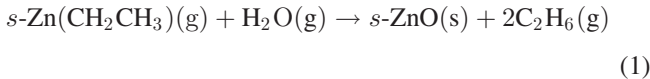
The thickness and optical properties of all films were measured using a spectroscopic ellipsometer (J. A. Woollam, M2000) in the visible and near-infrared (NIR) wavelength range from 250 to 1700 nm at 65°, 70°, and 75°. The pSemi-M0 oscillator model was used to fit the measured SE data [21]. The refractive indexes (n) and extinction coefficients (k) obtained from the ellipsometry measurements were used to simulate transmission and absorption spectra of the *i*-ZnO films for a typical CIGS solar cell structure, using the transfer matrix method in Mathematica. The n and k values of magnesium fluoride (MgF₂), cadmium selenide (CdS), and CIGS were obtained from the literature [22]–[24], while those of Al-doped ZnO (AZO) were measured in-house. Electrical properties were measured using a four-point probe mapping system (4PP, Advanced Instrument Technology, CMT-SR2000N-PV). The structural characterization of the films was done by scanning electron microscopy (SEM, Carl Zeiss, Auriga) and X-ray diffraction (XRD, PANalytical, X'Pert Pro MRD). X-ray photoelectron spectroscopy (XPS, Thermo Scientific K-Alpha, KA1066) with monochromatic Al K _{α} X-ray radiation ($h\nu = 1486.6 \text{ eV}$, spot

size = 400 μm) was used to assess the chemical composition of the films. All samples were corrected for sample charging with a flood gun. The sensitivity factors used for quantification were calibrated with Rutherford backscattering spectrometry as reported elsewhere [25]. For depth profiling, an Ar gun was applied to sputter the *i*-ZnO films.

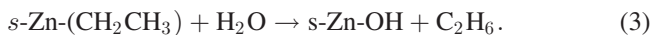
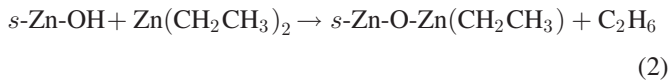
III. RESULTS AND DISCUSSION

A. Film Growth Characteristics

DEZ and H_2O used for ALD of *i*-ZnO react as follows:



where *s*- denotes surface species. The above reaction can be split further into two self-limiting half reactions that occur on exposure to each precursor [26]:



The reaction of DEZ with H_2O is exothermic and can be used to deposit *i*-ZnO thin films at substrate temperatures much lower than are required for other processes with different precursors. In Fig. 2(a), the relation between film thickness and the number of ALD cycles is shown for an *i*-ZnO film deposited at a substrate temperature of 225 $^\circ\text{C}$. It can be seen that the film thickness increases linearly with the number of ALD cycles, as is expected for an ALD process. The linear fit of the data in Fig. 2 indicates a short nucleation delay, which is quite typical for ALD of *i*-ZnO. The increase in thickness was linear after this initial lag. From the slope of the line, the growth per cycle (GPC) was found to be ~ 0.19 nm/cycle for a DEZ partial pressure of 2 mbar, an H_2O partial pressure of 65 mbar, and a pulse and purge time of ~ 8 ms each. This GPC value is found to correlate well with the interplanar distances in a hexagonal unit cell of *i*-ZnO, which range from 0.25 nm for the (1 0 1) planes to 0.26 and 0.28 nm for the (0 0 2) and (1 0 0) planes, respectively, indicating that the *i*-ZnO GPC is very close to ideal monolayer growth values [27]. This observation corresponds well to existing reports of *i*-ZnO GPC values by temporal ALD as well as spatial ALD in this deposition temperature range [5], [17]. Each cycle takes ~ 0.3 s, resulting in a deposition rate of ~ 35 nm/min for a GPC of 0.19 nm/cycle.

