

Optical spectroscopy of single beryllium acceptors in GaAs/AlGaAs quantum well

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Optical spectroscopy of single beryllium acceptors in GaAs/AlGaAs quantum wellP. V. Petrov,^{1,*} I. A. Kokurin,^{1,2} G. V. Klimko,¹ S. V. Ivanov,¹ Yu. L. Ivánov,¹
P. M. Koenraad,³ A. Yu. Silov,³ and N. S. Averkiev¹¹*Ioffe Institute, St. Petersburg, Russian Federation*²*Institute of Physics and Chemistry, Mordovia State University, Saransk, Russian Federation*³*Department of Applied Physics, Eindhoven University of Technology, The Netherlands*

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We carry out microphotoluminescence measurements of an acceptor-bound exciton (A^0X) recombination in the applied magnetic field with a single impurity resolution. In order to describe the obtained spectra we develop a theoretical model taking into account a quantum well (QW) confinement, an electron-hole and hole-hole exchange interaction. By means of fitting the measured data with the model we are able to study the fine structure of individual acceptors inside the QW. The good agreement between our experiments and the model indicates that we observe single acceptors in a pure two-dimensional environment whose states are unstrained in the QW plain.

DOI: [10.1103/PhysRevB.94.115307](https://doi.org/10.1103/PhysRevB.94.115307)**I. INTRODUCTION**

Studies of single impurities in solids is one of the most rapidly developing fields of experimental physics in recent years [1–3]. Such experiments are attractive since it makes it possible to verify the fundamental theoretical approaches that were based on macroscopic measurements. In the field of applied science a device based on single impurities is the ultimate limit of electronics miniaturization. At present, two main techniques are exploited in order to reach a single impurity resolution: scanning tunneling microscopy [4] and microphotoluminescence. On the one hand the optical methods have advantage over tunneling measurements due to the absence of the surface influence. On the other hand the resolution of optical measurements is fundamentally restricted by the diffraction limit. The photon wavelength has to be smaller than the average distance between impurities. The band gap of typical semiconductors is about 1 eV, and therefore the corresponding doping concentration should not exceed 10^{12} cm^{-3} for three-dimensional (3D) or 10^8 cm^{-2} for two-dimensional (2D) systems. At present, the spectroscopy of single semiconductor nanostructures such as quantum dots (QDs) is well developed [5]. The obvious approach is to dope a single QD with an impurity atom. Experiments of this kind were realized for CdTe [2] and InAs [3] QDs doped with Mn. However, an interpretation of experimental results in QD systems is hampered by the fact that such parameters as a dot size, shape, chemical composition, as well as an impurity position inside the QD are randomly distributed across QD ensemble. It makes it necessary to use a lot of additional parameters in the theoretical description of experimental results [6,7].

In the present paper, we study a narrow GaAs/AlGaAs quantum well (QW) doped with beryllium in order to optically explore single impurities. Usually, the single emitters in such systems are studied via submicron apertures or mesa structures formed on a sample surface [8,9]. To reach the single impurity resolution here, we do not use any preprocessing of the samples but optimize the doping process instead. The smallest

controllable sheet impurity density in our experiments is about 10^{10} cm^{-2} . This number does not meet the diffraction limit condition, but nonetheless can serve a purpose in the same way as was first realized in the spectroscopy of single organic molecules [10]. The point is to put emitters in a media that randomly changes the emitters energy and to employ a spectral resolution in addition to the spatial one. It is well known that fluctuations in a QW width lead to an inhomogeneous spectral broadening of the exciton photoluminescence due to significant variations of the effective band gap [11]. Assuming that the energy broadening corresponds to a Gaussian shape of photoluminescence line, let us consider the low-energy tail of spectrum. For the Gaussian distribution a probability that the transition energy is in the range between two and three standard deviations from the distribution maximum is about 1%. Therefore we can reach the necessary small sheet density 10^8 cm^{-2} of impurity related single optical emitters, if we examine a lower-energy tail in the photoluminescence of an inhomogeneously broadened ensemble of the impurities.