The reported optimum ALD temperature window and the temperature at which growth saturation for *i*-ZnO is achieved varies largely [5]. It is likely that reactor design and other factors such as the position of temperature sensors, reactor shape, and reactor size influence the values reported for ALD temperature windows. A temperature range from 100 to 250 $^\circ\text{C}$ was chosen in this study. Fig. 2(b) illustrates the GPC of *i*-ZnO films as a function of the substrate temperature. Maximum GPC for *i*-ZnO growth was achieved at 200 $^\circ\text{C}$, after which the GPC remained constant at ~ 0.19 nm/cycle. The DEZ and H_2O partial pressures were optimized in previous work to identify the

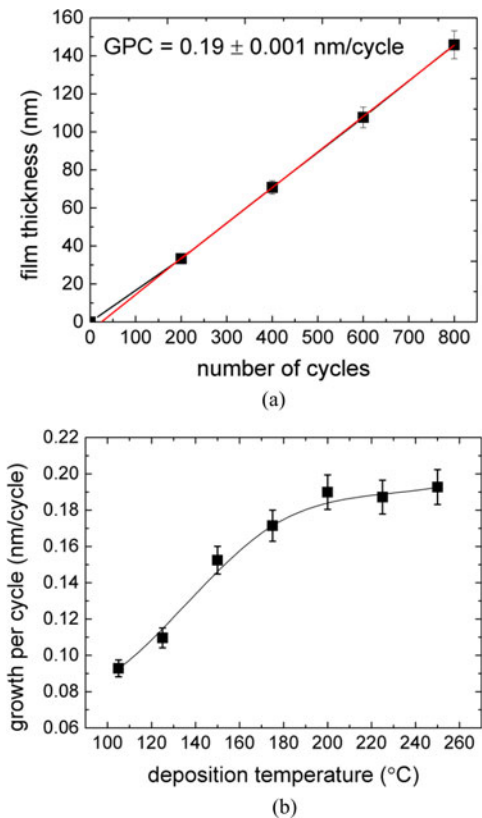


Fig. 2. (a) Film thickness as a function of the number of ALD cycles at a deposition temperature of 225 $^\circ\text{C}$. GPC was determined from the slope of the curve determined from a linear fit of the measured data. (b) GPC as a function of deposition temperature for a DEZ partial pressure of 2 mbar, a water partial pressure of 65 mbar, and a pulse and purge time of ~ 8 ms each.

precursor saturation regime [15]. Another point worthy to note is that some studies show a drop in the GPC of ALD deposited *i*-ZnO for substrate temperatures above 175 $^\circ\text{C}$ [28], [29]. This is attributed to “precursor desorption” or specifically to the decrease in hydroxyls at the surface at higher temperatures. A detailed study by Nelson *et al.* showed that this desorption is influenced by the purge time, whereby longer purge cycles result in a decreased GPC [30]. In comparison, spatial ALD offers the advantage of very short purge durations, maintaining a high GPC even at higher deposition temperatures.

The spatially resolved thickness was measured for a polished 125 mm \times 125 mm *c*-Si wafer and the thickness nonuniformity across the wafer was calculated using the following formula:

$$\text{Nonuniformity} = \frac{d_{\text{max}} - d_{\text{min}}}{2d_{\text{avg}}} \times 100 \quad (4)$$

where d_{max} , d_{min} , and d_{avg} are the maximum, minimum, and average thicknesses, respectively. With this formula the average nonuniformity value for a 100-nm-thick *i*-ZnO film was $\sim 3.5\%$. This value is $\sim 1\%$ lower than the film nonuniformity values reported in the literature for *i*-ZnO films deposited by temporal ALD [31]. By virtue of its self-saturating nature, ALD films should have high uniformity; nevertheless, nonuniformity in

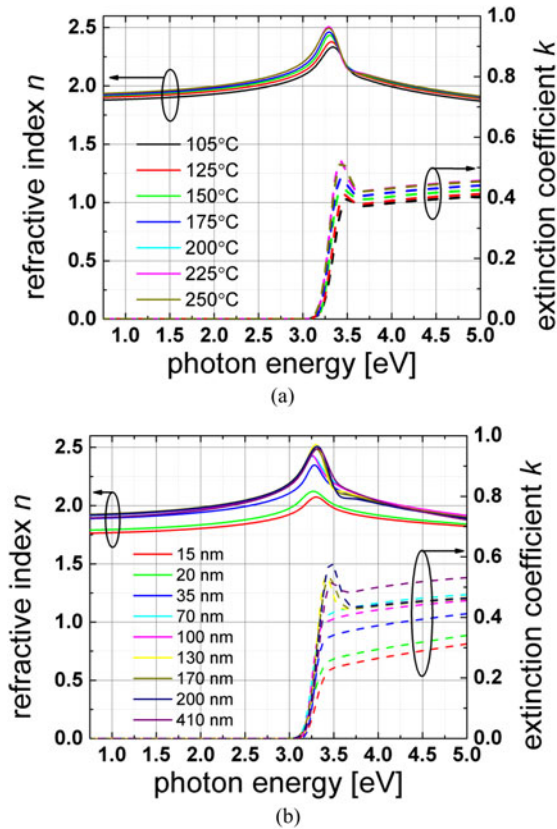


Fig. 3. Refractive index (n) and extinction coefficient (k) as a function of photon energy of i -ZnO films deposited (a) at various temperatures and (b) with various thicknesses using a fixed deposition temperature of 230 °C.

the thickness of ALD films have been consistently observed, reported, and studied [31].

B. Optical and Structural Properties

Along with film thickness, the SE measurements were also used to extract n and k values of the i -ZnO films. Fig. 3 shows the n and k values as a function of photon energy for i -ZnO films (600 cycles) grown at different deposition temperatures. The film thicknesses vary with temperature due to the temperature dependence of the GPC. The measured dielectric functions are in close agreement with spectra reported in the literature for i -ZnO films [32], [33].

The refractive index of the deposited i -ZnO films was found to be between 1.8 and 2.0 at a wavelength of 632 nm, which is consistent with most reports in the literature [17]. The n curves show a peak at the bandgap energy, as shown in Fig. 3(a). A shift in this n peak is observed with increasing deposition temperature. This shift can be attributed to a change in the film quality with increasing deposition temperature. The absorption edge of the k spectra was found to be at ~ 3.3 eV, which corresponds closely to the i -ZnO bandgap. However, a distinct peak is observed at ~ 3.4 eV, which is attributed to excitonic energies in wurtzite ZnO [34], [35]. The excitonic interaction of electrons and holes has a strong influence on the absorption spectrum, preventing the accurate determination of the bandgap value by

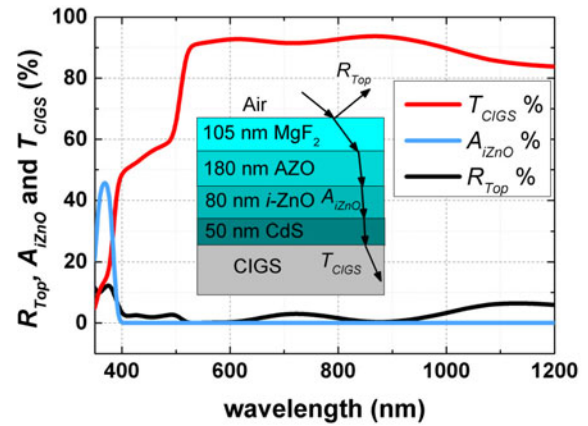


Fig. 4. Simulated reflection from the top surface (R_{Top}), absorption in the i -ZnO layer (A_{iZnO}), and transmission to the CIGS absorber (T_{CIGS}) for i -ZnO films applied as the window layer in a CIGS solar cell structure. (Inset) CIGS cell structure and interfaces used for the simulation.

extrapolation of the absorption edge. On the other hand, the presence of the excitonic peak indicates that the i -ZnO film is polycrystalline with possibly large crystal grains, indicating good film quality.

In Fig. 3(b), the shift in the n -spectra and the absence of the excitonic peak in the k -spectra for films with thickness of less than 50 nm could indicate that the thinner films are not completely developed in terms of crystallinity. However, it is also possible that the fit results for the very thin films are less reliable due to a relatively larger influence of the underlying SiO_2 layer, which has a very low refractive index. However, the presence of the excitonic peak for higher thicknesses clearly indicates a high crystalline content in the i -ZnO films. Structural investigations using SEM presented below for different i -ZnO thicknesses further support these observations.

The n and k values of i -ZnO deposited by spatial ALD obtained from SE measurements were further used to simulate the reflection, transmission, and absorption spectra for relevant interfaces in a CIGS solar cell structure with an i -ZnO layer between the Al-doped ZnO front contact and the CdS buffer layer, as shown in the inset of Fig. 4. The simulated spectra are presented in Fig. 4 for an 80-nm-thick i -ZnO film deposited at 230 °C. The i -ZnO film shows very low absorption (A_{iZnO}) at wavelengths above 400 nm. Hence, it gives high transmission to the CIGS absorber (T_{CIGS}) of $>90\%$ in the visible and NIR wavelength ranges, which makes the developed i -ZnO films well suited for application as transparent window layers.