The interface roughness leads to lateral asymmetry and affects the energy structure of excitons [8,12]. But if a radius of an impurity-bound exciton is smaller than a scale of the roughness, we can neglect the lateral asymmetry and consider such an exciton in a pure 2D environment. This allows us to significantly reduce the number of fitting parameters in comparison with the case of doped QDs [6,7].

This paper is organized as follows: we describe the sample growth, the characterization procedure, and microphotoluminescence data in Sec. II. In Sec. III we present a theoretical model of the acceptor-bound exciton which includes the QW confinement. A comparison of theoretical calculation with obtained and previously published experimental data is discussed in Sec. IV.

II. EXPERIMENT

We grew by molecular beam epitaxy three GaAs/Al_xGa_{1-x}As QW structures with a QW width of 3.7 nm and the Al content in the barriers $x = 0.25$. The samples were doped with Be acceptors inside the QW. They have similar design and differ only in the Be doping mode and sheet impurity density, as shown in Table I. In two samples the QW is δ doped in

*pavel.petrov@gmail.com

TABLE I. Sample parameters: T_{sub} and T_{Be} are the temperature of the substrate and Be source correspondingly, $\delta 3''$ stands for 3 sec of δ doping, 1 ML means one-half of the lattice constant, N_s is an expected acceptor sheet concentration inferred from the beryllium source calibration.

Sample no.	Substrate $T_{\text{sub}} = 560\text{--}580^\circ\text{C}$	GaAs			Al _x Ga _{1-x} As		Al _x Ga _{1-x} As		
		buffer	$x = 0.25$	GaAs	GaAs	GaAs:Be	GaAs	$x = 0.25$	GaAs
S1	undoped GaAs	0.25 μm	100 nm	7 MLs	$\delta 3''$, $T_{\text{Be}} = 660^\circ\text{C}$, $N_s = 5 \times 10^9 \text{ cm}^{-2}$	6 MLs	100 nm	20 nm	
S2	<i>p</i> -type GaAs	0.25 μm	100 nm	7 MLs	$\delta 3''$, $T_{\text{Be}} = 690^\circ\text{C}$, $N_s = 5 \times 10^{10} \text{ cm}^{-2}$	6 MLs	100 nm	20 nm	
S3	undoped GaAs	0.25 μm	100 nm	1 ML	11 MLs, $T_{\text{Be}} = 660^\circ\text{C}$, $N_s = 4 \times 10^{10} \text{ cm}^{-2}$	1 MLs	100 nm	20 nm	

the middle, while in the third we use uniform doping of the QW with 1-monolayer (ML)-thick undoped spacers at both interfaces. We adjusted the barrier height in order to ensure an effective band-to-band absorption of a pumping light inside the barriers.

We use macrophotoluminescence measurements at 4.2 K in order to characterize the grown samples and to establish the presence of beryllium inside the QW. The samples were pumped with a 660-nm diode laser via an optical fiber with a cross section of 0.1 mm²; macroluminescence spectra were collected through the same fiber. We expect that the beryllium dopant reveals itself as an additional low-energy broadening of the QW related photoluminescence line due to an appearance of the acceptor-bound excitons [13]. Samples S2 and S3 with more intensive doping indeed demonstrate the expected low-energy broadening as shown in Fig. 1. The microphotoluminescence measurements show that these tails consist of numerous narrow lines. In order to distinguish the microphotoluminescence lines due to the acceptor-bound excitons from the lines of different origin, we use the rich energy structure of A^0X complex [14] as a spectral fingerprint.