Fig. 5 shows SEM images of the top surface and cross section of a thick (400-nm) i -ZnO film on a SiO_2 covered c -Si wafer. Columnar crystal grains with a grain size of the order of 50 nm can be seen. The columnar crystal structure was not observable for very thin i -ZnO films of <50 -nm thickness. We believe that thin i -ZnO films do not have columnar grains, although it needs to be mentioned that clear SEM images were difficult to obtain for such samples due to their high resistivity, which resulted in electron charging. The large crystalline grains observed for the thicker i -ZnO films confirm their highly polycrystalline nature.

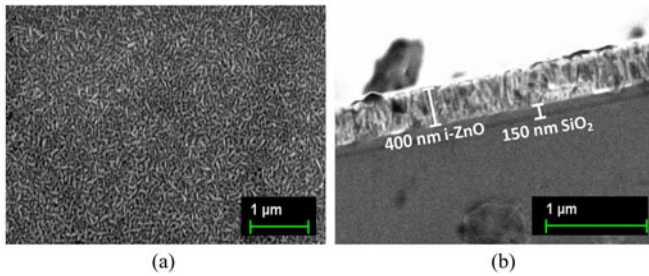


Fig. 5. (a) SEM micrograph (plan view) of a 400-nm-thick *i*-ZnO film deposited by spatial ALD on a *c*-Si/SiO₂ substrate. (b) SEM micrograph (cross section) of the sample of graph (a), showing a columnar grain structure.

TABLE I
PROPERTIES OF INTRINSIC ZNO FILMS DEPOSITED BY DIFFERENT TECHNIQUES

Deposition technique	Deposition temperature [°C]	Resistivity [$\Omega \cdot \text{cm}$]	Comments	Ref.
PECVD	250	3×10^{-3}	UV exposure ~ 1 min	[41]
PLD	200	$4\text{--}7 \times 10^{-3}$		[42]
Sputtering	60–120	$\sim 4 \times 10^{-1}$		[43]
MOCVD	620	42.7		[44]
Spray Pyrolysis	180	10^{-3}	Annealed in N ₂ at 400 °C	[45]–[46]
Sol-Gel	80	28.2	Annealed in air at 500–600 °C	[47]
Temporal ALD	200–220	$\sim 1 \times 10^{-3}$	Exposure ~ 1 s, Purge ~ 10 s	[5]
Spatial ALD-prototype	75–250	4×10^{-3}	Exposure ~ 30 ms Purge ~ 0.3 s	[17]
This work	200–225	$\sim 100\text{--}200$	Exposure ~ 8 ms Purge ~ 8 ms	

*PECVD = Plasma-enhanced chemical vapor deposition, PLD = Pulsed laser deposition, MOCVD = Metal-organic chemical vapor deposition, ALD = Atomic layer deposition.

This also corresponds well with the observations from the *n* and *k* spectra in Fig. 3.

C. Electrical Properties

i-ZnO films are intrinsically n-type and are frequently reported to be highly conductive, with resistivity values as low as $10^{-3} \Omega \cdot \text{cm}$ being reported for as-grown *i*-ZnO without any additional doping, as summarized in Table I. In contrast, the *i*-ZnO films deposited in this study all showed very high sheet resistances of the order of $10^7\text{--}10^8 \Omega/\text{sq}$ for 100-nm films, corresponding to a film resistivity in the range of 100–200 $\Omega \cdot \text{cm}$; see Fig. 6. Nelson *et al.* [30] recently showed that the resistivity of *i*-ZnO films exponentially increases with decreasing purge time after the DEZ exposure. They reported sheet resistance values of $\sim 10^5 \Omega/\text{sq}$ for a 100-nm-thick *i*-ZnO film deposited with pulse and purge times of 80 ms. Specifically, their study showed that for a fixed exposure time (2 s), the DEZ purge time influences the resistivity much more than the H₂O purge time.