We carry out microphotoluminescence measurements at 5 K on the setup with about 1 μm spatial and $\sim 60\text{-}\mu\text{eV}$ energy resolution using HeNe laser as a pump source. Figure 2 shows three characteristic spectra which were measured at different spots on the surface of our samples. The strong photoluminescence line at 1.65–1.66 eV corresponds to recombination of excitons inside the QW (X_{QW}) and it looks the same for all the samples, while the low-energy side of the spectrum presents a wide variety of results. We observe mostly spectra of type (a) on the sample S1 with the lowest doping concentration. There are no evidences of impurities at the low-energy tail in panel (a). Sample S3 shows a strong nonuniformity across the surface: most of the sample surface corresponds to type (a) while less than 10% of the surface gives the spectra of type (c). The spectrum in panel (c) contains an impurity related

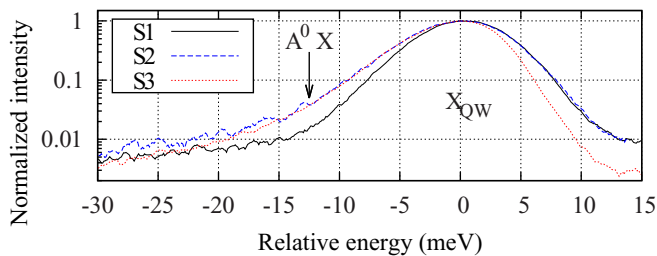


FIG. 1. Normalized macrophotoluminescence spectra of the studied samples. Spectra are centered at photoluminescence maxima for comparison; the pump density is 10 W/cm².

luminescence at the low-energy side, but the luminescence lines are quite broad and overlap each other. Sample S2 is the one most suitable for microphotoluminescence measurements. We observe the spectra of type (b) with strong narrow lines at the low-energy tail of the sample S2 photoluminescence. Below in the text we discuss results which were obtained on this sample.

In trying to find out if there are any distinctive peculiarities in the microphotoluminescence spectra, we carry out measurements in the applied magnetic field in Faraday geometry. Among manifold combinations of luminescence lines we observe a kind of repeating pattern in polarized photoluminescence spectra. It consists of one strong single luminescence line and two weak adjacent satellites which are split by magnetic field in doublets denoted as (1), (2), and (3) in Fig. 3. It is noteworthy that a Zeeman splitting of the satellites is 1.5–2 times stronger than a splitting of the main line. In order to give a reliable interpretation of the results we develop a theoretical model of an acceptor-bound exciton in which we take into account an interparticle exchange [15,16], the QW confinement [14], and the magnetic field.

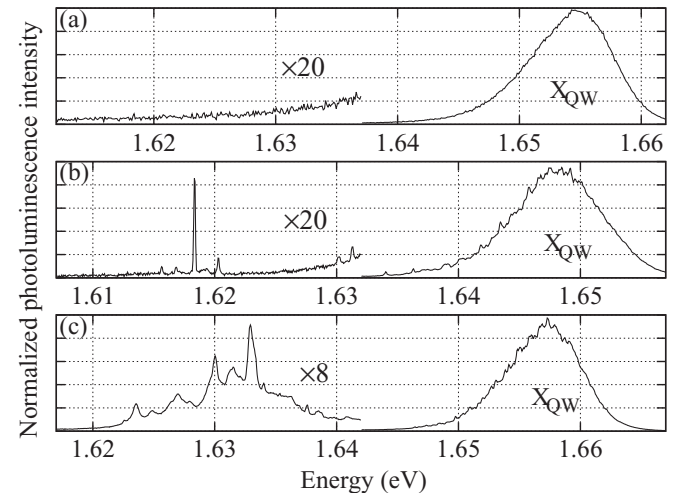


FIG. 2. Three typical microphotoluminescence spectra of the studied samples that were measured at different spots on the samples. The spectrum in panel (a) contains only broad line of excitons (X_{QW}) which recombine inside the QW; no extra features present at the low-energy tail. Narrow, well-resolved lines are present on the low-energy side of the spectrum of type (b). Panel (c) depicts spectrum which contains numerous overlapping lines at the region of interest.

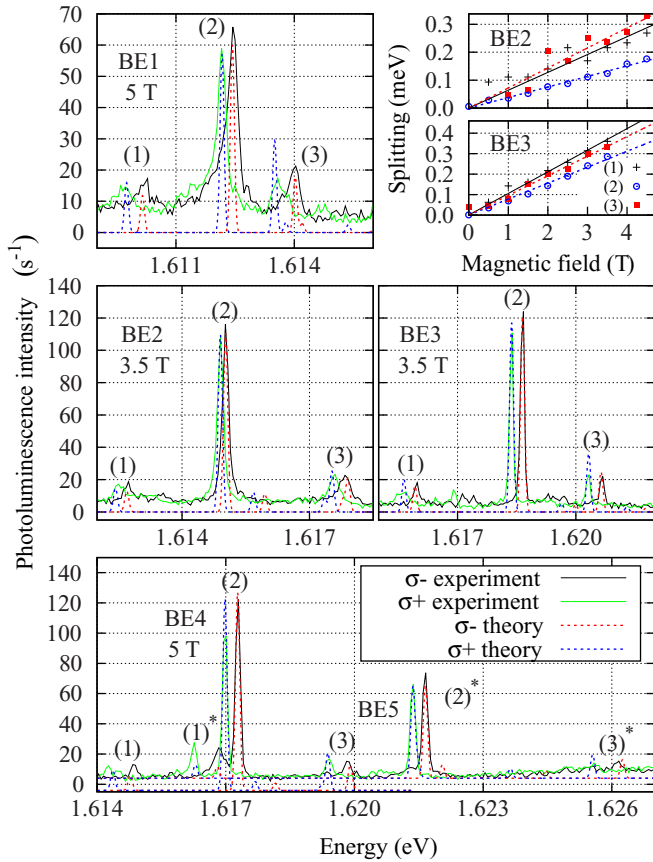


FIG. 3. Experimental (solid line) and theoretical (dashed line) polarized microphotoluminescence spectra of single acceptor-bound excitons in the applied magnetic field. Small panels display the Zeeman splitting. Symbols are experimental data for doublets (1), (2), and (3); solid lines are the linear least-squares fitting.

III. THEORY

In order to obtain an energy structure of acceptor-bound exciton A^0X inside a QW, we use a model Hamiltonian:

$$H = -\Delta^{\text{hh}} \mathbf{J}_1 \cdot \mathbf{J}_2 - \Delta^{\text{eh}} \mathbf{S} \cdot (\mathbf{J}_1 + \mathbf{J}_2) + \frac{\Delta^{\text{qw}}}{2} \left(J_{1z}^2 + J_{2z}^2 - \frac{5}{2} \right), \quad (1)$$

where Δ^{hh} and Δ^{eh} are the hole-hole and the electron-hole exchange energies, respectively, and Δ^{qw} is a splitting of the localized hole state due to a QW confinement. Here \mathbf{J}_i ($i = 1, 2$) and \mathbf{S} stand for an angular momentum of the holes and the electron, respectively. We use a spherical model of the localized hole states [17] and consider only the ground state with momentum $J = 3/2$. The wave function of two indistinguishable holes must be antisymmetric, therefore only states with total angular momentum $J = 0, 2$ are present. In diamondlike semiconductors the hole-hole exchange is a ferromagnetic interaction ($\Delta^{\text{hh}} > 0$), therefore a state with the largest total angular momentum is the ground one [18]. The electron-hole exchange interaction between these two holes and the electron with $S = 1/2$ leads to the emergence of a three-particle complex with total angular momentum $F = 1/2, 3/2, 5/2$. This interaction is also ferromagnetic ($\Delta^{\text{eh}} > 0$)