Here, it is worthy to note that the exposure and purge cycles for this spatial ALD are asymmetric, in contrast with conventional temporal ALD. The details of the spatial ALD process used in this study are shown in Fig. 7. Unlike in temporal ALD

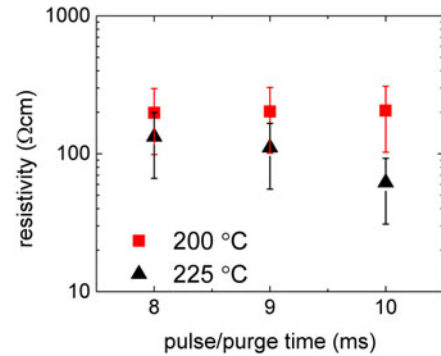


Fig. 6. Film resistivity as a function of pulse and purge duration (ms) for 100-nm-thick ZnO films deposited at 200 and 225 °C.

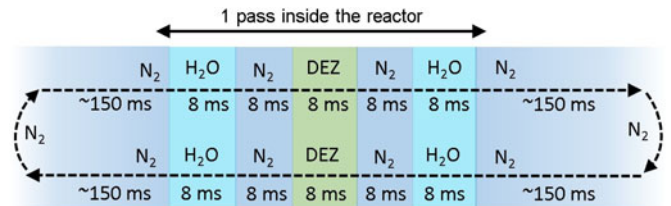


Fig. 7. Schematic highlighting the DEZ and H₂O pulse and purge durations experienced by the substrate during one back-and-forth motion for the spatial ALD technique used in this study.

(where a sample is exposed to consecutive H₂O-purge-DEZ-purge-H₂O-purge-DEZ cycles), in spatial ALD, the sample experiences two H₂O exposures after each DEZ exposure. The DEZ/H₂O exposure times (for a constant wafer speed) are each ~ 8 ms. The purge time between the H₂O and the DEZ exposures is also ~ 8 ms during each wafer pass. Upon completion of one pass, the wafer is pushed back by a pulse of N₂ (as seen in Fig. 1), and this results in a longer N₂ purge followed by another H₂O exposure.

As mentioned earlier, Nelson *et al.* also showed that longer purge times result in lower growth rates, most likely due to desorption of the hydroxyl (OH) groups. If this was the case, the 300-ms-long purge in our system would result in higher desorption of (OH) groups from the surface. However, the subsequent second H₂O exposure refreshes the surface with (OH) groups followed by only a short ~ 8 ms N₂ purge that prevents any further desorption.

Keeping this in mind, the high resistivity of the *i*-ZnO films deposited in this study is easily explained by the results of Nelson *et al.* The DEZ pulse/purge time in our spatial ALD system was ~ 8 ms, much shorter than the 80 ms used in [30]. This also explains the high resistivities obtained by spatial ALD compared with temporal ALD where the exposure and purge durations are of the order of ~ 1 and ~ 10 s, respectively. On the other hand, Illiberi *et al.* reported a low resistivity of $4 \times 10^{-3} \Omega \cdot \text{cm}$ for *i*-ZnO films that were deposited by spatial ALD. This low resistivity value can be explained by noting that in their spatial ALD setup the DEZ/H₂O exposure time is ~ 30 ms and the purge time is ~ 0.3 s, both significantly longer than the time durations used in this study. From an industrial point of view,

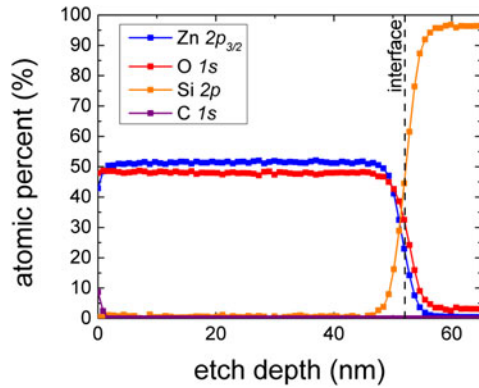


Fig. 8. XPS depth profile of a 50-nm *i*-ZnO film deposited at 225 °C with a DEZ partial pressure of 2 mbar, a water partial pressure of 65 mbar, and pulse and purge durations of ~ 8 ms each.

it should also be noted that longer pulse/purge durations would result in lower deposition rates; hence, shorter pulse/purge durations are more desirable for high volume manufacturing.

The additional carriers in less resistive *i*-ZnO films are often associated with defects due to oxygen vacancies, while higher oxygen content and oxygen interstitials are associated with high resistivity [4], [5], [29], [36], [37]. The observations in this study fit this line of thinking very well. With shorter exposure/purge times, it is expected that the O atoms have a lower probability to desorb from the surface, resulting in fewer oxygen vacancies and hence a lower defect density, which corresponds to a higher resistivity. This is further supported by the lower resistivity obtained for films deposited at 225 °C, as compared with those deposited at 200 °C. It is expected that higher temperature also results in higher densities of vacancies and defects.