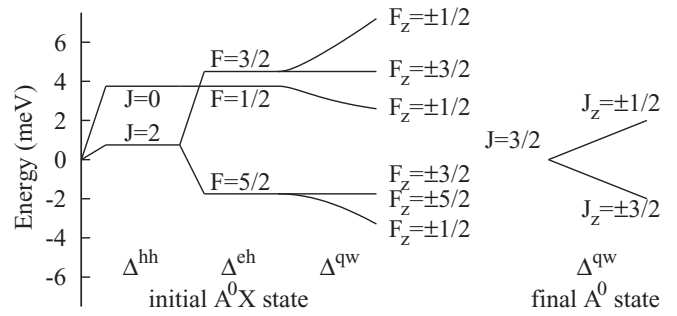


FIG. 4. The scheme of energy levels for initial A^0X and final A^0 states in presence of exchange interaction and QW confinement.

in GaAs/AlGaAs QWs [19]; it means that the “dark” state of the free exciton is the ground state. An energy splitting due to the QW confinement is negative ($\Delta^{\text{qw}} < 0$) which corresponds to $J_z = \pm 3/2$ as a ground hole state. The order of levels in the bulk A^0X complex ($\Delta^{\text{qw}} = 0$) depends on the ratio between Δ^{hh} and Δ^{eh} . We obtain all the three-particle wave functions $\Psi_{\pm 5/2}^{5/2}$, $\Psi_{\pm 3/2}^{5/2}$, $\Psi_{\pm 1/2}^{5/2}$, $\Psi_{\pm 3/2}^{3/2}$, $\Psi_{\pm 1/2}^{3/2}$, and $\Psi_{\pm 1/2}^{1/2}$ analytically using the usual procedure of angular momentum coupling [20]. Here the upper index is a full angular momentum of the state while the lower one is its projection. The energy levels of A^0X complex are given by solution of Schrödinger equation with Hamiltonian (1) at $\Delta^{\text{qw}} = 0$,

$$E_{1/2} = \frac{15}{4} \Delta^{\text{hh}}, \quad E_{3/2} = \frac{3}{4} \Delta^{\text{hh}} + \frac{3}{2} \Delta^{\text{eh}}, \quad E_{5/2} = \frac{3}{4} \Delta^{\text{hh}} - \Delta^{\text{eh}}. \quad (2)$$

The QW potential leads to the mixing of levels with angular momentum projection $F_z = \pm 1/2$ keeping other levels constant. Using $\Psi_{\pm 1/2}^{1/2}$, $\Psi_{\pm 1/2}^{3/2}$, $\Psi_{\pm 1/2}^{5/2}$ functions as a basis we can write the Hamiltonian for three states with $F_z = \pm 1/2$:

$$H = \begin{pmatrix} \frac{15}{4} \Delta^{\text{hh}} & -\sqrt{\frac{2}{5}} \Delta^{\text{qw}} & \sqrt{\frac{3}{5}} \Delta^{\text{qw}} \\ -\sqrt{\frac{2}{5}} \Delta^{\text{qw}} & \frac{3}{4} \Delta^{\text{hh}} + \frac{3}{2} \Delta^{\text{eh}} & 0 \\ \sqrt{\frac{3}{5}} \Delta^{\text{qw}} & 0 & \frac{3}{4} \Delta^{\text{hh}} - \Delta^{\text{eh}} \end{pmatrix}. \quad (3)$$

We obtain the energy levels E_i ($i = 1, 2, 3$) and the corresponding wave functions $\Psi_{\pm 1/2}^{(i)} = a_{1/2}^i \Psi_{\pm 1/2}^{1/2} + a_{3/2}^i \Psi_{\pm 1/2}^{3/2} + a_{5/2}^i \Psi_{\pm 1/2}^{5/2}$ as a solution of the Hamiltonian (3) eigenvalue problem. Figure 4 depicts the obtained energy scheme of A^0X complex.

Assuming that Zeeman energy is much smaller than all the energy parameters of the system, we find the Zeeman splitting of A^0X levels in the first order of perturbation theory. For simplicity's sake we take the electron g factor $g_e = 0$, which is true for narrow GaAs/AlGaAs QWs [21]. Zeeman splitting is described by the Hamiltonian

$$H_Z = \mu_B g_{hi} B (J_{1z} + J_{2z}), \quad (4)$$

where g_{hi} is a hole g factor of the initial A^0X state. Let us obtain all the g factors of the acceptor-bound exciton states

normalized to the angular momentum $1/2$:

$$g_{3/2}^{3/2} = \frac{18}{5}g_{hi}, \quad g_{5/2}^{5/2} = 4g_{hi}, \quad g_{3/2}^{5/2} = \frac{12}{5}g_{hi}.$$

The g factors of mixed states depend on the eigenvector coefficients a_j^i ($i = 1, 2, 3, j = 1/2, 3/2, 5/2$):

$$\frac{g_i}{g_{hi}} = \frac{2}{5}(\sqrt{3}a_{3/2}^i + \sqrt{2}a_{5/2}^i)^2. \quad (5)$$

The final state after the A^0X recombination is a neutral acceptor A^0 . The final state is also split by the QW potential [22] with the same Δ^{qw} :

$$H = \frac{\Delta^{qw}}{2} \left(J_z^2 - \frac{5}{4} \right). \quad (6)$$

The four nondegenerate states of A^0 produced by a magnetic field are

$$E_{\pm 3/2} = \frac{\Delta^{qw}}{2} \pm \frac{3}{2}g_{hf}\mu_B B, \quad (7)$$

$$E_{\pm 1/2} = -\frac{\Delta^{qw}}{2} \pm \frac{1}{2}g_{hf}\mu_B B, \quad (8)$$

where g_{hf} is a g factor of the final A^0 state.

Knowing the energy of the initial E_i and final E_f states we can establish all transition energies as

$$\hbar\omega = E_g + E_i - E_f, \quad (9)$$

where E_g is an effective band gap including all the confinement shifts and the exciton binding energy. In order to obtain oscillator strengths and polarizations of the transitions we use the usual selection rules combined with Clebsch-Gordan coefficients that couple spins of A^0X complex.

IV. DISCUSSION

Figure 3 shows a set of the circularly polarized microphotoluminescence spectra which were measured at different spots on the sample in applied magnetic field. As mentioned above, all the spectra match a repetitive pattern: a strong line with two accompanying satellites. According to our model the strongest photoluminescence lines which are denoted as (1), (2), and (3) correspond to transitions $\Psi_{\pm 3/2}^{5/2} \rightarrow J_z = \pm 1/2$, $\Psi_{\pm 1/2}^{(1)} \rightarrow J_z = \pm 3/2$, and $\Psi_{\pm 5/2}^{5/2} \rightarrow J_z = \pm 3/2$, respectively. The most intense line (2) originated from the ground $\Psi_{\pm 1/2}^{(1)}$ state of A^0X complex while satellites are due to the subsequent degenerate $\Psi_{\pm 3/2}^{5/2}, \Psi_{\pm 5/2}^{5/2}$ state. The energy spacing between (1) and (3) lines is equal to the Δ^{qw} parameter of our model. Sets of the fitting parameters of all spectra presented in Fig. 3 are compiled in Table II under the labels BE1–BE5.

Another characteristic feature of our model comes from the fact that a radiative recombination of acceptor-bound exciton occurs via transitions between a multiplet of initial states and only two available final states. It means that a few pairs of the spectral line with an equal spacing of Δ^{qw} can be present in the spectrum of the acceptor-bound exciton recombination. Figure 5 depicts an experimental spectrum which contains three pairs with the similar energy splitting of $\Delta^{qw} = 4.47$ meV. Assuming that the most intensive line

TABLE II. Fitting parameters used for calculation of spectra in Figs. 3 and 5.