D. Material Properties

In Fig. 8, the atomic percentage (at%) of oxygen, zinc, silicon, and carbon as obtained from the XPS depth profiles are presented for an *i*-ZnO film deposited at 225 °C. First, no carbon contamination was detected in the bulk of the film. Second, the relative at% of O is seen to be comparable to that of Zn indicating that the *i*-ZnO films are nearly stoichiometric. This corresponds well with earlier reports for ZnO films prepared by ALD [38] and indicates that there is no significant O/OH deficiency in the films due to the short pulse and purge durations.

The variation in crystal orientation with changing deposition temperature of *i*-ZnO films was measured with XRD; see Fig. 9. The predominant grain orientation of the films changed from (0 0 2) at low deposition temperature (105 °C) to (1 0 0) for high deposition temperature (250 °C). Most reports in the literature on *i*-ZnO deposited by temporal ALD show a predominant grain orientation of (0 0 2) [39], [40]. Nelson *et al.* also found that with shorter exposure/purge times, the (1 0 0) and (1 0 1) peaks seemed to dominate [30], whereas at longer exposure/purge durations, the (0 0 2) orientation dominates. This is consistent with the dominance of the (1 0 0) peak over the (0 0 2) peak in the *i*-ZnO films deposited in this study. The preferred (1 0 0) orientation is consistent with earlier results for

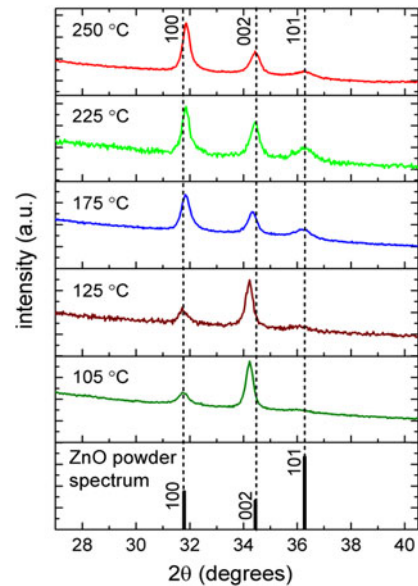


Fig. 9. XRD 2θ scan of *i*-ZnO films grown at different deposition temperatures, showing a change in the preferential crystal orientation. The *i*-ZnO powder spectrum has been included for reference.

i-ZnO deposited by spatial ALD [17]. It was proposed that the longer exposure/purge times increase the ability of the precursor molecule to find low-energy sites on the surface.

IV. CONCLUSION

Using spatial ALD, the deposition of intrinsic ZnO films was scaled up resulting in a high deposition rate of ~ 35 nm/min on large-area (125 mm \times 125 mm) *c*-Si substrates. The *i*-ZnO films have a high resistivity ($\sim 100 \Omega \cdot \text{cm}$) and a refractive index that results in a high transparency ($>90\%$) in the visible range. The high resistivities of the *i*-ZnO films deposited in this study are attributed to the short precursor exposure and purge times of the spatial ALD technique, as compared with other deposition techniques. XPS results showed that the *i*-ZnO films were nearly stoichiometric. The preferential crystallographic orientation of the *i*-ZnO films was found to be (1 0 0) at higher deposition temperatures, rather than (0 0 2) as reported for conventional temporal ALD ZnO films. This is also attributed to the short exposure/purge times. These properties make these *i*-ZnO films ideal for application as the window layer between the CdS buffer and the conductive front contact in CIGS solar cells in order to reduce the influence of absorber layer defects from affecting the cell performance. However, one should take into account that ZnO would grow differently on CdS as compared to *c*-Si, resulting in potentially varying material properties. For applications of ZnO as a transparent conductive oxide in CIGS and silicon heterojunction solar cells, lower resistivities are desired and can be achieved by impurity doping, e.g., Al-, B-, Ga-, or In-doped ZnO. This preliminary study on *i*-ZnO paves the way for its application in existing and new PV technologies, while the advantages of the spatial ALD technique described in this study make it an attractive option for the PV industry.

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