No.	Δ^{hh} (meV)	Δ^{eh} (meV)	Δ^{qw} (meV)	g_{hi}	g_{hf}	T (K)
BE1	1.0	2.3	-3.81	0.75	0.4	20
BE2	0.8	2.0	-5.15	0.65	0.25	30
BE3	1.2	2.3	-4.70	0.8	0.55	20
BE4	1.0	2.0	-5.07	0.75	0.4	20
BE5	2.5	3.0	-9.35	0.9	0.4	50
BE6	0.37	2.27	-4.47			60

in the spectrum is a $\Psi_{\pm 1/2}^{(1)} \rightarrow J_z = \pm 3/2$ transition we successfully describe all other optical transitions using our three-parameter fit at zero magnetic field. In order to fit the transition intensities we also take into account an equilibrium probability $\sim \exp(-E_i/kT)$ to find the A^0X complex in a certain initial state using an effective bath temperature T as a fourth parameter. The corresponding fit parameters are listed in Table II denoted by BE6.

Let us compare the values of parameters with published results of other experiments. Using Eqs. (2) we extract the values $\Delta_{3D}^{hh} = 0.11$ meV and $\Delta_{3D}^{eh} = 0.06$ meV from A^0X photoluminescence data obtained on the bulk GaAs material [15]. It is well known that a quantum confinement significantly enhances the electron-hole exchange in nanostructures [23,24], therefore our fitted values of exchange parameters Δ^{hh} and Δ^{eh} seem reasonable. The typical effective temperature of recombining excitons is about 20 K in narrow GaAs/AlGaAs QWs [25] in accordance with our results. We obtain relatively high effective temperature $T \sim 50$ K for a couple of measured spectra which means that the local exciton lifetime can be comparable to the time of energy relaxation. The QW splitting Δ^{qw} and g factor of neutral acceptors were directly measured via spin-flip Raman scattering [26]. Our g factor values are comparable with those from [26]; the discrepancy is due to using the model fit instead of direct measurement. We have a

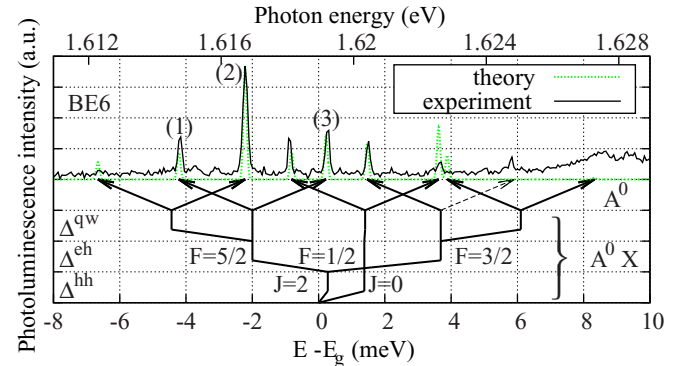


FIG. 5. Experimental (solid line) and theoretical (dashed line) microphotoluminescence spectra of an exciton bound to the single beryllium acceptor. The level scheme depicts splitting of the initial A^0X state due to the hole-hole Δ^{hh} , the electron-hole Δ^{eh} exchange, and the QW confinement potential Δ^{qw} ; arrows indicate a splitting in the final A^0 state. The dashed arrow marks a forbidden transition, which becomes available if we take into account the cubic anisotropy of the crystal.

good agreement of Δ^{qw} values with the data from [26] if we take into account the strong fluctuation of Δ^{qw} depending on the acceptor position with respect to the barrier.

Such a strong dependence of Δ^{qw} on z coordinate of an individual acceptor makes it possible to establish a position of impurity in the growth direction. Lateral coordinates of an impurity in the quantum well could be established within 1-nm accuracy using a super-resolution optical technique which is well developed for the single molecule spectroscopy [27]. An application of these methods could provide unprecedented possibilities to establish exact atomic coordinates of impurities inside the crystal lattice and to explore its spin and energy structure, combining advantages of optical spectroscopy with the ultimate spatial accuracy of the scanning tunneling microscopy.

In conclusion, we report photoluminescence measurements of excitons bound to single beryllium acceptors in

GaAs/AlGaAs QWs. In order to describe our results we use a simple theoretical model of an acceptor-bound exciton confined in the QW. The model includes the interparticle exchange. The obtained parameter values of our model are in a good agreement with previously published data and accurately describe a complex spectral signature of the single impurity in radiative recombination.

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